

Opposing Intermolecular Tuning of Ca^{2+} Affinity for Calmodulin by Neurogranin and CaMKII Peptides

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ABSTRACT We investigated the impact of bound calmodulin (CaM)-target compound structure on the affinity of calcium $(Ca²⁺)$ by integrating coarse-grained models and all-atomistic simulations with nonequilibrium physics. We focused on binding between CaM and two specific targets, Ca^{2+}/CaM -dependent protein kinase II (CaMKII) and neurogranin (Ng), as they both regulate CaM-dependent Ca²⁺ signaling pathways in neurons. It was shown experimentally that Ca²⁺/CaM (holoCaM) binds to the CaMKII peptide with overwhelmingly higher affinity than Ca^{2+} -free CaM (apoCaM); the binding of CaMKII peptide to CaM in return increases the Ca^{2+} affinity for CaM. However, this reciprocal relation was not observed in the Ng peptide $(Ng_{13–49})$, which binds to apoCaM or holoCaM with binding affinities of the same order of magnitude. Unlike the holoCaM-CaMKII peptide, whose structure can be determined by crystallography, the structural description of the apoCaM-Ng_{13–49} is unknown due to low binding affinity, therefore we computationally generated an ensemble of apoCaM-Ng_{13–49} structures by matching the changes in the chemical shifts of CaM upon Ng_{13-49} binding from nuclear magnetic resonance experiments. Next, we computed the changes in Ca^{2+} affinity for CaM with and without binding targets in atomistic models using Jarzynski's equality. We discovered the molecular underpinnings of lowered affinity of Ca²⁺ for CaM in the presence of Ng_{13–49} by showing that the N-terminal acidic region of Ng peptide pries open the β -sheet structure between the Ca²⁺ binding loops particularly at C-domain of CaM, enabling Ca^{2+} release. In contrast, CaMKII peptide increases Ca^{2+} affinity for the C-domain of CaM by stabilizing the two Ca²⁺ binding loops. We speculate that the distinctive structural difference in the bound complexes of apoCaM-Ng_{13–49} and holoCaM-CaMKII delineates the importance of CaM's progressive mechanism of target binding on its Ca^{2+} binding affinities.

INTRODUCTION

Calcium (Ca^{2+}) is exquisitely used by a cell for transducing external stimuli through its gradient of extracellular (~1000 μ M) and intracellular (~0.1 μ M) concentration ([1\)](#page-12-0). A broad spectrum of Ca^{2+} signals are encoded by protein calmodulin (CaM) (2) (2) (the structure is shown in [Fig. 1](#page-1-0) a) through specific binding with various targets regulating CaM-dependent Ca²⁺ signaling pathways in neurons ([3\)](#page-12-0). The targets enhance CaM's affinity for Ca^{2+} by up to 30 times or accelerate dissociation of Ca^{2+} from CaM ([4\)](#page-12-0). Therefore, this mechanism tunes CaM's capacity of encoding global Ca^{2+} signals by its differential binding affinity for Ca^{2+} . Among those targets, Ca^{2+}/CaM -dependent protein kinase II (CaMKII) $(5-7)$ (the structure and sequence of the CaM-binding domain are shown in [Fig. 1](#page-1-0) b) and neurogranin (Ng) (8) (8) (the sequence of the Ng_{13–49} peptide is shown in [Fig. 1](#page-1-0) c) play an essential role in synaptic plasticity

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([9\)](#page-12-0), which is critical to learning and memory formation shown by early behavior experiments on Ng-mutated or Ng-knocked-out mice ([10,11](#page-12-0)). We summarized several experimental observations [\(3,7](#page-12-0)) in Fig. S1 on how CaM sensitizes downstream CaM-binding enzymes to Ca^{2+} regulating synaptic plasticity in memory formation: CaM activates CaMKII at high Ca^{2+} levels [\(6](#page-12-0)) while the availability of CaM is controlled by a group of postsynaptic proteins including Ng at a lower Ca^{2+} level ([8,12](#page-12-0)).

Waxham et al. $(13,14)$ $(13,14)$ $(13,14)$ used fluorescence experiments to show that the CaMKII kinase increases the affinity of Ca^{2+} for CaM while Ng significantly decreases the affinity of Ca²⁺ for Ca²⁺-binding loop III and IV ([Fig. 1](#page-1-0) *a*) in C-domain of CaM (cCaM) by up to 60-fold. These experiments raised questions about the dependence of CaM's affinity for Ca^{2+} ions on structural models of CaM [\(13,15\)](#page-12-0). Recent studies by Le Novère et al. $(16,17)$ $(16,17)$ $(16,17)$ demonstrated the alteration of Ca^{2+} affinity for CaM by plotting the binding affinity curves with elaborated allosteric models and fitted rates. They investigated the effect of competing targets on CaM by showing the enhancement of binding affinity in

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FIGURE 1 Illustration of structures of CaM and sequences of the target peptides. (a) Solution structure of Ca^{2+} -free CaM (apoCaM) (PDB: 1CFD) and crystal structure of Ca^{2+}/CaM (holoCaM) (PDB: 1CLL). CaM is colored as follows, red \rightarrow nCaM (residue 1–76), gray \rightarrow central linker (residue 77–82), blue \rightarrow cCaM (residue 83–148), and the Ca²⁺ ions are colored in green; Ca^{2+} binding sites are labeled from I to IV. (b) Structure of Ca^{2+}/CaM -dependent kinase II (CaMKII) peptide from the holoCaM-CaMKII crystal structure (PDB: 1CDM) and the sequence of the peptide. (c) X-ray structure of holoCaM-CaMKII (PDB: 1CDM). CaM and the Ca^{2+} ions are colored as in (*a*) and the CaMKII peptide is colored in green. (d) Sequence of the neurogranin peptide (Ng₁₃₋₄₉). The two minimum compositions of the Ng_{13-49} peptide are marked as "acidic region" and "IQ motif.'' To see this figure in color, go online.

the presence of some targets (such as CaMKII), as well as the diminishment in the presence of other targets (such as Ng). Dieckmann et al. ([18\)](#page-12-0) used steady-state fluorescence and nuclear magnetic resonance (NMR) experiments and found that the binding of endothelial nitric oxide synthase (eNOS) to CaM enhances CaM's affinity for Ca^{2+} while phosphorylated eNOS does not due to electrostatic repulsion. They speculated that the phosphorylation on eNOS diminished its helical propensity, possibly affecting CaM's affinity for Ca^{2+} . On the other hand, the structural basis of the regulation of Ca^{2+} affinity for CaM by Ng and CaMKII remains elusive.

CaM has two domains separated by a flexible central linker. Each domain consists of two helix-linker-helix motifs (marked as I–IV in Fig. 1 a), which can bind a $Ca²⁺$ ion. The flexible central linker allows adjustable CaM conformations, thus accommodates the binding of distinct CaM binding targets (CaMBTs) ([19–23\)](#page-12-0). In a canonical complex of Ca^{2+}/CaM (holoCaM) and CaMKII peptide, the two domains of holoCaM wrap around the CaMKII peptide, which forms a helical structure [\(24](#page-12-0)). In our previous work, by using a combined approach of coarse-grained molecular simulations and experiments, we have established the binding mechanism of holoCaM and CaMKII peptide [\(25,26](#page-12-0)) that follows ''conformational and mutually induced fit" [\(27](#page-12-0)). Both holoCaM and CaMKII undergoes large conformational changes in the two-step binding mechanism.

Unlike the CaMKII peptide in complex with holoCaM, which was solved by X-ray crystallography decades ago ([24\)](#page-12-0), the weak binding between Ng and Ca^{2+} -free CaM (apoCaM) makes the determination of its three-dimensional complex structure challenging [\(28](#page-12-0)). Ng belongs to the IQ motif family. In this family, the CaMBTs share a common motif (IQXXXRGXXXR), where X stands for any amino acid. Recently, Kumar et al. [\(29](#page-12-0)) determined the crystal structure of apoCaM and IQ motif of Ng (NgIQ), however, NgIQ was artificially covalently linked with cCaM in the bound structure. Furthermore, a previous experimental study [\(30](#page-12-0)) revealed that the IQ motif alone cannot represent the biophysical characteristics of the intact Ng protein. Inclusion of the acidic region in the N-terminal before the IQ motif (or Ng_{13-49} , which are largely composed of Asp or Glu residues, reproduced the Ng-mediated affinity of Ca^{2+} to apoCaM, as well as the intermolecular interactions between the intact Ng and apoCaM. However, the main obstacle to investigate apoCaM-Ng₁₃₋₄₉ binding is a lack of experimentally solved structures because of their weak binding affinity (approximately micromolar) [\(30](#page-12-0)). To determine the apoCaM-Ng13–49 complex structures, we modeled the interaction between intrinsically disordered Ng_{13-49} and apoCaM by introducing statistical dihedral angle potentials of Ng_{13-49} ([31\)](#page-12-0). After obtaining the appropriate coarse-grained models of apoCaM-Ng_{13–49}, we reconstructed them into all-atomistic structures (32) (32) . We noted that in comparison to the canonical complex of CaM and CaMBT that includes the two domains of CaM wrapping around a helical CaMBT, the conformation of CaM in the complexes of apoCaM- Ng_{13-49} remains extended instead. The conformations of Ng in the bound complexes are not entirely helical.

The reconstructed atomistic models of the holoCaM- Ng_{13-49} complex and the crystal structure of holoCaM-CaMKII complex (Protein Data Bank (PDB): 1CDM) allow us to develop a hypothesis that the conformation of the Ca^{2+} binding loops from cCaM in a holoCaM-CaMBT complex dictates cCaM's affinity for Ca^{2+} . To validate this hypothesis, we computationally applied all-atomistic steered molecular dynamics (MD) simulations of holoCaM-CaMBT [\(33](#page-12-0)) and used Jarzynski's equality (JE) ([34\)](#page-12-0) to calculate the change in Ca^{2+} binding free energy. As a reference, the same simulations were performed on a conformation of holoCaM without a CaMBT (PDB: 1CLL).

One of the challenges is the modeling of the bivalent Ca^{2+} ions. It has been shown that the charge on Ca^{2+} is

 \lt +2e due to charge transfer and polarization effects in the solution and especially in the Ca^{2+} -bound state ([35–37\)](#page-12-0). Moreover, Jungwirth et al. ([36,38](#page-12-0)) suggested that both the size and the charge require adjustment to account for the polarization effects. Following our previous work ([39\)](#page-12-0), we computed the charges of calcium ions according to the protonation states of CaM by employing a semiempirical quantum chemistry program MOPAC [\(40](#page-12-0)). Thereby we revealed a mechanism that the presence of Ng peptide destabilizes the Ca^{2+} binding to cCaM, whereas the CaMKII peptide stabilizes the Ca^{2+} binding to cCaM. We speculate that CaM's progressive mechanism of target binding, which yields distinctive complexes of apoCaM-Ng13–49 and holoCaM-CaMKII, modulates Ca^{2+} affinities for CaM.

MATERIALS AND METHODS

Coarse-grained protein or peptide models

We performed coarse-grained molecular simulations of the intact Ng protein as well as the binding between apoCaM and Ng13–49 peptide. Sample preparation of the intact Ng or the Ng_{13-49} peptide is provided in Section I of the [Supporting Material.](#page-11-0) The Hamiltonian of the coarse-grained protein models including Ng protein and the apoCaM-Ng13–49 complex is described in Section II of the [Supporting Material.](#page-11-0) Ca^{2+} ions are absent from the coarse-grained models. The details about the coarse-grained molecular simulations are included in Section III of the [Supporting Material](#page-11-0).

Free energy calculation using JE

Preparation of the structures

The preparation of the coordinates as initial conditions for the JE calculation is highly nontrivial for the computation of Ca^{2+} binding free energy ΔG . For both systems, we determined the accurate protonation states by using H_{++} server ([41\)](#page-12-0) according to the protein conformation, pH, and the ionic strength of the solution. We then assigned the partial charges based on the geometry of the proteins by using a semiempirical quantum chemistry program MOPAC ([40\)](#page-12-0). We applied those protonated protein structures and partial charges for minimization and further procedures to carry out all-atomistic calculations of the Ca^{2+} binding free energy (find the details in Section IV of the [Supporting Material\)](#page-11-0). For holoCaM and holoCaM-CaMKII, which have experimentally determined structures, we fixed the heavy atoms during the preparation; for holoCaM-Ng, we constrained the atoms with harmonic potentials for the preparation of the structures (holoCaM: calcium-bound calmodulin).

holoCaM-CaMKII and holoCaM. We fixed the positions of Ca^{2+} atoms and backbone heavy atoms from holoCaM (PDB: 1CLL) or holoCaM-CaMKII (PDB: 1CDM) during the minimization, solvation, ionization, and the equilibration procedures (see Section IV of the [Supporting Material](#page-11-0) for details).

 $holoCaM-Ng_{13–49}$. Since there is no experimentally determined structure, we reconstructed the all-atomistic structures from the selected coarsegrained side-chain- C_{α} models of apoCaM-Ng_{13–49} using SCAAL method [\(32](#page-12-0)) (see Section IV.1 of the [Supporting Material](#page-11-0) for the criteria of selecting apoCaM-Ng₁₃₋₄₉ models). We estimated the position of each $Ca²⁺$ ion as the center of mass of the side chains of the corresponding Ca^{2+} -coordinating residues from the four Ca^{2+} binding loops. We inserted the Ca^{2+} ions into these positions as an initial condition for the next minimization process with the all-atomistic AMBER99SB-ILDN force field to define their final positions. During the minimization, only the backbone heavy atoms (excluding Ca^{2+}) of the holoCaM-Ng_{13–49} were constrained. It was then followed by solvation of TIP3P water molecules and ionization of Na^+/Cl^- ions, which required another round of energy minimization and equilibration. The backbone heavy atoms and the Ca^{2+} ions were constrained (but not fixed) for the holoCaM-Ng₁₃₋₄₉ complexes. After preparation, the positions of the Ca^{2+} ions were robust and did not drift away because 1) there were ~5–7 oxygen atoms from the side chains in the loop within 4 \AA of each Ca²⁺ ion; 2) our calculation showed that the average stability of the Ca²⁺ ions at binding site III and IV is $\Delta G =$ -16.0 and -12.8 kcal/mol, respectively.

Pulling simulations and computation of the free energy differences

We performed all-atomistic steered molecular dynamics (MD) simulations (see Sections IV.3 and IV.4 in the [Supporting Material](#page-11-0) for the procedures and details of the steered MD simulations) to calculate the Ca^{2+} binding free energy (ΔG) using JE [\(34](#page-12-0)). JE relates ΔG and the irreversible work along an ensemble of non-equilibrium trajectories pulling the Ca^{2+} from the bound state to the unbound state (Eq.uation S10). The nonequilibrium trajectories (number $M = 100~150$) were generated by steered MD simulations with all-atomistic models. The difference in the Ca^{2+} binding free energy ΔG between holoCaM-CaMBT (CaMBT: calmodulin binding targets) and CaM is defined as $\Delta\Delta G = \Delta G^{holoCaM-CaMBT} - \Delta G^{holoCaM}$. $\Delta\Delta G$ allows us to evaluate the influence of different CaMBTs on Ca^{2+} binding to CaM, where $\Delta G^{\text{holoCaM-CaMBT}}$ and $\Delta G^{\text{holoCaM}}$ were estimated from independent simulations of pulling Ca^{2+} from the holoCaM-CaMBT and the holoCaM, respectively.

