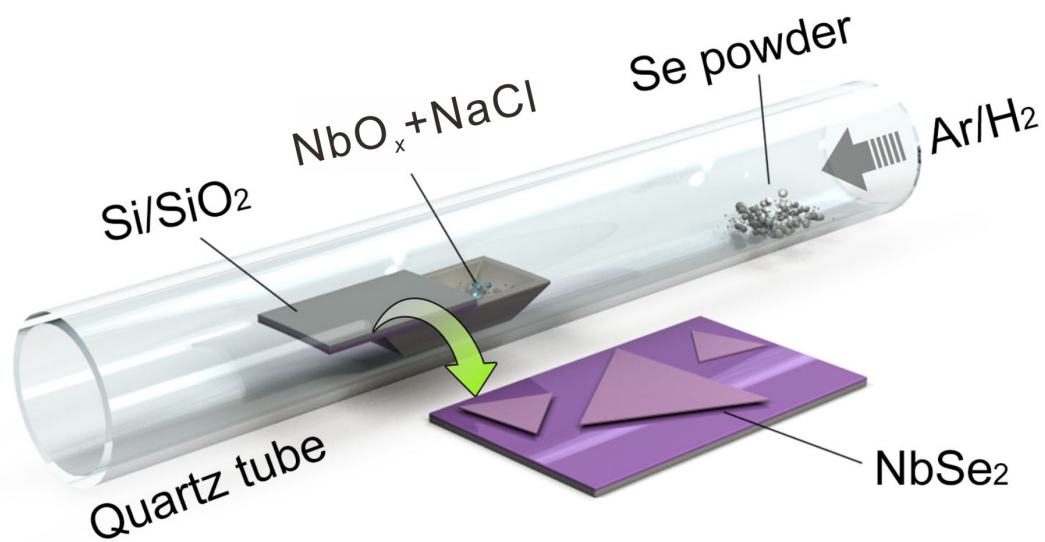
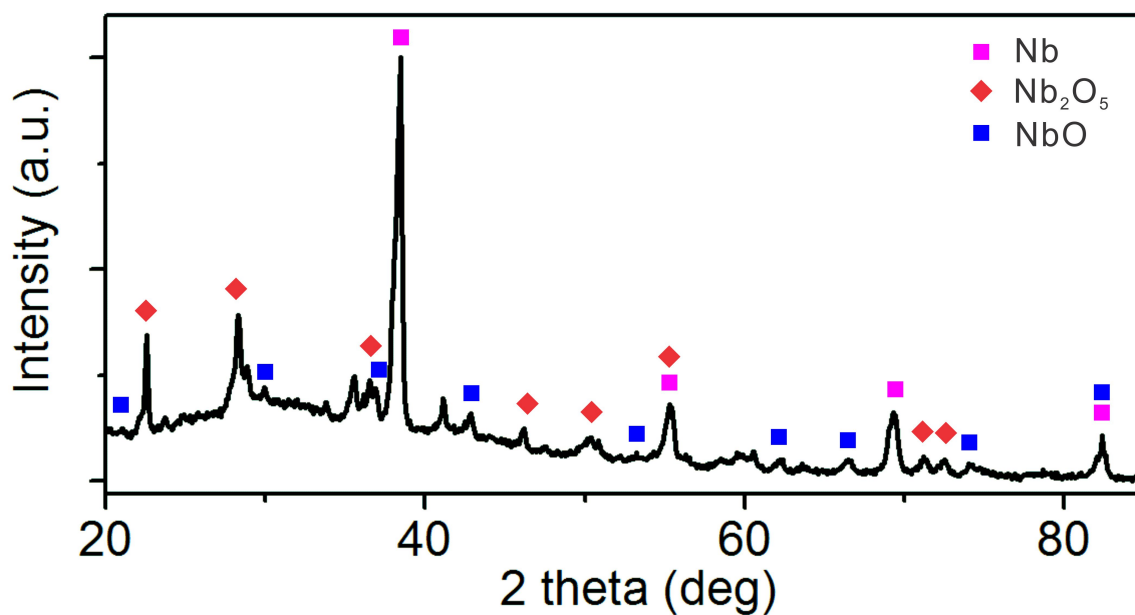


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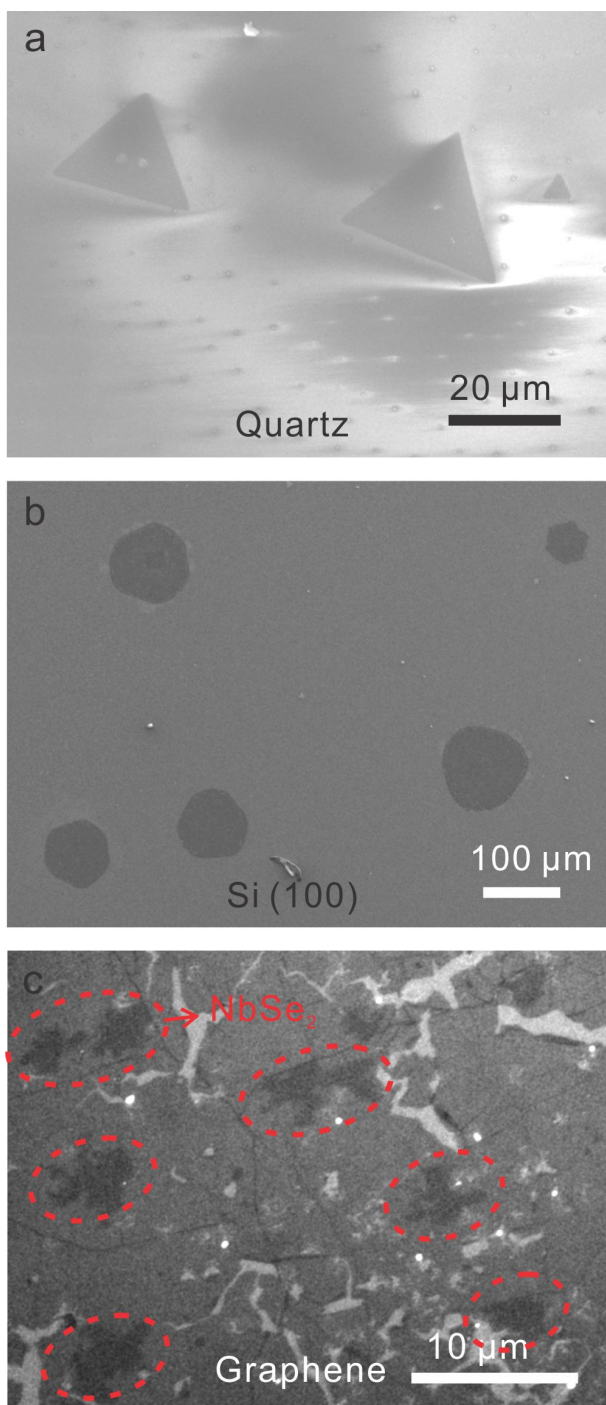
**Description:** Supplementary Figures, Supplementary Notes, Supplementary Tables, Supplementary References