Analyses: calculation of ''apparent chemical shifts''

In the NMR experimental work by Hoffman et al. [\(30](#page-12-0)), a semiquantitative comparison of amide chemical shifts were presented for residues in apoCaM (calmodulin without calcium ions) that are affected by binding of Ng13–49. In the NMR spectroscopy, the chemical shift is the relative resonant frequency in the local induced magnetic field and is diagnostic of the molecular structure. A threshold was applied to generate a set of Boolean numbers, with "1" representing a significant change in the chemical shifts and "0" indicating no change. We refer to the Boolean series as "apparent chemical shifts'' throughout this study. Berjanskii and Wishart ([42\)](#page-12-0) showed that the inverse absolute chemical shifts roughly correlates with root mean square fluctuations (RMSF), which reflects the flexibility of the structure. Therefore, to compare with the change of chemical shifts from unbound to bound apoCaM in the experiment, we approximated the amide chemical shifts by the inverse of RMSF of C_{α} beads and defined "apparent chemical shifts" as $\Delta = \Theta$ (RMSF^{unbound}/RMSF^{bound} - τ). $\tau = 50\%$ was used. Δ measures the change of local flexibility of the apoCaM: larger values of Δ indicate stabilization of the corresponding residue upon binding with Ng_{13-49} .

RESULTS

Building weakly bound apoCaM-Ng13–49 complexes in coarse-grained models

Because of the lack of experimentally determined apoCaM- Ng_{13-49} complex structures, we used low-resolution information on the bound complex of apoCaM-Ng_{13–49} from NMR measurements ([30\)](#page-12-0) for building plausible models from coarse-grained molecular simulations. The NMR experiments showed the change in the backbone chemical shifts of apoCaM upon Ng_{13-49} binding [\(Fig. 2](#page-3-0) *a*). It revealed the residual information on the structural changes

FIGURE 2 Reweighted PMF of apoCaM-Ng₁₃₋₄₉ binding and changes in chemical shifts for apoCaM upon Ng₁₃₋₄₉ binding. (a and b) The residues that experienced significant changes in chemical shifts are projected on the apoCaM structure (PDB: 1CFD) as a ball-and-stick representation from experimental measurement in (a) [\(30](#page-12-0)) and from our calculations according to the coarse-grained molecular simulations in (b). (c) The PMF from the coarse-grained molecular simulations is plotted against the center of mass distance between apoCaM and Ng_{13-49} (d_{COM}) and the number of intermolecular contacts (Z). The coarse-grained molecular simulations were performed at $pH = 6.3$ and ionic strength $= 0.1$ M according to the conditions of the nuclear magnetic resonance experiments ([30\)](#page-12-0). We used a reduced unit of length $\sigma = 3.8\text{Å}$. The color is scaled in k_BT, T = 1.1 $\varepsilon/\text{k}_\text{B}$, where ε is reduced unit of energy and k^B is Boltzmann constant. For visual guidance, we provided bars above the schematic representation of secondary structures in red and blue segments of apoCaM on the half circles in $(a \text{ and } b)$. To see this figure in color, go online.

of apoCaM upon binding with Ng_{13-49} [\(30](#page-12-0)), but did not guarantee a direct contact between apoCaM and Ng_{13-49} . The experiment showed that the major structural changes occur in cCaM as well as Helices B and C in nCaM, as in Fig. 2 a.

One of the prominent reasons that we used coarse-grained simulations is to achieve efficient sampling of the conformational space (see Sections I–III in the [Supporting Material](#page-11-0) for details about the coarse-grained simulations). We implemented the statistical potentials into the structural Hamiltonian of the Ng_{13-49} , as well as the intermolecular interaction between apoCaM and Ng_{13-49} . For the modeling of apoCaM, we kept limited structural information that favors an extended conformation and matches the structural description of the experimental measurements $(25,43)$. We found that apoCaM-Ng13–49 complex samples an ensemble of varying conformations through nonspecific intermolecular interactions, indicated by the broad distribution of the bound complex in the potential of mean force (PMF) (see Fig. $2 c$).

Once we generated the ensemble of bound complexes, we strategically selected the most probable complex structures of apoCaM-Ng_{13–49} with the previously mentioned NMR measurement as a guide (Fig. 2 a) by following the procedures: 1) we reweighed the ensemble of apoCaM-Ng₁₃₋₄₉ structures from the umbrella sampling simulations [\(44](#page-13-0)) using Weighted Histogram Analysis Method [\(45,46\)](#page-13-0), and employed an importance sampling method [\(32](#page-12-0)) (see Section V.4 in the [Supporting Material](#page-11-0)) to select a subset of 23,722 structures from five million independently sampled conformations. The surprisal value between the sample distribution $P^{\text{sample}}(d_{\text{COM}})$ and the original unbiased distribution $P^{\text{ori}}(d_{\text{COM}}) = 0.14$ to show that the sample distribution represents the original distribution well. 2) We performed a clustering analysis on the sampled structures (see Section V.5 in the [Supporting Material](#page-11-0) for the clustering method); 3) we computed the ''apparent chemical shifts'' (see [Analyses: calculation of ''apparent chemical](#page-2-0) shifts" for the definition, Fig. $2 b$) from the major clusters. The population of the major four clusters (comprising 86%, Table S2) and the correlation coefficients between computationally obtained ''apparent chemical shifts'' from the four clusters and the experimental NMR data are shown in Table S2. We identified the cluster (the dominant cluster) of complex structures, which best matched the experiment ([30\)](#page-12-0), with a correlation coefficient of 0.36. From the dominant cluster (i.e., cluster 1 in Table S2), cCaM is mostly stabilized by binding with Ng_{13-49} and helices B and C are partially stabilized (Fig. 2 c).

We next analyzed the structures of the weakly bound apoCaM-Ng13–49 by comparing the probability of contact formation from the unbound state (defined by $d_{COM} \approx 20$ σ in [Fig. 2](#page-3-0) c) in Fig. S3. In the selected cluster (cluster 1) of the apoCaM-Ng13–49 complexes, contacts within apoCaM (triangles in solid lines in Fig. S3 c) became less probable upon Ng_{13-49} binding. In addition, the probability of interdomain contacts between Helix A from nCaM and inter-domain linker, helices G and H from cCaM (the solid rectangles in Fig. $S3 c$) decreased to allow interaction with Ng_{13-49} (the solid rectangles in Fig. S3 *a*). In contrast, after binding with apoCaM, Ng_{13-49} diminished interactions between IQ motif and the acidic region (Fig. $S3 b$, the rectangle in dotted lines); subsequently, probability of forming α -helix structure in the IQ motif increases (Fig. S3 b, the rectangle in solid line). Provided that there was no bias to the formation of α -helix structure in Ng_{13–49}, our model predicted formation of secondary structures in the IQ motif after binding with apoCaM.

For the intermolecular interaction of the weakly bound complex from the main cluster, IQ motif of Ng_{13-49} interacted with Helix A, B/C Helix linker, and Helix D from nCaM, and interdomain linker, Helix E and Helix H, and the F/G Helix linker from cCaM (Fig. S3 a , ellipse in dotted lines). These regions from CaM (either apoCaM or holoCaM) also participate in canonical binding of holoCaM with CaMKI or CaMKII [\(25](#page-12-0)). Although the IQ motif of Ng13–49 formed a helical structure and interacts with cCaM of high probability, the acidic region of Ng_{13-49} remains unstructured (Fig. S3 b , no change in contact formation along diagonal for the acidic region). It is highly probable for the acidic region to interact with both nCaM and cCaM, especially with F/G helix linker and the central linker (Fig. $S3$ *a*, ellipse in solid lines). We further calculated the correlation among these pairwise contacts in Section V.6 of the [Supporting Material](#page-11-0), and found that the interaction between acidic region of Ng_{13–49} and Ca²⁺ binding loops from apoCaM anticorrelated with those within cCaM (shown in magenta ellipse in Fig. S4).

Molecular mechanism of apoCaM-Ng₁₃₋₄₉ binding

We next investigated the conformational changes of apoCaM-Ng13–49 during the association process along their center-of-mass separation (d_{COM}) for better understanding of the molecular mechanism of binding. The unbound structure was taken from $d_{COM} = 20.0 \sigma$, where the interaction between the two proteins is negligible; the encounter state structure ([25\)](#page-12-0) was taken from $d_{COM} = 10.0 \sigma$, where the two proteins began to make contact; the bound state structure was take from $d_{COM} = 2.8 \sigma$, where the PMF(d_{COM}) reached the minimum (see the magenta curve Fig. S2 for the PMF). Superposed structures of apoCaM-Ng13–49 complex at $d_{COM} = 20.0$, 10.0, and 2.8 σ are shown, representing unbound, encounter, and bound states, respectively (Fig. 3).

FIGURE 3 Illustration of the structural changes in unbound, encounter, and bound ensemble of CaM-Ng13–49 complexes. Structures of CaM in (a) was taken from the unbound state when CaM and Ng_{13-49} are well separated at $d_{COM} = 20.0 \sigma$. $\sigma = 3.8 \text{ Å}$. (b) Structures from the encounter of binding when apoCaM and Ng is separated at $d_{COM} = 10.0 \sigma$. (c) Structures from the bound state at $d_{COM} = 2.8 \sigma$. For visual guidance, we superposed 20 sets of structures in each panel. The CaM is colored in black; the residues (residues 99~101 in Ca²⁺ binding loop III and residues 135~137 in Ca²⁺ binding loop IV), which form EF-hand β -scaffold in cCaM are colored in blue; the acidic region and IQ motif of Ng_{13-49} are colored in pink and yellow, respectively. To see this figure in color, go online.

There was no apparent structural change in apoCaM from the unbound state (Fig. 3 a) to the encounter complex (Fig. 3 b) that the conformation of apoCaM remains extended. The onset of binding in Fig. $3 b$ shows that the IQ motif of the Ng_{13-49} mainly interacted with cCaM, which was not involved in a global conformational change of apoCaM. From the encounter complex $(Fig. 3 b)$ to the bound state (Fig. 3 c), Ng_{13-49} remained partially structured and mainly gained interactions with nCaM and local conformational change of apoCaM was observed. In particular, the EF-hand β -scaffold [\(47](#page-13-0)), which stabilizes the two helixloop-helix motifs (EF-hand) in cCaM (see Fig. S5 for illustration of the EF-hand β -scaffold), is broken by insertion of the acidic region of Ng_{13-49} .

To understand their binding at the residual level, we plotted the probability of contact formation at d_{COM} = 20.0, 10.0, and 2.8 σ in Fig. S6. At d_{COM} = 20.0 σ , or the unbound state, there was no contact formed between apoCaM and Ng_{13-49} (in Fig. S6 *a*). The contacts within Ng13–49 were ubiquitous and of low probability, corresponding to the transient nature of conformations of an intrinsically disordered protein/peptide (IDP). Upon initial binding with apoCaM at $d_{COM} = 10.0 \sigma$, only IQ motif from the Ng_{13–49} gained contacts with Helix F and Ca^{2+} binding loop IV from cCaM. This indicates a crucial role of the IQ motif in the recognition by apoCaM. Those contacts were kept when d_{COM} became 2.8 σ , which suggests that the IQ motif also facilitated the stabilization of the apoCaM- Ng_{13-49} complex. At this stage, the contacts between the acidic region of Ng_{13-49} and both of nCaM and cCaM formed broadly but with low probability. For apoCaM, from unbound to bound states, most of the contact formation within apoCaM remained the same (of high probability), except when $d_{COM} = 2.8 \sigma$, the side chain-side chain contacts between Helix E and Helix H, between Helix F and Ca^{2+} binding site IV, and the β -sheet contacts between Ca^{2+} binding sites III and IV (encircled in Fig. S6). In particular, the probability of the contacts that form the EF-hand β -scaffold (encircled) became less when the acidic region of Ng_{13-49} formed contacts with cCaM. Ng became partially helical in the bound form. We will relate this finding to its importance of specific functions under Discussion.

Binding of the CaMBTs to CaM modulates the $Ca²⁺$ binding affinity validated by atomistic pulling simulations

We found that interactions between the acidic region of Ng_{13–49} and the Ca²⁺ binding loops from cCaM-apoCaM competed with those within apoCaM. We hypothesized that, due to this competition, the EF-hand β -scaffold, which is shown to control the opening and closing of the EF-hands ([47,48](#page-13-0)), become less stable, thus facilitating the release of Ca^{2+} from CaM. To test this hypothesis, we strategically selected apoCaM-Ng₁₃₋₄₉ complex structures from the most dominant cluster (i.e., cluster 1 in Table S2) in the coarse-grained molecular simulations for the calculation of Ca^{2+} binding free energy. The selection was based on the characteristics of apoCaM binding with either Ng protein or Ng peptides according to the results from several experimental studies: 1) EF-hands in cCaM are open and EF-hands in nCaM are closed from X-ray crystallography and NMR experiments [\(29,49\)](#page-12-0); 2) Ng_{13-49} has more interactions with cCaM as suggested from NMR studies, $\Delta Z = Z_n - Z_c < 0$ ([30\)](#page-12-0). Z_n (Z_c) is the number of intermolecular contacts between Ng and nCaM (cCaM). (See Section V in the [Supporting Material](#page-11-0) for the definition of Z.) This was further shown by the EF-hand angles in CaM in forms of Ca^{2+} -absent, Ca^{2+} -loaded, apoCaM-NgIQ, and holo-CaM-CaMKII in Table S3. As a result, four coarse-grained structures were selected from the major cluster (cluster 1) for further reconstruction of all-atomistic protein models with four Ca^{2+} ions included (see ([50\)](#page-13-0) for the protocol of reconstructing the all-atomistic model from a side-chain- C_{α} model).

We evaluated the free energy difference between Ca^{2+} -unbound and Ca^{2+} -bound state ($\Delta G = G_B - G_U$,

see Fig. 4 for illustration) using JE $(34,51)$ by pulling the two Ca²⁺ ions independently from the Ca²⁺ binding sites III and IV of the cCaM (see simulations details in IV.5 from the [Supporting Material\)](#page-11-0). We further computed the free energy changes in the absence or the presence of a target ($\Delta \overline{\Delta G} = \Delta G^{holocam-TaMBT} - \Delta G^{holocam}$, see Fig. 4 for illustration). $\Delta\Delta G$ allows us to evaluate the influence of CaMBT on Ca²⁺ binding to CaM. If $\Delta\Delta G > 0$, it means that the CaMBT destabilizes the bound state and thus decreases the Ca²⁺ affinity, as illustrated in Fig. 4. If $\Delta\Delta G$ <0, CaMBT enhances Ca^{2+} affinity for CaM.

We found that the distribution of work values does not follow Gaussian distribution (Fig. S14). The estimation of free energy difference ΔG could be very inaccurate by directly applying JE or its second-order cumulants approximation ([51\)](#page-13-0) (we reported the direct estimation from JE in Table S6). As shown in Fig. S13, the binding free energy ΔG was estimated for Ca²⁺ at binding site III or site IV of holoCaM and holoCaM-CaMKII using a running JE estimate or a block-average method (BA, i.e., using subsets of all available work data using Jarzynski's equality). ΔG converged to values, which deviate from experimental values $(-5 \text{ to } -4 \text{ kcal/mol})$. To improve the efficiency of the free energy estimation, we implemented the cumulative integral (CI) extrapolation method developed by Ytreberg and Zuckerman ([51\)](#page-13-0). CI uses an integral for more accurate estimations than the linear extrapolation method by extrapolating to $1/n \rightarrow 0$, where n is total number of work values. It was shown that CI extrapolation could reduce the required data by 5- to 40-fold $(51,52)$ $(51,52)$ $(51,52)$. Therefore, we adopted the CI extrapolation method to estimate the free energy difference.

We summarized the binding free energies of Ca^{2+} estimated by the CI extrapolation in [Table 1](#page-6-0) for a total of six

FIGURE 4 Illustration of the definitions of the binding free energy ΔG and $\Delta\Delta G$. $\Delta G = G_B - G_U$. $\Delta\Delta G = \Delta G^{holoCaM-CaMBT} - \Delta G^{holoCaM}$. B and U stand for bound and unbound states of the Ca^{2+} , respectively. $\Delta\Delta G > 0$ means that the CaMBT destabilizes the bound state and thus decreases the Ca^{2+} affinity.

TABLE 1 Difference in Binding Free Energy of Ca^{2+} Calculated from Nonequilibrium Molecular Simulations and from the Experiments at $pH = 7.4$.

		holoCaM- Ng_{13-49}	holoCaM-CaMKII
$\Delta\Delta G^{CI}$ (kcal/mol)	Site III	$9.2 + 2.2$	-2.5
	Site IV	$22.4 + 0.9$	-1.7
$\Delta\Delta G^{exp}$ (kcal/mol)	Site III/IV	2.5	-3.3

The CI extrapolation was used in the calculation of binding free energy of $Ca^{2+}\Delta\Delta G^{CI}$ from the simulations. The averaged value of $\Delta\Delta G$ statistical errors from the four conformations of the holoCaM-Ng13–49 complex are provided. The experimental binding free energy $\Delta\Delta G^{exp}$ for holoCaM- Ng_{13-49} was from literature ([30\)](#page-12-0) and the experimental values of the holo-CaM-CaMKII were from ([14\)](#page-12-0).

conformations: four holoCaM-Ng13-⁴⁹ (showing averaged values), one holoCaM and one holoCaM-CaMKII. For all of the four conformations of holoCaM-Ng₁₃₋₄₉, the binding of Ng₁₃₋₄₉ to CaM destabilized the Ca^{2+} in site III and site IV, respectively, by showing an averaged positive $\Delta\Delta G$ (Table 1). In comparison, we examined the importance of the canonical bound structure for the release of Ca^{2+} from the complex structure of holoCaM-CaMKII (PDB: 1CDM) where the CaM wraps around and binds with a well-formed helical structure of CaMKII peptide at an antiparallel position. From CI extrapolation, each site presented a negative $\Delta\Delta G$, implying that the presence of CaMKII stabilizes the Ca^{2+} in both site III and site IV from CaM.

Conformation of bound CaM-CaMBT complex dictates Ca^{2+} release

HoloCaM-CaMKII forms a canonical bound complex, while apoCaM- Ng_{13-49} does not. In the following, we unveiled the molecular mechanism that governs varying $Ca²⁺$ binding affinity for cCaM in complexes of CaM and distinct CaMBTs. We examined one holoCaM-Ng13-49 model reconstructed from the coarse-grained structure that led to a maximal increase of Ca^{2+} binding free energy, as well as the crystal structure of holoCaM-CaMKII (PDB: 1CLL).

In Fig. 5 a , Ng_{13–49} is bent from the middle between residues P24 and G25 from the acidic regions (in green and in gray on the amino acid sequence). On the one hand, the IQ motif interacts with Helix F, Helix G, and F/G helix linker from cCaM to stabilize the holoCaM-Ng₁₃₋₄₉ complex. On the other hand, the acidic N-term interacts with both nCaM and cCaM. Surprisingly, the four residues D22-D23-P24-G25 (DDPG) (in ball-and-stick representation) before the IQ motif on Ng_{13-49} stick out and insert into the middle of the EF-hand β -scaffold (colored in *blue* segments) between Ca^{2+} binding loops III and IV from cCaM, disrupting the stability of EF-hand β -scaffold. As a result, the distance between the two β -strands becomes 7.9 \pm 1.2 Å (average over the four conformations)

FIGURE 5 Differential intermolecular interactions in holoCaM-Ng₁₃₋₄₉ and holoCaM-CaMKII complex structure tune Ca^{2+} binding affinity for cCaM. (a) The atomistic structure of holoCaM-Ng₁₃₋₄₉ complex was reconstructed from the coarse-grained simulations that led to the largest decrease in Ca^{2+} affinity for site III and site IV. (b) The structure of holoCaM-CaMKII complex, showing a wrap-around binding pattern, is from the X-ray crystallography (PDB: 1CDM). The complete structures of the complexes are on the left, CaM is in black, Ca^{2+} ions are represented by yellow beads, and the CaMBTs are in white. The sequences of the CaMBTs are provided above the figures. For visual guidance, the cCaM and the CaMBTs are enlarged and recolored on the right. The cCaM is in ribbon representation and colored in orange. The residues forming EF-hand β -scaffold in the crystal structure, i.e., residues Y99, I100 from site III, and residues Q135, V136 from site IV, are colored in blue. The invariant Glu residue in the 12th position of each Ca^{2+} -binding loop contributes two oxygen atoms to the coordination of the Ca^{2+} ion and is shown in ball-and-stick representation. The dotted lines show the distance between Ca^{2+} and the two oxygen atoms from Glu12 and the distance between the two β -strands in cCaM. The four residues that stick out from Ng13–49 are D22 (red) D23 (red) P24 (green) G25 (magenta) in (a). The residues from CaM, which form contacts with CaMKII peptides, are colored in light green in (b) . To see this figure in color, go online.

comparing to \sim 5.0 Å in the intact holoCaM (PDB: 1CLL). In contrast, the distance decreases slightly to 4.8 Å in holoCaM-CaMKII (PDB: 1CDM).

Moreover, the average distance between the Ca^{2+} and the Ca^{2+} coordinating oxygen atoms from the 12th residue (Glu-12), or $\overline{d}^{\text{Glu-Ca}}$, in site III and IV from cCaM is 4.5 \pm 0.7 and 3.7 \pm 0.5 Å, respectively (Fig. 5 *a*). As a reference, for Ca^{2+} in site III and site IV of the crystal structure of holoCaM without a target (PDB: 1CLL), $\overline{d}^{\text{Glu-Ca}} = 2.3$ and

2.5 Å, respectively. The 12th residue in the Ca^{2+} binding loop, or Glu-12, plays a crucial role in stabilizing Ca^{2+} in the binding loop $(1,47)$ $(1,47)$ $(1,47)$. However, the presence of the acidic region from Ng_{13-49} increases \overline{d}^{Glu-Ca} in the holoCaM- Ng_{13-49} complex [\(Fig. 5](#page-6-0) *a*), especially for site IV, which is slightly closer to the center of mass of the DDPG motif than site III. The motif DDPG from the acidic region of Ng₁₃₋₄₉ pries open the EF-hand β -scaffold and facilitates the release of the Ca^{2+} from cCaM. In comparison, we examined the importance of the canonical bound structure for the release of Ca^{2+} from the crystal structure of holoCaM-CaMKII (PDB: 1CDM) in [Fig. 5](#page-6-0) b. In contrast, $\overline{d}^{\text{Glu-Ca}}$ for sites III and IV [\(Fig. 5](#page-6-0) b) are similar to those in the crystal structure of holoCaM without a target (PDB: 1CLL).

To understand the evolutionary role of the acidic region of Ng_{13-49} and its biological significance in modulating Ca^{2+} binding affinity of CaM, we performed sequence alignment of Ng from different species available in the protein knowledgebase (UniProtKB, [http://www.uniprot.org/](http://www.uniprot.org/uniprot/) [uniprot/](http://www.uniprot.org/uniprot/)) (Table S14). Sequence alignment results of Ng revealed that the amino acids in the acidic region are either conserved across the species or found as amino acids with similar physicochemical properties (see Fig. S15). In particular, we noted that the cluster of residues ''DIPLDDP'' from the acidic region of Ng that show anti-correlation with the binding loops III and IV of CaM based on the pairwise contacts (see Fig. S4; Table S8) are conserved across all the species (see Fig. S15). Moreover, the ''DDPG'' motif from the acidic region of Ng is conserved (see Fig. S15) across the higher vertebrate species (from rodent to human; see Table S14). Interestingly, in the lower vertebrate species the Gly (G) residue of this motif is found as Asp (D) (in reptiles and birds) or Glu (E) (in fish) or Ala (A) (in zebrafish), whereas the rest of the residues (DDP) are conserved across all the vertebrates (see Fig. S15). We also found that the DDPG motif, which maintains close proximity to the middle part of both the Ca^{2+} binding loops III and IV in the holoCaM-Ng_{13–49} complex (see Fig. S16) plays an important role in disrupting the EF-hand β -scaffold in CaM as we mentioned before. We speculate that the acidic region of Ng and especially the DDPG motif is evolutionarily important for their biological role in suppressing and/or fine-tuning the Ca^{2+} binding affinity of CaM.

DISCUSSION

The energy landscape of apoCaM- Ng_{13-49} binding is rugged

The weakly bound complexes are commonly observed in protein-protein associations involving intrinsic disordered proteins (IDPs) in both experiments [\(53,54\)](#page-13-0) and computer simulations [\(55–60](#page-13-0)). Long-range electrostatic interactions enhance the rate of association of several weakly bound complexes [\(61](#page-13-0)). Generally, the nonspecific interactions, including hydrophobic and electrostatic interactions, are known to be a major contributor to the weakly bound complexes involving IDPs. Sugase et al. ([54](#page-13-0)) showed from NMR experiments that the weakly bound transient complexes of the phosphorylated kinase inducible domain (pKID) of transcription factor CREB and the kinase-induced domain of the CREB-binding protein were predominantly stabilized by nonspecific hydrophobic contacts instead of electrostatic interactions. The importance of nonspecific intermolecular interactions was further addressed by computer simulations on phosphorylated kinase inducible domain and kinaseinduced domain by Turjanski et al. [\(55](#page-13-0)), as well as on nuclear-receptor coactivator binding domain of CREB-binding protein and the p160 steroid receptor coactivator ACTR by Chen et al. ([57\)](#page-13-0). In both cases structure- or topology- based models ([62](#page-13-0)) were used for the intermolecular interactions that define a specific bound structure.

Determination of the structure of the apoCaM-Ng_{13–49} complex experimentally is challenging due to its low stability [\(28](#page-12-0)). To overcome this problem, we incorporated statistical nonbonded interactions and electrostatic interactions, which do not rely on a specific complex structure, in our protein model. In addition, to model the disordered Ng_{13-49} , we implemented a sequence-based statistical dihedral angle potential ([31\)](#page-12-0) in the Hamiltonian. Our study reveals a rugged free energy landscape of apoCaM-Ng13–49 binding ([Fig. 2](#page-3-0) c). The conformations of apoCaM-Ng₁₃₋₄₉ thermally fluctuate between shallow basins separated by 1~2 k_BT (T = 1.1 ε/k_B, where ε is reduced unit of energy and k^B is Boltzmann constant.).

The weakly bound complexes comprises ''fuzzy'' structures, in which either one of the binding partners or both stay dynamic [\(63](#page-13-0)). According to our simulations of the apoCaM-Ng_{13–49} complex, Ng_{13–49} remains partially structured as shown in the low probability of contact formation within Ng_{13–49} (Fig. S3 b and S6 c). Our model of the CaM model ([25\)](#page-12-0) samples a wide range of states by incorporating non-native interactions between residues. This idea of modeling beyond a simple structure-based model meets the experimental exploration of CaM with single-molecule experiments that its energy landscape is rather rugged and complex with plenty of intermediates [\(19\)](#page-12-0). Our approach differs from the approaches adopted by those studies that the conformational changes of the target were ignored. For example, in the investigation of the interactions between the transcriptional activation domain of an oncoprotein and its target by Chan et al. ([64](#page-13-0)), the target protein was modeled as a sphere of distributed charges on the surface.

To generalize, the structural flexibility and nonspecific interactions, corresponding to entropic and enthalpic effects, respectively, result in the rugged energy landscape of apoCaM-Ng_{13–49} that presents multiple weakly bound complexes without a global minimum.

IQ motif of Ng_{13-49} is crucial for the initial binding of apoCaM

It was postulated that the residual structure of an IDP is necessary for molecular recognition in CaM-CaMBT complex formation ([65\)](#page-13-0). We found that the residual structure in Ng is crucial for recognition and binding with CaM. We found that the IQ motif started to interact with cCaM at a distance of $d_{COM} = 10 \sigma$ and remained during the binding process, indicating its key role in binding with apoCaM. Creamer et al. ([66\)](#page-13-0) using CaM-CaMBT systems as well as Clarke et al. [\(67](#page-13-0)) using the BH3 motif of the largely disordered protein p53 upregulated modulator of apoptosis (PUMA) and a folded protein induced myeloid leukemia cell differentiation protein (MCL-1), however, postulated that the residual structure of the IDP did not affect binding kinetics, but weakened the binding affinity upon helixbreaking mutation. It is important to note that in Creamer's study, most of the CaMBTs bind CaM in a canonical form where CaM undergoes large global conformational changes. In such scenarios, the "conformational and mutually induced fit'' mechanism ([25\)](#page-12-0) is needed for the recognition and binding (see The binding mechanism of CaM-CaMBT might be key to regulating Ca^{2+} binding affinity on the binding mechanism of CaM-CaMBT). However, Ng predominantly interacts with cCaM and does not form a full helical structure, therefore the binding does not require global conformational changes of CaM.

The acidic region is key to Ng_{13-49} for tuning CaM's affinity for Ca^{2+} comparing to other IQ-motif peptides

Our computation of the change in free energy of Ca^{2+} binding to cCaM from nonequilibrium atomistic pulling simulations reveals that the target peptides Ng_{13-49} and CaMKII have opposing effects on Ca^{2+} binding affinity to cCaM. It is important to note that the CaMKII peptide binds to CaM in a canonical form in which the two domains of CaM wrap around a helical CaMKII peptide ([24\)](#page-12-0), whereas Ng_{13-49} forms weakly bound complexes with a rather extended CaM. Such a difference results in distinct conformations of the EF-hand motifs and the Ca^{2+} binding loops in CaM.

The binding and dissociation of Ca^{2+} to/from cCaM essentially depend on the two EF-hand motifs, which are connected by the F/G helix linker and are rigid in holoCaM (Fig. S5). This construct of EF-hand motif pair is often referred as the "EF-hand β -scaffold" ([47\)](#page-13-0) proposed by Grabarek ([47\)](#page-13-0). The EF-hand β -scaffold immobilizes the Ca^{2+} ion during its initial binding to the Ca^{2+} binding site. Our findings with regard to target binding that modulate CaM's calcium affinity is supported by experimental observations [\(1](#page-12-0)). An enhanced structural stability of an EF-hand β -scaffold through target binding in a canonical CaM- CaMBT form increases CaM's calcium affinity, which is discussed in the detailed review by Gifford et al. [\(1](#page-12-0)). In those experiments, a severe disruption in the Ca^{2+} binding loops due to the repacking of cCaM causes the Ca^{2+} coordinating oxygen atoms from the Glu-12 of the Ca^{2+} binding loop to move away from their Ca^{2+} binding positions and the release of the Ca^{2+} [\(68\)](#page-13-0). In addition, several experiments showed that when CaMBT peptides bind to CaM canonically, the Ca^{2+} affinity for cCaM increases ([69\)](#page-13-0): $\Delta\Delta G$ varies from -1.5 kcal/mol for phosphorylase kinase, -3.3 kcal/mol for CaMKII [\(70](#page-13-0)) to -4.8 kcal/mol for β -calcineurin [\(71\)](#page-13-0).

We found from this work that during the binding process, on one hand, IQ motif of Ng_{13-49} carries the function of recognition and stabilization. On the other hand, the role of the acidic region of Ng_{13-49} might relate to the biological functions of the protein because of its influence on the stability of EF-hand β -scaffold (Fig. S6 c) regulating Ca²⁺ binding to CaM. Ng belongs to the IQ motif CaMBTs family. IQ-motif CaMBTs such as the sodium channel (Na_V1.2), the calcium channel ($Ca_V1.1$ and $Ca_V1.2$), PEP-19, and Ng present distinct effects on CaM's affinity for Ca^{2+} .

The presence of the acidic region rather than the IQ motif determines whether the interactions between CaM and the IQ-motif CaMBT are disruptive or constructive. For $Na_V1.2$ and PEP-19 (Purkinje cell protein 4; PCP 4) that contain an IQ motif and a previous acid region, both of them bind to cCaM. Binding of them lowers the Ca^{2+} binding affinity for cCaM: binding of the $\text{Na}_{\text{V}}1.2$ causes an increase in $\Delta\Delta G$ of about +2 kcal/mol [\(72](#page-13-0)) and PEP-19 an increase of about $+0.2$ kcal/mol [\(49](#page-13-0)). However, for $Ca_V1.1$ and $Ca_V1.2$, which contain an IQ motif but no such previous acidic region, they behave like a canonical binding motif, where the CaMBT forms a helical structure wrapped by the compact CaM, and stabilize the Ca^{2+} binding by $\Delta\Delta G = -1$ and -2.6 kcal/mol, respectively ([73\)](#page-13-0). The results support the active role of the acidic region, but do not imply the molecular mechanism of how it regulates Ca^{2+} binding to CaM. By carrying out multiscale molecular simulations, we identified the key residues DDPG in the acidic region of Ng_{13–49} that "pry" open the β -sheet structure between the two Ca^{2+} binding loops.

The binding mechanism of CaM-CaMBT might be key to regulating Ca^{2+} binding affinity

A number of recent advancements have been made in experimental and computational approaches to examine proteinprotein interactions involving IDPs [\(74–78](#page-13-0)). A key emerging concept is the idea of ''specificity on demand'' where protein interactions are dictated by concerted interactions between both binding partners. From this view, the process of binding cannot be inferred even if the structures of the initial reactants and final complex are known. The importance of structural flexibility motivated the develop-ment of theories for binding mechanisms ([79–83](#page-13-0)), such as fly casting $(61,84)$ $(61,84)$, folding upon binding $(58,85)$ $(58,85)$ $(58,85)$, "dockand-coalesce'' ([74,86](#page-13-0)), and conformational selection ([87,88](#page-13-0)). These multiple steps mechanisms underlie possible kinetic bottlenecks due to subsequent conformational changes in a post-collision event $(74,86)$ $(74,86)$. Despite differences in these theories, past efforts that focus on IDPs binding to known structures have not dealt with the complexities of the conformational variability inherent in the binding process that can ultimately lead to different final products involving potentially large changes in both partners. This view is highlighted by our efforts that revealed that CaM binding to its targets is mediated by a conformational and mutually induced fit model [\(25](#page-12-0)).

In our previous study (25) (25) , we discovered that the binding between CaM and CaMBTs that forms a canonical bound complex, such as the CaMKII peptide, require a ''conformational and mutually induced fit'' mechanism. In this mechanism, the binding of the CaMKII peptide and CaM presents a two-stage process: at the onset of binding (as in an early stage), the binding is driven by diffusion and electrostatic steering effect that both CaM and CaMKII do not undergo large conformational changes. At the late stage, both experience large conformational changes that the two domains of CaM wrap around the emerging rod-like CaMKII peptide. The latter event of binding dictates a small, but significant, difference in the association rates of CaM with various CaMBTs. However, this is not the case for apoCaM-Ng13–49 binding. After the onset of binding where, the IQ motif of Ng_{13-49} binds to cCaM, there is no subsequent large conformational change of CaM.

The difference in the bound complexes of CaMBT-CaM along its progressive binding mechanism delineates their effects on Ca^{2+} binding affinity. The Ca^{2+} affinity is increased by binding of CaMKII peptide. The EF-hand motifs of CaM wrap around the rod-like CaMKII peptide; hence, they enhances the stability of the "EF-hand β -scaffold" that retains Ca^{2+} in Ca^{2+} binding loops. In contrast, the Ca^{2+} affinity is decreased by binding of Ng peptide because of lack of such wrap-and-enhance pattern of CaM. In addition, insertion of the DDPG motif into the "EF-hand β -scaffold" further destabilizes the Ca^{2+} binding loops. Therefore, we speculate that the benefit of a progressive binding mechanism of CaM that can be modulated by CaMBTs underlies their distinct effects on CaM's Ca^{2+} binding affinities.

Force fields affect the magnitude of free energy estimation

From [Table 1,](#page-6-0) the signs of $\Delta\Delta G$ from our computed values agreed with the experimental measurements. For holoCaM-CaMKII, they were in the same order of magnitude, while for the holoCaM-Ng_{13–49} complexes, the computed values were about one order of magnitude larger than the experimental measured ones. Majorly the force field contributed to the accuracy in the calculation of $\Delta\Delta G$. We break the reasoning down into three categories relating to force fields:

Initial structures might matter

We found that the current all-atomistic force fields, including AMBER99SB-ILDN and CHARMM27, alter the initial structures after minimization and equilibration if the positions of the heavy atoms were not strictly fixed in space. This is true for coordinates from the crystal structures of holoCaM and holoCaM-CaMKII, as well as the allatomistic structures of holoCaM-Ng13–49 reconstructed from coarse-grained models. The shifts in the positions are especially prominent for the Ca^{2+} ions. According to the study by Shukla et al. [\(89](#page-14-0)) on the same structure of holoCaM (PDB: 1CLL) using the AMBER99SB-ILDN force field, the root mean-square deviation of the equilibrated structure from the crystal structure was around 1 A. This is in agreement with our simulations using the same force field that after minimization and equilibration without fixing the heavy atoms, the displacement between the Ca^{2+} atom and the center of mass of the corresponding Ca^{2+} binding loop (d^{Ca}) changed by $\delta d^{Ca} \sim 0.5$ Å for holoCaM, $\delta d^{Ca} \sim 1.0$ Å for holoCaM-CaMKII, and $\delta d^{Ca} \sim 1.0-2.1$ Å for holoCaM-Ng₁₃₋₄₉. This shift in d^{Ca} affected the free energy of any Ca^{2+} bound state. For a crude estimation of the free energy differences due to a shift in d^{Ca} , suppose the average force is 500 kJ/mol/nm, an underestimation of $\Delta G^{\text{holoCaM}}$ would be by 6 kcal/mol if $\delta d^{Ca} = 0.5$ A^{\dot{A}} . Likewise, an underestimate of $\Delta G^{\text{holoCam-CaMKII}}$ would be by 12 kcal/mol if $\delta d^{Ca} = 1$ Å. Such deficiency would cause a net underestimation of $\Delta\Delta G^{\text{holoCaM-CaMKII}}$ by 6 kcal/mol.

Therefore, we adopted the following strategy in our estimation of $\Delta\Delta G$: for a system that has experimentally determined structure, such as holoCaM or holoCaM-CaMKII, we fixed the positions of the Ca^{2+} atoms and the backbone heavy atoms during the energy minimization and equilibration simulations; otherwise, for holoCaM-Ng₁₃₋₄₉ which has no experimentally determined structure, we reconstructed the all-atomistic coordinates from coarse-grained models of apoCaM-Ng_{13–49} and constrained the positions of the Ca^{2+} ions and the backbone heavy atoms.

We determined the initial positions of the Ca^{2+} ions carefully. We estimated the position of each Ca^{2+} ion as the center of mass of the side chains of the corresponding Ca^{2+} coordinating residues and minimized the potential energy with the all-atomistic AMBER99SB-ILDN force field for optimization. During the energy minimization, only the backbone heavy atoms (excluding Ca^{2+}) of holoCaM-Ng_{13–49} were constrained. Afterward, 1) there were \sim 5–7 oxygen atoms in the loop within 4 \AA of each Ca²⁺ ion; 2) later calculation showed that the Ca^{2+} ions at binding site III and IV had average $\Delta G = -16.0$ and -12.8 kcal/mol, respectively. In summary, the final positions of the Ca^{2+} ions of holoCaM-Ng_{13–49} in the initial structures were robustly determined. Our results with respect to the sign of the differences in the free energy calculation whether targets are present or not were not affected.

The protonation states of the proteins might matter

ApoCaM is a highly charged protein with a net charge of $-24e$, where *e* is the elementary charge. The IDP target peptide Ng_{13-49} contains 12 charged residues. Therefore, accurate determination of the partial charges according to the protonation states particularly at the Ca^{2+} -binding loops impacts the estimation of the binding free energy of Ca^{2+} for CaM. The standard procedure of assigning protonation states according to a neural pH and its pKa eludes the correct prediction of the actual protonation states for Glu and Asp amino-acids in the Ca^{2+} -binding loops of CaM. Therefore, it requires a careful treatment on the protonation states. We employed the H_{++} program [\(41](#page-12-0)) to predict the protonation states based on pH and electrostatic energies that vary with the protein conformation and ionic strength of the solution.

After determining the protonation states of the proteins (complexes) at $pH = 7.4$ and ionic strength of $= 0.15$ M using H_{++} server [\(41](#page-12-0)), we found that there was a significant change in the net charges in comparison with the outcome from the standard protonation protocol of the MD package GROMACS [\(90](#page-14-0)). In the standard protocol, the Asp and the Glu amino acids in the Ca^{2+} -binding are deprotonated at neutral pH and they carry negative charges. However, from the predictions by $H++$, they are protonated instead. For holoCaM, the newly protonated residues by H_{++} included D20, D56, E67, D93, and D129 from the Ca^{2+} binding loops I–IV. For holoCaM-CaMKII, the newly protonated residues by $H++$ included D20, D56, E67, D93, D133, and E140. For holoCaM-Ng₁₃₋₄₉, the results depended on the conformations: most of the protonated residues by H_{++} included D22, E67, and D93. Because the acetyl groups $(-COO^{-})$ from the Ca^{2+} binding loops stabilize the Ca^{2+} -bound state, the numerical estimate of free energies was way off from the experimental values. Once the side chains of these residues in the Ca^{2+} binding loop were adequately protonated, the order of magnitude in Ca^{2+} binding free energy from computer simulations became comparable to the experimentally measured ones (shown in [Table 1\)](#page-6-0).

As a comparative study, we explored the robustness of our results by changing the pH from 7.4 to 6.8. According to H_{++} , the protonation states of holoCaM remained the same at the range of pH $\in [6.9-7.8]$. However, at $pH \in [6.5–6.8]$, E104 at Ca²⁺ binding loop III was additionally protonated. This changed the net charge of the loop III as well as the charge of the corresponding Ca^{2+} ion. We found that the Ca²⁺ binding free energy of loop III (ΔG) at $pH = 6.8$ increased by 3.15 kcal/mol comparing to that at $pH = 7.4$. ΔG for holoCaM-CaMKII remains the same at $pH = 6.8$ as the protonation states did not change. Therefore, $\Delta\Delta G$ for holoCaM-CaMKII still kept the same negative sign at $pH = 6.8$ as at $pH = 7.4$. For holoCaM-Ng13-49, depending on different structures, the gain of additional protonated residues at $pH = 6.8$ varied: residues D58, D11 (from Ng) for the first structure, H107 for the second structure, E84 for the third structure, and D95, H107 for the fourth structure. We found that $\Delta\Delta G$ for holoCaM- Ng_{13-49} remained the same positive sign at $pH = 6.8$ as at $pH = 7.4$. Our results remain robust and can be further tested by experiments.

We provided the calculated $\Delta\Delta G$ values of Ca^{2+} binding site III for both holoCaM-CaMKII and holoCaM-Ng₁₃₋₄₉ in Table S7. When the pH was lowered to 6.8, the $\Delta\Delta G$ values decreased for holoCaM-CaMKII (favoring retaining of the Ca^{2+}) and increased for holoCaM-Ng₁₃₋₄₉, (favoring release of Ca^{2+}) as illustrated in Fig. S17. The range of pH we explored here is typical for CaM to carry its specific function in a cell. We concluded that our finding is consistently robust with regard to how CaM encodes calcium incidents by binding to distinctive targets.

The charges of the Ca^{2+} ions depend on the environment

It has been shown that the charge on Ca^{2+} can be $\lt +2e$ due to charge transfer and polarization in the solution and especially in the Ca^{2+} -bound state [\(35–37](#page-12-0)). We computed the charges of calcium ions according to the protonation states of CaM by employing a semiempirical quantum chemistry program MOPAC ([40\)](#page-12-0). Similar to our previous work ([39\)](#page-12-0), we found that the charges on Ca^{2+} ions in holoCaM were approximately in the range of $+1.2e$ to $+1.6e$. However, we did not rescale the size of the calcium ions in this study, which can be critical to the development of proper effective polarizable force fields for Ca^{2+} according to Jungwirth et al. ([36,38](#page-12-0)).

Despite those issues, our calculated free energy successfully matched the signs of the experimental values, whereas they differed by an order of magnitude for the holoCaM- Ng_{13-49} complexes. By using classical MD simulations, we were able to capture that Ng_{13-49} greatly reduces affinity of Ca^{2+} for cCaM whereas CaMKII peptide increases the affinity. More importantly, we provided a molecular mechanism of how the CaMBTs tune CaM's capacity of sensing Ca^{2+} .

The systematic error from the parameters used in steered MD simulations was minimized for the estimation of free energy changes

Computationally, free energy difference can be estimated from equilibrium ensemble using methods including umbrella sampling (44) (44) , free energy perturbation (91) (91) , and thermodynamic integration ([92\)](#page-14-0), as well as from nonequilibrium ensemble by using JE based on nonequilibrium work ([51\)](#page-13-0). We employed the nonequilibrium work method by carrying out many steered MD simulations for its simplicity and a particular reaction coordinate is not required. The drawback of this method is that the accuracy is highly dependent on sufficient sampling of the small work [\(34](#page-12-0)). Although theoretically JE does not depend on the pathways of pulling when removing Ca^{2+} from its binding site, in practice, pulling along some relevant pathways leads to small work and small variance of the work distribution; while pulling along irrelevant pathways usually yields larger work and larger variance [\(51](#page-13-0)). In the latter case, significantly more trajectories are required for ample sampling. Therefore, in several studies, the pulling was only along one or more predetermined pathways repeatedly [\(52,93–95](#page-13-0)). However, for complex systems, it becomes even more difficult to find the relevant pathways. To solve this problem, we first searched the pathways leading to small work by pulling the Ca^{2+} from each binding site for \sim 100 times along randomly generated directions. Second, we increased the number of simulations by pulling along those specific pathways that generate smallest work until the differences in the free energy converges between the direct JE estimation and the CI extrapolation estimation. We minimized the systematic error from sampling by running the simulations until the extrapolation from CI as well as JE converges. The signs of $\Delta\Delta G$ for both systems we studied agreed with the experimental measured ones.

There are other computational studies on the mechanisms of Ca^{2+} releasing/binding in CaM [\(96,97\)](#page-14-0) arguing that the deformation of protein structures attributes the discrepancies in the free energy calculation by allowing backbones to sample wide conformations in the pulling simulations ([97\)](#page-14-0). Even though the protein conformation can deform during pulling, we argue that this part of the work associated with protein deformation from the course of pulling Ca^{2+} should be part of the free energy calculation according to the concept of the JE: only the initial state and the final state determine the free energy difference.

CONCLUSIONS

CaM coordinates the activation of a family of Ca^{2+} -regulated proteins that are crucial for synaptic plasticity at the molecular and cellular level associated with learning and memory in neurons ([3](#page-12-0)). Ng is a neural-specific postsynaptic IDP, for which Ca^{2+} decreases binding to CaM [\(12](#page-12-0)). Due to the low binding affinity of apoCaM-Ng₁₃₋₄₉ interactions, experimentally it is challenging to determine the bound structures. Our coarse-grained molecular simulations of apoCaM-Ng13-⁴⁹ binding guided by experimental measurements demonstrated the importance of conformational flexibility of both the binding partners in various states of the bound complex. Our study revealed that the acidic region of Ng_{13-49} (the DDPG residues right before the IQ motif; refer to [Fig. 1](#page-1-0) for the sequence of Ng_{13-49}) sticks out into the middle and "pries open" the EF-hand β -scaffold of cCaM. We speculate that this feature of destabilizing Ca^{2+}/CaM at low Ca^{2+} concentration could possibly explain CaM in apo-form in the resting cell. CaMKII on the other hand plays a pivotal role in learning and memory formation for both long-term potentiation and mechanisms for the modulation of synaptic efficacy ([98\)](#page-14-0). Our study demonstrated that in an apoCaM-Ng13-⁴⁹ complex, CaM dominantly remains in an extended conformation. This is in contrast to a canonical bound complex where the two domains of CaM wrap around a rod-like CaMKII peptide that essentially stabilizes the Ca^{2+} binding loops of CaM. The embedded CaMKII peptide by two domains increases the binding affinity of Ca^{2+} to CaM. This proposed mechanism leads to a broader understanding of the reciprocal relation between Ca^{2+}/CaM binding and $CaM/CaMBT$ binding ([30\)](#page-12-0) that involves major conformational changes in both partners. To the best of our knowledge, this is the first computational study providing a detailed description at atomistic level about how binding of CaM with two distinct targets (Ng and CaMKII) influences the release of Ca^{2+} from cCaM as a molecular underpinning of CaM-dependent Ca^{2+} signaling in neurons, which has been investigated for decades ([27,30,99,100](#page-12-0)). Molecular basis for learning and memory formation has aroused attention of the neuroscience community dating back to 1980s [\(101](#page-14-0)). We believe this study allows opportunities to connecting the molecular regulations in atomistic detail to the understanding of cellular process cascade of learning and memory formation.

SUPPORTING MATERIAL

Supporting Material, seventeen figures, and fourteen tables are available at [http://www.biophysj.org/biophysj/supplemental/S0006-3495\(17\)30116-9](http://www.biophysj.org/biophysj/supplemental/S0006-3495(17)30116-9).

AUTHOR CONTRIBUTIONS

M.S.C. and P.Z. designed the research; P.Z., S.T., and H.T. performed the research; P.Z., S.T., and H.T. analyzed data; M.S.C., P.Z., and S.T. wrote the article.

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SUPPORTING CITATIONS

References ([102–124\)](#page-14-0) appear in the Supporting Material.

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Supplemental Information

Opposing Intermolecular Tuning of Ca²⁺ Affinity for Calmodulin by Neurogranin and CaMKII Peptides

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Supporting Text

I. Sample preparation

I.1 Intact Ng

Ng is a 78-residue intrinsically disordered protein with IQ motif [\(1,](#page-77-0) [2\)](#page-77-1) composed of hydrophobic and basic amino acids. Available experimental measurement for residual structure of Ng is for the full-length protein Ng from mouse [\(3\)](#page-77-2), therefore, the sequence is used in the modeling and parameterization of Ng, as shown below,

1 10 20 30 40 50 MDCCTESACSKPDDDILDIPLDDPGANAAAAKIQASFRGHMARKKIKSGE 60 70 CGRKGPGPGGPGGAGGARGGAGGGPSGD 78

Following the work by Hoffman et al. [\(4\)](#page-77-3), we divided the Ng protein into several units, including the acidic N terminal (residues 13-25), the IQ motif (residues 26-49) and the poly-Glycine C terminal (residues 50-78). The underscored sequence stands for the IQ motif which partially retains residual structure (residues 25-42) in the unbound state. To reproduce the residual structure, replica exchange method (REM) [\(5,](#page-77-4) [6\)](#page-77-5) was employed for calculation of average helicity of the residual structure and the backbone model-free order parameter S^2 .

I.2 Ng13-49 peptide

Hoffman et al. showed from calorimetry and NMR experiments that the IQ motif alone does not reproduce similar Ng-mediated affinity of Ca^{2+} for Ca^{2+} -free CaM (apoCaM), or the pattern of intermolecular interaction between Ng and apoCaM [\(4\)](#page-77-3). A combination of the acidic region in the N-terminal and the IQ motif yields the minimum composition of Ng for its function.

Therefore the Ng13-49 peptide is used to study the interaction with apoCaM (PDB ID: 1CFD) that allows the comparison of the results from computer simulations with several experimental measurements, including dissociation constant of apoCaM and Ng_{13-49} , the affinity of Ca^{2+} to CaM, and the changes in the chemical shifts of apoCaM upon Ng₁₃₋₄₉ binding [\(4\)](#page-77-3). The sequence of the Ng13-49 peptide is provided below,

20 30 40 13 DDDILDIPLDDPGANAAAAKIQASFRGHMARKKIKSG 49

II. Coarse-grained protein or peptide models

II. 1 Hamiltonian and parametrization of the coarse-grained protein models

As described above, Ng is partly ordered in solution. The fragment G25-A42 has a fraction of 22% and 28% residual structure calculated from the C_{α} and the H_α chemical shifts in the nuclear magnetic resonance (NMR) experiment [\(3\)](#page-77-2), respectively. The rest part of the peptide remains unstructured. In order to compare with results from several experiments $(3, 4)$ $(3, 4)$, peptide Ng₁₃₋₄₉ comprising the IQ domain with adjacent acidic region as well as the full sequence Ng were used in the simulations (please find the sequences in session I from the *Supporting Information*.).

We adopted a side-chain- C_{α} coarse-grained model to represent protein/peptides developed by Cheung et al [\(7\)](#page-77-6). In this model, each amino acid (except glycine, which is represented by a C_α bead) is represented by two beads: the C_α bead is located at the C_α position to represent the backbone atoms; the side-chain bead is located at the center of mass of the sidechain atoms to represent the side-chain atoms. The use of side-chain- C_{α} model enables us to capture large structural fluctuations of the proteins at a large timescale that warrants a wide search of the phase space.

The total potential energy E of Ng or Ng₁₃₋₄₉ is given by,

$$
E^{Ng} = E_{bond} + E_{angle} + E_{dihedral} + E_{chirality} + E_{elec} + E_{HB} + E_{LI}
$$
 Eq. (S1)

The bonded interactions include the bond energy (E_{bond}) , the bond-angle energy (E_{angle}) , and the dihedral-angle energy $(E_{dihedral})$ the L-isomer restraint of chirality $(E_{chirality})$. E_{bond} , E_{angle} , and *Echirality* constrain the bond length, bond angle and side-chain chirality through harmonic potentials. The formula can be found in our previous work [\(8\)](#page-77-7). The equilibrium bond length, bond angle and side-chain chirality parameters were taken from the crystal structure of apoCaM and IQ motif of Ng (PDB ID: 405E), and those of the missing segment were obtained from the structures predicted by the Sparks-X protein structure prediction server [\(9\)](#page-77-8). Since these terms are mainly constrained by chemical rules, they do not vary significantly by the conformations.

The electrostatic energy (E_{elec}) between each two beads (C_{α} or side-chain bead) *i* and *j* with partial charges is described by Debye-Hückel potential [\(10\)](#page-77-9) to include the screening effect of the electrolyte solution.

$$
E_{elec}^{ij} = \frac{q_i q_j}{4\pi \epsilon_r \epsilon_0 r_{ij}} \exp\left(-r_{ij} / \sqrt{\frac{\epsilon_r \epsilon_0 k_B T}{2e^2 I}}\right)
$$
 Eq. (S2)

where q_i (q_j) is the partial charges on bead *i* (*j*); *e* is the elementary charge (see the method of generating the partial charges in the next section II.2 from the *Supporting Information*); ϵ_r is the relative dielectric constant and is set to 80 for aqueous solution; ϵ_0 is the permittivity of the vacuum; r_{ij} is the distance between beads *i* and *j*; k_B is Boltzmann constant; T = 1.1 ε / k_B is the temperature; $I = 0.15$ M is the ionic strength.

For the backbone hydrogen bonding interactions, we adopted an angular‐dependent function that captures directional properties of backbone hydrogen bonds [\(7\)](#page-77-6),

$$
E_{HB}^{ij} = \sum_{|j-i|>3} A(\rho) \varepsilon \left[\left(\frac{\rho_{ij}^0}{r_{ij}} \right)^6 - 2 \left(\frac{\rho_{ij}^0}{r_{ij}} \right)^{12} \right]
$$
 Eq. (S3)

$$
A(\rho) = \frac{1}{\left[1 + (1 - \cos^2 \rho)(1 - \frac{\cos \rho}{\cos \rho \alpha})\right]^2}
$$
 Eq. (S4)

where the strength $\varepsilon = 0.6$ kcal/mol; r_{ij} is the distance between backbone beads *i* and *j*; and $\rho_{ij}^0 =$ 4.6 Å is the typical length of a hydrogen bond; $A(\rho)$ measures the structural alignment of two interacting strands; ρ is the pseudo dihedral angle between backbone beads of the two interacting strands [\(7\)](#page-77-6); ρ_{α} is the pseudo dihedral angle of a canonical helical turn, 0.466 rad. *A*(*ρ*) preserves the tendency to form β-strands (*ρ* = 0 or π) or α-helices (*ρ* = *ρ_α*) as it is maximized to 1 in either of the two cases.

For the non-bonded interactions (E_{LJ}) , between C_{α} bead *i* from the backbone (bb) and side-chain (sc) bead *j*, a pure repulsive interaction was considered,

$$
E_{LJ}^{sc-bb} = \varepsilon \left(\frac{\rho_{ij}^0}{r_{ij}}\right)^{12} \qquad \text{Eq. (S5)}
$$

where the strength $\varepsilon = 0.6$ kcal/mol; i $\rho_{ij}^0 = 0.9(\rho_i^0 + \rho_j^0)$, ρ_i^0 is the size of backbone bead *i*, 0.5 $σ (σ = 3.8 Å)$ and $ρ_j^0$ is the size of side-chain bead *j*, which is the van der Waals radius of the side-chain. 0.9 is a scaling factor to remove clashes between bulky side-chains. For the nonbonded interactions between side-chain bead *i* and side-chain bead *j*, we applied the Lennard-Jones potential,

$$
E_{LJ}^{sc-sc} = \varepsilon_{ij} \left[\left(\frac{\rho_{ij}^0}{r_{ij}} \right)^{12} - 2 \left(\frac{\rho_{ij}^0}{r_{ij}} \right)^6 \right]
$$
 Eq. (S6)

where ε_{ij} is solvent mediated interactions between the involved amino acids obtained from Betancourt-Thirumalai's study [\(11\)](#page-77-10). $\rho_{ij}^0 = 0.9(\rho_i^0 + \rho_j^0)$, $\rho_{i(j)}^0$ is the van der Waals radius of side-chain bead i (j). r_{ij} is the distance between beads i and j in Eqs. S3, S5 and S6. Since the electrostatic interactions were explicitly included in the Hamiltonian, the Lennard-Jones potential coefficient ε_{ij} was rescaled as described in our previous work [\(8\)](#page-77-7).

To best describe the intrinsic disorder nature of Ng or Ng13-49, we employed a backbone dihedral angle potential that is sequence-specific but independent on protein topology. For this we used the model introduced by Karanicolas and Brooks [\(12\)](#page-77-11), or the KB model. The dihedral angle between four adjacent α -carbons depends on the backbone dihedral angles of the two middle residues. Brooks' group produced a probability distribution for 400 possible ordered pairs of amino acid residues from a survey of the Protein Data Bank and thus related the probability distribution to potential energy ignoring the entropy contribution. The dihedral angle potential presents two minima corresponding to local α-helical and β-strand geometries. The statistical potential is modeled as a 4-term cosine series,

$$
E_{\text{dihedral}}^{\text{ijkl}} = \varepsilon_{\text{KB}} \sum_{n=1}^{4} K_{ij,n} \left[1 + \cos(n\varphi_{ijkl} - \delta_{jk,n}) \right] \tag{S7}
$$

where φ_{ijkl} is the dihedral angle formed by four consecutive C_α beads *i, j, k, l* with beads *j* and *k* in the middle, $K_{jk,n}$ and $\delta_{jk,n}$ are statistically determined constants. ε_{KB} is a factor to adjust the strength in relative to other interactions in the current model.

We tested a total of three types of dihedral potentials for a comparative study. The first dihedral angle potential is a sequence-based model as described above. The second dihedral angle potential is a structure-based potential, or SB model, where the structure of a specific segment of Ng obtained from the crystal structure (PDB code: 4E50 for Ng) was used as the reference (please see Table S4 for the residues in this segment). The SB model is composed of two-term cosine-series,

$$
E_{dihedral}^{ijkl} = \varepsilon_{SB} \sum_{n=1,3} k_{\varphi}^{n} \left[1 - \cos \left(n \times \left(\varphi_{ijkl} - \varphi_{ijkl}^{0} \right) \right) \right]
$$
 Eq. (S8)

where *i, j, k* and *l* are four consecutive C_α beads. φ_{ijkl} is the dihedral angle formed by those four beads. The equilibration position of the dihedral angle φ_{ijkl}^0 for this specific segment G25-A42 was taken from the crystal structure (PDB code: 4E50 for Ng). $k_{\varphi}^1 = 2k_{\varphi}^3 = 2\varepsilon$. ε_{SB} is used to adjust the barrier of the dihedral angle potential. For the unstructured segment $\varepsilon_{SB} = 0$.

The third dihedral potential is the hybrid of the two models, or Hybrid model, as shown in Table S4. We replaced the dihedral angle potential of the unstructured segment (beyond G25- A42) from the SB model with a KB model.

Different barrier values of the structure-based dihedral potential ε_{SB} and the statistical dihedral potential ϵ_{KB} were tested. Fig. S7 shows the helicity of the fragment G25-A42 decreases with temperature in all the cases. At the temperature $T = 1.1 \epsilon/k_B$, for the SB model, the helicity fell within the experimental range when $\varepsilon_{SB} = 0.3$ (Fig. S7A); for the Hybrid model, the helical fraction fell within the experimental range with when $\varepsilon_{SB} = 0.3$, and ε_{KB} did not make any influence in the range from 1.8 to 2.2 (Fig. S7B); for the KB model, the helical fraction was best matched when ε_{KB} is in the range from 1.5 to 2.2 (Fig. S7C).

To determine which best represents the properties of Ng among all the three models, we further investigated them at the residue level. Therefore, for those models with parameters that matched the overall helicity (Fig. S7), we computed the nuclear magnetic resonance (NMR) model free order parameter S^2 for the backbone beads from our coarse-grained simulations (please find the details of calculation in V.1 from the *Supporting Information*), and compared with data from the NMR experiments [\(3\)](#page-77-2). The computed S^2 (Table S5) positively correlated with the experimental values in all cases. Interestingly, for the KB model, the correlation was overall higher than the SB and Hybrid model, especially for $\varepsilon_{KB} = 1.8$. Because Ng protein is intrinsically disordered [\(13\)](#page-77-12), The KB dihedral angle potential of no bias to any specific structure enables sampling of a broad spectrum of Ng conformations, whereas the SB model and the Hybrid model of full or partial bias to a specific structure limit the flexibility to explore more conformations. Therefore, for modeling intrinsically disordered protein Ng, we adopted the sequence-based KB dihedral potential model. $\varepsilon_{KB} = 1.8$ was used in the following study of apoCaM-Ng13-49 binding simulations.

II.2 Determination of partial charges for the coarse-grained models

II.2.a Neurogranin protein

We adopted a multi-scale method developed by Cheung group [\(14\)](#page-77-13) to assign partial charges to the intact Ng protein. Firstly, we ran REMD simulations without electro-static interactions (please find the details in III.1 in the *Supporting Information*) and obtained the free energy surface $F(\Delta, \chi)$ as a function of asphericity Δ and overlap function χ . Δ measures the shape of the protein: it is like a rod or a sphere when $\Delta = 1$ or 0, respectively [\(15\)](#page-77-14). χ measures the similarity to the reference structure [\(16\)](#page-77-15). Using the free energy, we thus selected about 400-600 frames from the REMD simulations of Ng protein through importance sampling [\(17\)](#page-78-0). We reconstructed the coarse-grained structures into atomistic protein models [\(17\)](#page-78-0), used H++ server [\(18\)](#page-78-1) to predict the protonation states and computed the partial charges using the semi-empirical

quantum chemistry program MOPAC [\(19\)](#page-78-2). We collected the partial charges on the backbone (side-chain) atoms as the charge for the C_α (sidechain) bead. We then obtained the averaged partial charges on the C_{α} and side-chain beads over all the structures. We repeated the same process for all the three models shown in Table S4 and the three sets of charges we obtained were highly similar to each other (the correlation coefficients were \sim 1.00). We therefore used only one set of charges for all the three models. The charges on Ng protein are provided in Table S9.

II.2.b Apo-calmodulin and neurogranin peptide system

We followed the same procedures as for the Ng protein. The free energy surface $F(d_{COM}, Z_{inter})$ was obtained from US simulations of apoCaM and Ng₁₃₋₄₉ without electrostatic interactions by using WHAM [\(20,](#page-78-3) [21\)](#page-78-4) (please see details about the US simulations in III.2 in the *Supporting Information*). d_{COM} is the distance between center of mass of apoCaM and center of mass of $Ng₁₃₋₄₉$, and Z_{inter} is the total number of intermolecular contacts (please see the definition of a contact in V.3 from the *Supporting Information*). Two groups of partial charges were generated for apoCaM and Ng13-49 according to the experimental conditions: the fluorescence experiments for measurement of the dissociation constant of apoCaM and Ng_{13-49} were conducted at pH = 7.2 and at ionic strength $I = 0.15$ M; the NMR experiments for determining the change in chemical shifts of apoCaM upon binding Ng₁₃₋₄₉ were conducted at $pH = 6.3$ and at ionic strength I = 0.10 M. The protonation states of histidine residues were determined by the H++ server [\(18\)](#page-78-1) before using MOPAC [\(19\)](#page-78-2). The calculated partial charges of apoCaM and Ng13-49 in the above two conditions were provided in Table S10-S13*.*

III. Coarse-grained Molecular Simulations

III.1 Replica Exchange Molecular Dynamics simulations

To study dynamics of unbound Ng protein, which has a rugged energy landscape because of its intrinsic disorder, we applied replica exchange molecular dynamics (REMD) [\(5,](#page-77-4) [6\)](#page-77-5) to enhance the sampling. For the three dihedral angle models, a range of barriers of dihedral angel potential were investigated: for the SB (structure-based) model, $\varepsilon_{SB} = 0.1, 0.2, 0.3, 0.4, 0.5, 0.6$ and 0.7; for the Hybrid model, $\varepsilon_{SB} = 0.3$, 0.4 for the part with residual structure (G25-A42), in combination with $\varepsilon_{KB} = 1.8$, 2.0 and 2.2 for the rest; for the KB model, $\varepsilon_{KB} = 1.0$, 1.5, 1.8, 2.0, 2.2, 2.5 and 3.0. For each of the models, 24 replicas were distributed at temperatures $T = 0.80$, 0.82, 0.83, 0.85, 0.87, 0.88, 0.90, 0.93, 0.97, 1.00, 1.03, 1.07, 1.10, 1.13, 1.17, 1.20, 1.23, 1.27, 1.30, 1.33, 1.40, 1.47, 1.57, 1.67 ε / k_B to produce ample exchanges between the replicas (ε is the reduced energy unit = 0.6 kcal/mol, and k_B is the Boltzmann constant). The average acceptance ratio of exchanging among the neighboring replicas ranges from 0.57 to 0.77. An exchange between neighboring replicas was attempted every 100,000 time steps. In order to explore more conformation space, in each REMD simulation, a total of 6 sets of distinct initial configurations were used, making up a total of 1200 exchange attempts for each model.

III.2 Umbrella sampling simulations

Umbrella sampling (US) method [\(22\)](#page-78-5) was used to explore the thermodynamic properties of the apoCaM-Ng system. The distance between the center of mass of apoCaM and the center of mass of Ng₁₃₋₄₉ d_{COM} was used as the reaction coordinate. d_{COM} was restrained by a harmonic force with spring constant 66.7 ε/σ^2 ($\sigma = 3.8$ Å is reduced length unit). The equilibrium positions of the harmonic force range from 0.2 σ to 20.0 σ with a bin size of 0.2 σ, making up a total of 100

windows. For each window, 10,000,000 time-steps of constrained molecular dynamic simulation were carried out.

To generate the initial structures at each window of d_{COM} for the US simulations, molecular dynamics simulations were carried out for the apoCaM-Ng13-49 complex from the bound form (PDB code: 4E50, in the coarse-grained model) at a high temperature $T = 1.33 \epsilon / k_B$ for dissociation ($\varepsilon = 0.6$ kcal/mol is reduced energy unit, and k_B is Boltzmann constant). A total of 5 sets of initial configurations for the following umbrella sampling simulations were generated from 5 dissociation trajectories.

US simulations were then performed without electrostatic interactions for generating ensemble of structures to generate ensemble of structures for partial charge determination (see 2.3 for the details). After we obtained the partial charges for apoCaM and Ng_{13-49} , US simulations were carried out with charge-charge interactions to study how the strength of nonelectrostatic intermolecular interactions influences the binding affinity and to determine the optimal strength by comparing with the experiments [\(4\)](#page-77-3). Using this strength of non-electrostatic intermolecular interactions, US simulations were carried up at appropriate experimental conditions to study the binding thermodynamic properties of apoCaM-Ng13-49 system. Each set of US simulations were performed from 5 different initial structures at the temperature $T = 1.1$ ε/k_B . We analyzed the data using Weighted Histogram Analysis Method (WHAM) [\(20,](#page-78-3) [21\)](#page-78-4).

IV. Calculation of difference in binding free energy using Jarzynski's equality from steered molecular dynamics simulations

IV.1 Selection of initial atomistic structures

holoCaM-Ng13-49 complex: We employed coarse-grained molecular simulations to efficiently sample a broad ensemble of complex structures. We used the experimental measurements as a guide to strategically select several structures from the major cluster of coarse-grained complex structures. According to the crystal structure of apoCaM and IQ motif of Ng (NgIQ) (PDB ID: 4E50), the two EF-hand motifs from cCaM are open (EF-hand angles ≈ 101) and the two EFhand motifs from nCaM are closed (EF-hand angles \approx 129) (Table S3). The EF-hand angles were computed as the angle between the vectors that define the orientation of the two helices in the EF-hand motif; the vectors were defined from the center of mass of the C_{α} beads of the first 4 residues to the center of mass of the C_{α} beads of the last 4 residues in a helix of the EF-hand motif. Moreover, NgIQ binds to cCaM in the crystal structure and the NMR experiments [\(4\)](#page-77-3) showed that Ng13-49 mainly interacts with cCaM. Therefore, we select the 4 structures from the simulations by following the criteria: (i) EF-hand motifs from cCaM must be open (EF-hand angles range from 85̊ to 105̊) and EF-hand motifs from nCaM must be closed (EF-hand angles are greater than 125°; (ii) cCaM has more interactions with Ng₁₃₋₄₉ than nCaM (Z_c – Z_n > 0). Using SCAAL method [\(17\)](#page-78-0), we reconstructed all-atomistic structures of apoCaM-Ng13-49 from the selected sidechain-C_α configurations, as shown in Fig. S8. Four Ca²⁺ ions were added and their positions were estimated as the center of mass of sidechain beads of the corresponding Ca^{2+} co-ordination residues. The Ca^{2+} positions were further optimized by energy minimization on the all-atomistic structures of $Ca^{2+}-CaM-Ng_{13-49}$ (holoCaM-Ng₁₃₋₄₉) using a gradient descent algorithm provided by the molecular dynamics software package GROMACS (version 5.0.6) [\(23\)](#page-78-6) with AMBER99SB-ILDN force-field [\(24\)](#page-78-7). The tolerance for the maximum force was set to 500 kJ / mol / nm to remove clashes between atoms.

holoCaM-CaMKII complex: We used the crystal structure of calcium-bound CaM-CaMKII complex (PDB ID: 1CDM) for Ca^{2+} binding free energy calculations.

holoCaM: We used the crystal structure of holoCaM (PDB ID: 1CLL) for assessing the Ca^{2+} binding free energy in the absence of the CaMBT as a reference.

IV.2 Protonation of the initial structures for pulling simulations

The charge distribution impacts accurate estimation of the binding free energy of the Ca^{2+} ions. We first took into consideration of the pH and ionic strength by using H++ server [\(18\)](#page-78-1) to predict the protonation states of the all-atomistic structure. We input the following parameters to $H++$ as used in the experiment: $pH = 7.4$, ionic strength I = 0.15 M, the external dielectric constant is 78.4 and the internal dielectric constant is 4. After we obtained the protonated states of all the residues, we performed an energy minimization using steepest descent algorithm and the tolerance for the maximum force was set to 500 kJ / mol / nm. Then we assigned partial charges according to the geometry of the proteins (input is the protonated structure in PDB format) by using a semi-empirical quantum chemistry program MOPAC [\(19\)](#page-78-2). We applied those protonated protein structures and partial charges for further all-atomistic calculations.

IV.3 Steered molecular dynamics simulations

Classical molecular dynamics simulations were carried out using GROMACS molecular dynamics package (version 5.0.6) [\(23\)](#page-78-6). We used the AMBER99SB-ILDN force field [\(24\)](#page-78-7) except the charge assignment. The partial charges were generated from MOPAC as explained in the

previous session. The rigid three-site TIP3P model [\(25\)](#page-78-8) was used for water molecules. We neutralized the system with $Na⁺$ and Cl⁻ ions and maintained an ionic strength of 150 mM.

The size of the box is about 10 nm x 10 nm x 10 nm. The proteins are placed at least 1 nm away from the edges of the cubic box. Periodic boundary conditions were employed to mimic the macroscopic settings for electrolytes. Electrostatic interactions between periodic images were treated using the particle mesh Ewald approach [\(26\)](#page-78-9), with a grid size of 0.16 nm, fourth-order cubic interpolation and a tolerance of 10^{-5} . A cutoff of 10 Å was used for van der Waals interactions and real space Coulomb interactions as well as for updating neighbor lists.

We adopted a gradient descent algorithm for energy minimization. Then we gradually heated the system temperature by 50 K per 0.1 ns to 298.15 K using NVT ensemble. We carried out 1 ns NPT equilibration with heavy atoms of the proteins (including the four Ca^{2+} ions) fixed. The proteins as well as the Ca^{2+} ions were afterwards released and were further equilibrated for another 1 ns. All NPT simulations maintained a constant pressure of 1 bar and temperature of 298.15 K using the Parrinello-Rahman barostat [\(27\)](#page-78-10). The bond lengths in proteins were constrained using the LINCS algorithm of Hess [\(28\)](#page-78-11). The equation of motion was integrated using a 2-fs time steps. As discussed in the main text, the positions of the Ca^{2+} ions in the bound state change during the minimization and equilibration stage, which yields inaccurate estimation of the binding free energy, therefore, we froze the positions of the Ca^{2+} ions as well as the backbone heavy atoms of the protein (or protein complex) during these preparation stage. They were free to move afterwards in the pulling simulations.

For each of the initial structures including 1 structure for each of holoCaM and holoCaM-CaMKII and four selected structure of holoCaM-Ng₁₃₋₄₉, I pulled the Ca²⁺ from site III and IV independently to 2 nm away where the interaction between the Ca^{2+} and the corresponding Ca^{2+}

binding loop is negligible. During the pulling simulations, the positions of the proteins may shift and this would cause inaccurate estimation of the distance between Ca^{2+} and the corresponding Ca^{2+} binding loop. Therefore, we fixed the C_α atom of the 100th or 136th residue in CaM, which showed smallest RMSD in a separate equilibration simulation, while pulling Ca^{2+} from the binding site III or IV, respectively. The force constant k and speed of pulling ν of the reference position are described in the next section. The direction of the pulling force was randomly assigned and pointed away from the center of mass of CaM to avoid clashes between CaM and the Ca²⁺. The pulling direction was selected if the angle Ω between the pulling vector and the vector connecting center of mass of CaM and center of mass of the corresponding Ca^{2+} binding loop was within 90 degrees since $\Omega > 90$ leads to a large work (as shown in Fig. S12). The displacement of the Ca^{2+} ion and the pulling forces were output every 20 fs for calculation of the work as shown in Eq. S9. The coordinates and velocities of the system were saved every 1 ps.

The setup of the pulling simulations is illustrated in Fig. S9. The Ca^{2+} is constrained to a reference position that is moving along \vec{x} direction at a speed of $|\vec{v}| = 0.001$ nm/ps. The force constant of the spring *k* is set to 10000 kJ / mol / nm² to guarantee that the Ca^{2+} strictly follows the reference position. Therefore, the pulling force is calculated as,

$$
\vec{f}(t) = -k(\vec{x}(t) - \vec{v}t - \vec{x}_0)
$$
 Eq. (S9)

where \vec{x} (\vec{x}_0) is the instantaneous (initial) displacement of the Ca²⁺ from the center of mass of the Ca^{2+} binding loop.

IV.4 Justification of parameters used in steered molecular dynamics simulations

We set the pulling speed ν and the spring constant of the external force k to effectively estimate Ca^{2+} binding free energy from steered molecular dynamics. We explored different

combinations of spring constant and pulling speed to rationalize the parameter setting. For example, to pull the Ca^{2+} from site III of holoCaM, we used $k = 100$, 1000, 10000, 100000 kJ/mol/nm² and $v = 0.001$, 0.01, 0.02 nm/ps. The simulation time was 4 ns, 0.4 ns and 0.2 ns to ensure that Ca^{2+} was pulled the same distance. In Fig. S10A, with fairly slow pulling speed, using $k = 1000$, 10000 and 10000 kJ/mol/nm², the Ca²⁺ follows the motion of the dummy bead. For $k = 100 \text{ kJ/mol/nm}^2$, the pulling force is so weak to pull the Ca²⁺ out during the entire simulation. For $k = 1000 \text{ kJ/mol/nm}^2$, there exists a lag in the beginning of the simulation, indicating a favored interaction between CaM and $Ca²⁺$ at site III. However, the tail that deviates from the straight line indicates that the thermal fluctuations in the unbound state of Ca^{2+} dominate its motion. By increasing *v* to 0.01 or 0.02 nm/ps, $k = 1000$, 10000 or 100000 kJ/mol/nm² meet the stiff spring approximation, shown in Figs. S9B and S9C. Comparing to $v =$ 0.001 nm/ps, the displacement curves are much smoother.

During the exploration of factors that influence work values, we found that the pulling speed ν is more sensitive. Jarzynski's group [\(29\)](#page-78-12) showed that JE does not depend on the pulling speed. In their study, they showed for the same amount of simulation time, the estimation of the chemical potential of a simple Lennard-Jones fluid did not vary for different switching time (comparable to pulling speed v). However, in practice, the free energy calculation is more efficient with fewer slow pulling trajectories than more fast pulling trajectories given the same amount of simulation time [\(39\)](#page-79-0). For a complex system of pulling Ca^{2+} form a holoCaM-CaMBT compound, the efficiency of convergence of distribution of work values depend on the pulling speed ν . Moreover, it has been shown that slower pulling speed ν reduces the perturbation from the pulling force compared to the level of the thermal fluctuations of the binding site [\(30\)](#page-78-13). Therefore, in this study, we set a relatively slow pulling speed $v = 0.001$ nm/ps with a stiff spring constant k = 10,000 kJ/mol/nm² to guarantee that the Ca^{2+} follows the moving reference bead that accounts for small work. To improve the efficiency of the free energy estimation, furthermore, we also employed cumulative integral (CI) extrapolation method developed by Zuckerman's group [\(31\)](#page-78-14), which is discussed in the next session.

In Fig. S11, the typical force profiles lead to the same conclusion. At $v = 0.001$ nm/ps, *k* $= 10,000$ and 100,000 kJ/mol/nm² demonstrated the same trend in the force profiles except more thermal fluctuations in the latter, indicating a converged behavior of the dissociation of Ca^{2+} in this parameter range. Therefore, $k = 10,000 \text{ kJ/mol/nm}^2$ were used for the pulling simulations.

IV.5 Estimation of binding free energy using Jarzynski's equality and Cumulative Integral extrapolation

After obtaining the work (*w*) distribution from the pulling simulations as explained in IV.3 in the *Supporting Information*, we estimated the binding free energy $\Delta G^{\text{holoCaM-CaMBT}}$ between a one- Ca^{2+} -missing complex of holoCaM-CaMBT (G_U) and a Ca²⁺-loaded complex of holoCaM-CaMBT (G_B) according to JE [\(32\)](#page-78-15),

$$
\Delta G^{\text{holoCaM-CaMBT}} = G_{\text{B}} - G_{\text{U}}
$$

$$
= k_B T ln \left\langle e^{-\frac{W}{k_B T}} \right\rangle
$$

$$
\approx k_B T ln \sum_{n=1}^{M} e^{-w_n / k_B T} / M
$$
 Eq. (S10)

where k_B is Boltzmann constant, T is the temperature, M is number of pulling simulations (M = $100 \sim 150$, w_n is the work in nth pulling calculated as,

$$
w_n(\mathbf{B} \to \mathbf{U}) = \int_0^t \overrightarrow{f_n}(\tau) \cdot \vec{v}_n d\tau
$$
 Eq. (S11)

B and U stands for bound and unbound states of the Ca^{2+} ion, respectively.

However, accurate estimation of the binding free energy ΔG relies on the sufficient sampling of small work [\(32\)](#page-78-15). Zuckerman's group [\(31\)](#page-78-14) showed that using a cumulative integral (CI) to extrapolate the free energy estimation to $n \rightarrow \infty$ can help reduce the required simulation data by 5 to 40 fold [\(30,](#page-78-13) [31\)](#page-78-14), where n is the number of work values. In Table S6, we showed the estimation of change in Ca^{2+} binding free energy ($\Delta\Delta G$) by using JE or CI estimation. We found that JE or CI produced the same signs of the $\Delta\Delta G$, however, the magnitudes were not the same. The discrepancy in the free energy calculation lies in the requirement of sufficient sampling of small work for direct JE calculations, which can be overcome by CI extrapolation [\(28\)](#page-78-11).

V. Analyses

V.1 Calculation of model free order parameter S2

A number of $N \approx 500$ representative coarse-grained structures were selected from the free energy surface of unbound Ng protein using importance sampling method [\(17\)](#page-78-0). The NMR model-free order parameter S_i^2 [\(33\)](#page-78-16) for the backbone N-H bond of residue *i* is calculated on the selected structures by relating to the Root Mean Square Fluctuations (RMSF) of the C_{α} bead through the following empirical relation [\(34\)](#page-78-17),

$$
S_i^2 = 1 - 0.5 * \ln(1 + \frac{RMSF_i}{23.6\text{\AA}} * 10.0)
$$
 Eq. (S12)

$$
RMSF_i = \sqrt{\frac{1}{N} \sum_{j=1}^{N} (r_i^j - \frac{1}{N} \sum_{k=1}^{N} r_i^k)^2}
$$
 Eq. (S13)

where r_i^j is the position of the C_α bead of residue *i* at frame *j* and *N* is the total number of frames. RMSF*ⁱ* was computed after structural alignment.

V.2 Definition of helicity

The fraction of helical structure, or helicity, of Ng (G25-A42) was calculated using the following formula [\(35\)](#page-79-1),

$$
f_H = \frac{1}{N-3} \sum_{i=1}^{N-3} < \Theta(\Delta \phi - |\phi_i - \phi_0|) > \tag{S14}
$$

where $\Theta(x)$ is the Heaviside step function taking value 1 if $x \ge 0$ and value 0 otherwise. *N* is the total number of residues, *i* is the residue index, ϕ_i is the dihedral angle about the residues $i-i+3$ from the simulation, $\phi_0 = 49.50^{\circ}$ is the dihedral angle in a perfect alpha helix and $\Delta \phi = 12.07^{\circ}$. <•••> denotes ensemble average.

V.3 Definition of Z

We defined an order parameter Z to calculate the total number of intermolecular contacts between apoCaM and Ng13-49. For each residue *i* from apoCaM and residue *j* from Ng13-49, the sidechain-sidechain and backbone-backbone contacts $(Z_{ij}^{ss}$ and Z_{ij}^{bb} , respectively) are determined as

$$
Z_{ij}^{ss} = \Theta\left(c - \frac{d_{ij}^{ss}}{\rho_i + \rho_j}\right) \qquad \qquad \text{Eq. (S15)}
$$

$$
Z_{ij}^{bb} = \Theta(c - \frac{d_{ij}^{bb}}{\rho_{HB}})
$$
 Eq. (S16)

 $d_{ij}^{ss}(d_{ij}^{bb})$ is the distance between side-chain (backbone) beads of residue *i* of apoCaM and residue *j* of Ng₁₃₋₄₉ in simulation, ρ_i (ρ_j) is van der Waals radius of residue *i(j)*, $\rho_{HB} = 4.66 \text{ Å}$ is the typical length of a hydrogen bond and cutoff $c = 1.3$. The total number of contacts Z is the summation of backbone-backbone and sidechain-sidechain contacts over all residues $Z =$ $\sum_{ij} (Z_{ij}^{ss} + Z_{ij}^{bb}).$

V.4 Sampling protein configuration for structure analyses

The structures from US simulations are biased and not appropriate for structural analyses directly. We therefore used the Boltzmann reweighting method to sample an ensemble of unbiased structures. The probability of selecting a configuration x is

$$
w(x) = \begin{cases} 1 & \text{if } p(x) \ge \rho \\ 0 & \text{else} \end{cases} \quad \text{where } p(x) = e^{-G(d_{COM}(x))/k_B T} \quad \text{Eq. (S17)}
$$

 $p(x)$ is the probability of the configuration *x* in reweighted ensemble; $G(d_{COM})$ is the reweighted/unbiased free energy obtained from the WHAM analyses along reaction coordinate d_{COM} (G_{min} is scaled to 0); ρ is a random number generated between 0 and 1; k_B is Boltzmann constant and *T* is the temperature. The ensemble generated after Boltzmann reweighting thus obeys the canonical distribution and is employed for subsequent analyses. Each structure from the biased ensemble was challenged by acceptance probability $w(x)$ and a total of 23,722 structures were sampled. This sample achieves a distribution of P^{sample} (d_{COM}). In order to assess the quality of the sampling, we computed the surprisal value compared with the original unbiased distribution $P^{\text{ori}}(d_{\text{COM}})$ defined by the following formula [\(36\)](#page-79-2),

$$
surprisal = \sum_{r} -P^{ori}(r) \ln \left[P^{sample}(r) / P^{ori}(r) \right]
$$
 Eq. (S18)

where the summation is over all the order parameter r (in this case d_{COM}). A surprisal value of 0.14 ensures that the sampled ensemble can well represent the original distribution.

V.5 Self-organized clustering algorithm

We applied the self-organized neural-net clustering algorithm [\(37-39\)](#page-79-3) to determine the structures of the apoCaM-Ng13-49 complexes from the umbrella sampling simulations. In this clustering method a vector with *M* elements represents each conformation *j*, $x_j = [x_{1j}, x_{2j}, ..., x_{Mj}]$, where $j = 1, 2, ... N$ and N is the number of conformations selected for clustering analysis. The element x_{ij} ($i = 1, 2, ..., M$) stands for the Euclidean distance between side-chains of the

polypeptide chain in the conformation *j*. To partition the various conformations into distinct clusters, the clusters are described by the cluster center and the size of the cluster is determined by a radius R_c . A given conformation is assigned to cluster k if the distance between the vector x_j and the center of the k^{th} cluster,

$$
C_k = \frac{1}{M_k} \sum_{j=1}^{M_k} x_j,
$$
 Eq. (S19)

where $C_k = [C_{1k}, C_{2k}, ..., C_{Mk}]$ and M_k is the total number of conformations in C_k .

Conformation *j* belongs to C_k if the Euclidean distance between conformation *j* and the cluster center *k*, $d_{jk} = |x_j - C_k| < R_c$, where R_c is a preassigned value. In the current study we used the native contact pairs from the unbound structure of apoCaM and Ng₁₃₋₄₉ as the M elements and a following cutoff distance $R_c = 25 \sigma (\sigma = 3.8 \text{ Å})$ to categorizes the structures into seven distinct clusters.

V.6 Correlation between intra- and inter- molecular interactions in apoCaM-Ng13-49

binding

In order to better understand the relation between interactions within the two proteins and intermolecular contacts, we built a correlation map between contacts (Fig. S4). The correlation between two contact pairs m and n is computed as in the following equation,

$$
corr_{mn} = \frac{q_m q_n > - \langle q_m \rangle}{\sqrt{\langle q_m^2 \rangle - \langle q_m \rangle^2} \sqrt{\langle q_n^2 \rangle - \langle q_n \rangle^2}}
$$
 Eq. (S20)

where $q_{m(n)}$ is the contact state of the contact pair m (n): 1 if m (n) is a contact, 0 if not. The list of contacts are provided in Table S8. The correlation between those contacts falls into several categories.
- a) Correlation: as shown in diagonal, the contact formation within apoCaM, between apoCaM and Ng_{13-49} , and within Ng_{13-49} are correlated.
- b) Anti-correlation: contact formation between nCaM and Ng₁₃₋₄₉ anti-correlates with contact formation within nCaM; contact formation between cCaM and Ng13-49 anticorrelates contact formation within cCaM. The anti-correlation tells that the contact formation within nCaM or that within cCaM competes with its interaction with Ng13-49. We infer that binding of the target is responsible for the repacking of the nCaM or cCaM by direct competitive interaction with the Ca^{2+} binding loops.
- c) No correlation: there is no apparent correlation among contact formation within nCaM and contact formation within cCaM, indicating the two domains of CaM are relative independent.
- d) Mixed correlation: contact formation between nCaM and cCaM has a mixed correlation with contact formation between $nCaM$ and Ng_{13-49} as well as contact formation between cCaM and Ng_{13-49} .

Supporting Figures

Fig. S1. Overview of the effects of CaM-dependent Ca2+ signaling and effects of CaM binding targets (CaMBTs) on changes in synaptic plasticity. Many of the effects of intracellular Ca^{2+} on synaptic plasticity are mediated through CaM -regulated proteins. Increase in intracellular Ca^{2+} , generated through the activity of NMDA (N-methyl-Daspartate) receptors or voltage-sensitive Ca^{2+} channels, results in the release of CaM that is bound to Ng. CaM mediates the Ca^{2+} stimulation of CaMKII which is required for changes in synaptic plasticity. The structures of Ca^{2+} -free CaM (apoCaM, PDBID: 1CFD) and Ca^{2+} -CaM-CaMKII peptide (PDB ID: 1CDM) are provided. CaM is colored as follows, red \rightarrow nCaM (residue 1 to 76), gray \rightarrow central linker (residue 77 to 82), blue \rightarrow cCaM (residue 83 to 148) and the CaMKII peptide is colored in green.

Fig. S2 Reweighted potential of mean force of apoCaM and Ng13-49. The PMF was reweighted from umbrella sampling simulations using WHAM at varying scaling factors of the inter-molecular nonbonded interaction λ (excluding electrostatic interactions). KB statistical dihedral angle potential was employed. d_{COM} is the distance between center of mass of apoCaM and center of mass of Ng₁₃₋₄₉. σ equals 3.8 Å. T = 1.1 ε / k_B. ε = 0.6 kcal / mol.

Fig. S3 Difference in probability of contact formation between the bound ensemble of the most dominant cluster (cluster 1) and the unbound ensemble. The difference maps for apoCaM intramolecular (a), Ng_{13-49} intramolecular (b) and apoCaM-Ng₁₃₋₄₉ intermolecular (c) contacts are provided. $d_{COM} = 20.0$ s is used for the unbound state. The schematic representation of the helices of apoCaM and Ng13-49 are provided along the axes. The representative structure, which has minimal root mean square deviation from the averaged structure, is in ribbon representation and colored according to the secondary structures: nCaM is in gray, cCaM is in green, acidic region of Ng_{13-49} is in pink and IQ motif of Ng_{13-49} is in orange. The regions of particular interest are encircled by ellipses, rectangles, circles and triangles. The color bar represents the difference in probability of contact formation ΔP between the bound and the

unbound conformations where positive ΔP indicates an increase of contact probabilities in the bound ensemble.

Fig. S4 Correlation map between contacts formation involving apoCaM and Ng13-49 in bound ensemble (cluster 1). The contact pair list is provided in Table S8. The blue color shows strong anti-correlation between involved contact pairs; the red color shows strong correlation; white means no correlation.

Fig. S5 Illustration of the EF-hand β**-scaffold.** The structure is from C-terminus of the crystal structure of holoCaM (PDB: 1CLL). The calcium ions are shown in yellow beads. The $5th$ (+Z) coordinating ligand) and $12th$ (-Z coordinating ligand) residues in the Ca²⁺ binding loops are shown in ball-and-stick representation.

Fig. S6 Probability of contact formation in unbound, encounter and bound ensemble of apoCaM-Ng13-49. In panel (a), the probability map of contact formation for the unbound state was calculated using the ensemble when apoCaM and Ng₁₃₋₄₉ are well separated at $d_{COM} = 20.0$ σ. σ = 3.8 Å. (b) the probability map of contact formation for the encounter complex at d_{COM} = 10.0 σ and (c) in the bound state ($d_{COM} = 2.8$ σ) are provided. The upper and lower triangles show contacts maps for backbone-backbone conatcts and sidechain-sidechain contacts, respectively. The schematic representation of the helices of apoCaM and Ng13-49 are provided along the axes. The region encircled by a circle marks the interaction between the two β-strands in the EF-hand β-scaffold. The color bar represents the probability of contact formation P.

Fig. S7 Calculated helicity of the Ng protein fragment with residual structure matches with experiment using three dihedral angle potentials. The helicity of the fragment G25-A42 was computed from REMD simulations of the intact Ng protein with the coarse-grained side-chain C_{α} model. Three models of dihedral angle potential were used: (A) Structure-Based (SB) potential; (B) Hybrid potential; (C) Karanicolas-Brooks (KB) statistical potential. The two gray lines mark the upper and lower limit of the helicity of the segment from the experiment [\(3\)](#page-77-0); the yellow shade marks the corresponding reduced temperature 1.1 ε / k_B . Helicity was estimated based on the dihedral angles between four consecutive C_{α} beads (see definition in V.2 in *Supporting Information*).

Fig. S8 Illustration of protein reconstruction from coarse-grained configuration to atomistic configuration for the holoCaM-Ng13-49 complex. The backbone of CaM and Ng13-49 is shown in ribbon representation. CaM is colored from red (N-Domain) to blue (C-Domain); Ng₁₃₋₄₉ is shown in magenta. (left) The Sidechain-C_{α} model (SCM), where the backbone atoms are represented by the C_{α} beads (blue balls) and sidechain atoms are represented by the sidechain beads (red balls). (right) All-atomistic (AA) configuration of the holoCaM-Ng13-49 complex with Ca^{2+} (yellow balls) added to the Ca^{2+} binding loops. The other atoms are colored according to the atom names (oxygen atoms are in red, carbon in cyan, hydrogen in white, nitrogen in blue, etc.).

Fig. S9 Schematic illustration of pulling Ca2+ from one of the calcium-binding loop of CaM. The rest of the system including CaM and the solvent molecules is not shown in this illustration for better visualization.

Fig. S10 The displacement of the Ca^{2+} from the Ca^{2+} binding loop during the pulling **simulations of** \bar{Ca}^{2+} **from the** Ca^{2+} **binding site III** of holoCaM. The unit of *k* is kJ/mol/nm² and the unit of *v* is nm/ps. (A) (B) (C) show the case pulling speed $v = 0.001$, $v =$ 0.01 and $v = 0.02$ nm/ps.

Fig. S11 The profile of external force in pulling simulations of Ca^{2+} from the Ca^{2+} binding **site III of holoCaM.** The pulling speed $v = 0.001$ nm / ps. The unit of spring constant *k* is kJ/mol/nm² and the unit of pulling speed ν is nm/ps.

Fig. S12 Work done by external pulling force of pulling Ca2+ at different pulling angles. At each pulling angle zone Ω in [0, 30], [30, 60], [60, 90], [90, 120], [120, 150], [150, 180], the $Ca²⁺$ were pulled from site IV of holoCaM for 100 times. The average work, deviation (upper bar) as well as standard error (lower bar) are provided. The pulling speed $v = 10$ nm / ns and the spring constant of the pulling force $k = 1000 \text{ kJ} / \text{mol} / \text{nm}^2$.

Fig. S13 The running Jarzynski equality (JE) estimate and subsampled block-averaged (BA) estimate are plotted as a function of the number of trajectories used in the estimate. The binding free energies of Ca^{2+} at site III and IV of holoCaM using the running JE method and the BA method are shown in panel (a) and the binding free energies of Ca^{2+} at site III and IV of holoCaM-Ng₁₃₋₄₉ (using the conformation shown in Fig. 5) are shown in panel (b).

Figure S14 Distribution of work for dissociation of Ca2+. (a) and (b) represent the holoCaM and holoCaM-Ng13-49, respectively. (c-f) represent the four conformations of holoCaM-Ng13-49.

Fig. S15 Sequence alignment of neurogranin (Ng). 45 sequences of Ng (see Table S14 for the details of species associated with each sequence) were selected from UniProtKB (www.uniprot.org) by searching the keyword "neurogranin". The sequences were aligned using the "Align" tool in UniProt and visualized in Jalview (40) . The ClustalX color scheme in Jalview was used to highlight the residues and shading intensity of the color is based on the conservation over all the sequences in the alignment. The "acidic region" and the "IQ motif" of Ng are indicated by the grey and black line, respectively. The "DDPG (D/E/A)" motif is indicated inside the black area. The sequence of Ng from mouse (UniProt id P60761) that was used in this study is shown inside the red dotted area.

Fig. S16 Distance map between the Ca²⁺ binding loops (III and IV) of CaM and the "DDPG" motif of Ng. The distance here refers to the closest distance between the corresponding residues. The Ca^{2+} ions (shown as yellow spheres) coordinate (indicated by grey dotted lines) with the first, third, fifth, seventh, ninth (through water molecule (w), shown as a blue sphere) and twelfth residue of the loops. The distance to distance map is based on the holoCaM-Ng13-49 complex structure in Fig. 5a in the main text.

Fig. S17 Illustration of the changes in the binding free energy ΔG and ΔΔG from pH = 7.4 to pH = 6.8. $\Delta G = G_B - G_U$. $\Delta \Delta \tilde{G} = \Delta G^{holoCaM-CaMBT}$ - $\Delta G^{holoCaM}$. B and U stand for bound and unbound states of the Ca²⁺, respectively. The arrows show increase of ΔG for Ca²⁺ in all the three systems.

Supporting Tables

Table S1 Dissociation constant K_d of apoCaM and Ng_{13-49} with several values of strength of **non-electrostatic intermolecular hydrogen bonding and van der Waals interactions λ.** λ is shown in Eq 4. Experimentally measured $K_d = 680$ nM [\(4\)](#page-77-1).

Table S2 The population of major clusters from computer simulations and the correlation coefficient of their computational "apparent chemical shifts" with the data from NMR experiments [\(4\)](#page-77-1).

Table S3 EF-hand angles in several forms of CaM. Definition of EF hand angles is provided in IV.1 from the *Supporting Information*.

Table S4 The composition of the three dihedral potentials for modeling Ng or Ng13-49. SB stands structure-based and KB stands for Karanicolas-Brooks statistical potential [\(12\)](#page-77-2).

Table S6 Difference in binding free energy of Ca2+ (ΔΔG) calculated from non-equilibrium molecular simulations and from the experiments at $pH = 7.4$ **.** Direct use of Jarzynski's equality and a cumulants integral extrapolation was used in the calculation of binding free energy of Ca^{2+} from the simulations.

System		holoCaM-Ng ₁₃₋₄₉ (average)	holoCaM-CaMKII
$\Delta\Delta G$ (kcal/mol) Jarzynski's equality	Site III	8.5 ± 2.7	-0.1
	Site IV	23.9 ± 1.1	-3.4
$\Delta\Delta G$ (kcal/mol) CI extrapolation	Site III	9.2 ± 2.2	-0.5
	Site IV	22.4 ± 0.9	-7.0
$\Delta\Delta G$ (kcal/mol) experiment	Site III/IV	$+2.5$	-3.3

Table S7 Difference in binding free energy of Ca2+ (ΔΔG) for Ca2+ at site III calculated from non-equilibrium molecular simulations at $pH = 6.8$ **.** Direct use of Jarzynski's equality and a cumulants integral extrapolation was used in the calculation of binding free energy of Ca^{2+} from the simulations.

Table S8 Contact pair list used for correlation analysis. The non-specific contact pairs are determined from the difference contact map analysis. A contact pair is selected if the magnitude of the change in the probability of the contact pair in cluster 1 (Fig. S3) is greater than 0.2 from unbound state to the bound state. The contact pairs are listed in a sequence of categories: within nCaM, between nCaM and cCaM, within cCaM, between nCaM and Ng13-49, between cCaM and Ng_{13-49} and within Ng_{13-49} . The sequence of apoCaM is from NMR structure (PDB: 1CFD) and sequence of Ng13-49 is provided in the *Materials and Methods* in the main text.

Residue Index	Residue	Charge on	Error of the	Charge on side-chain	Error of the
1	Name MET	C_{α} 0.847	C_{α} charge 0.002	0.091	side-chain charge 0.001
\overline{c}	ASP	-0.028	0.002	-0.898	0.001
3	CYS	-0.037	0.002	0.035	0.001
$\overline{\mathcal{L}}$	CYS	-0.037	0.002		0.001
				0.033	
5	THR	-0.098	0.002	0.083	0.001
6	GLU	-0.052	0.002	-0.872	0.002
$\overline{7}$	SER	-0.098	0.002	0.066	0.001
8	ALA	-0.022	0.002	0.049	0.001
9	CYS	-0.042	0.002	0.025	0.001
10	SER	-0.077	0.002	0.062	0.001
11	LYS	-0.061	0.003	1.000	0.001
12	PRO	-0.307	0.003	0.316	0.002
13	ASP	-0.010	0.002	-0.926	0.002
14	ASP	-0.037	0.002	-0.927	0.002
15	ASP	-0.040	0.002	-0.930	0.002
16	ILE	-0.070	0.002	0.078	0.001
17	LEU	-0.043	0.002	0.061	0.001
18	ASP	-0.053	0.002	-0.925	0.002
19	ILE	-0.073	0.003	0.080	0.001
20	PRO	-0.301	0.003	0.305	0.002
21	LEU	-0.066	0.002	0.063	0.001
22	ASP	-0.034	0.002	-0.930	0.002
23	ASP	-0.035	0.003	-0.929	0.002
24	PRO	-0.335	0.003	0.320	0.002
25	GLY	-0.182	0.002	0.188	0.001
26	ALA	-0.041	0.002	0.063	0.001
27	ASN	-0.047	0.003	0.011	0.001
28	ALA	-0.043	0.002	0.067	0.001
29	ALA	-0.064	0.002	0.064	0.001
30	ALA	-0.068	0.002	0.068	0.001
31	ALA	-0.057	0.002	0.069	0.001
32	LYS	-0.079	0.002	1.020	0.002
33	ILE	-0.089	0.002	0.103	0.001
34	GLN	-0.077	0.002	0.069	0.001
35	ALA	-0.046	0.002	0.064	0.001
36	SER	-0.091	0.002	0.079	0.001
37	PHE	-0.052	0.002	0.062	0.001
38	ARG	-0.079	0.002	1.027	0.002
39	GLY	-0.181	0.002	0.192	0.001
40	HIS	-0.049	0.002	0.049	0.001
41	MET	-0.069	0.002	0.069	0.001
42	ALA	-0.062	0.002	0.068	0.001
43	ARG	-0.062	0.002	1.020	0.002
44	LYS	-0.074	0.002	1.031	0.002

Table S9 Charge distribution on a coarse-grained side-chain C^α model of full length Ng.

The unit of charge is $1.6*10^{-19}$ C.

Table S10 Charge distribution on a coarse-grained side-chain C^α model of apoCaM at pH = 7.2, I = 0.15 M.

Residue	Residue	Charge on	Error of the	Charge on	Error of the
Index	Name	C_α	C_{α} charge	side-chain	side-chain charge
	ALA	1.564	0.004	0.293	0.002
2	ASP	-0.038	0.003	-0.889	0.002
3	GLN	-0.050	0.003	0.051	0.001
$\overline{4}$	LEU	-0.068	0.002	0.081	0.001
5	THR	-0.093	0.002	0.073	0.001
6	GLU	-0.051	0.002	-0.915	0.002

The unit of charge is $1.6*10^{-19}$ C.

Table S11 Charge distribution on a coarse-grained side chain C_α **model of Ng₁₃₋₄₉ at pH = 7.2, I = 0.15 M.**

Residue	Residue	Charge on	Error of the	Charge on	Error of the
Index	Name	Ca	Ca charge	side-chain	side-chain charge
13	ASP	1.574	0.004	-1.659	0.004
14	ASP	-0.020	0.003	-0.921	0.002
15	ASP	-0.034	0.003	-0.911	0.002
16	ILE	-0.071	0.003	0.081	0.001
17	LEU	-0.060	0.002	0.061	0.001
18	ASP	-0.037	0.003	-0.916	0.002
19	$\rm ILE$	-0.539	0.004	0.540	0.002
20	PRO	-0.265	0.003	0.258	0.003
21	LEU	-0.054	0.003	0.064	0.001
22	ASP	-0.042	0.003	-0.931	0.002
23	ASP	-0.818	0.004	-0.133	0.003
24	PRO	-0.282	0.003	0.261	0.002
25	GLY	-0.181	0.003	0.187	0.001
26	ALA	-0.032	0.002	0.054	0.001
27	ASN	-0.046	0.003	0.005	0.001
28	ALA	-0.040	0.003	0.064	0.001
29	ALA	-0.059	0.003	0.060	0.001
30	ALA	-0.066	0.003	0.062	0.001
31	ALA	-0.055	0.003	0.064	0.001
32	LYS	-0.079	0.003	0.996	0.001
33	ILE	-0.089	0.003	0.100	0.001
34	GLN	-0.072	0.002	0.059	0.001
35	ALA	-0.042	0.003	0.062	0.001
36	SER	-0.092	0.003	0.076	0.001
37	PHE	-0.048	0.003	0.057	0.001
38	ARG	-0.077	0.002	1.002	0.002
39	GLY	-0.175	0.002	0.187	0.001

The unit of charge is $1.6*10^{-19}$ C.

Table S12 Charge distribution on a coarse-grained side chain C^α model of apoCaM at pH = 6.3, I = 0.1 M.

Residue	Residue	Charge on	Error on the	Charge on the	Error on the
Index	Name	C_{α}	C_{α} charge	side-chain	side-chain charge
1	ALA	1.565	0.004	0.292	0.001
$\overline{2}$	ASP	-0.041	0.003	-0.885	0.002
$\overline{3}$	GLN	-0.052	0.002	0.053	0.001
$\overline{4}$	LEU	-0.066	0.002	0.080	0.001
5	THR	-0.095	0.002	0.075	0.001
6	GLU	-0.051	0.002	-0.914	0.002
$\overline{7}$	GLU	-0.072	0.003	-0.871	0.002
8	GLN	-0.040	0.003	0.053	0.001
9	ILE	-0.098	0.003	0.088	0.001
10	ALA	-0.061	0.003	0.058	0.001
11	GLU	-0.061	0.003	-0.868	0.002
12	PHE	-0.047	0.003	0.047	0.001
13	LYS	-0.090	0.003	0.998	0.001
14	GLU	-0.055	0.003	-0.893	0.002
15	ALA	-0.035	0.002	0.058	0.001
16	PHE	-0.055	0.003	0.053	0.001
17	SER	-0.118	0.003	0.070	0.001
18	LEU	-0.060	0.002	0.076	0.001
19	PHE	-0.054	0.003	0.047	0.001
20	ASP	-0.049	0.003	-0.885	0.002
21	LYS	-0.078	0.002	0.986	0.001
22	ASP	-0.049	0.003	-0.914	0.002
23	GLY	-0.196	0.002	0.191	0.001
24	ASP	-0.064	0.002	-0.945	0.001
25	GLY	-0.150	0.002	0.184	0.001
26	THR	-0.062	0.003	0.065	0.001
27	ILE	-0.066	0.003	0.077	0.001
28	THR	-0.107	0.003	0.079	0.001

The unit of charge is $1.6*10^{-19}$ C.

The unit of charge is $1.6*10^{-19}$ C.

Table S14. Detail of the neurogranin (Ng) sequences obtained from UniprotKB (www.uniprot.org). 45 sequences of Ng with unique Uniprot ID were used for the sequence alignment (see Fig. S15).

Table S14(a)

Table S14(b)

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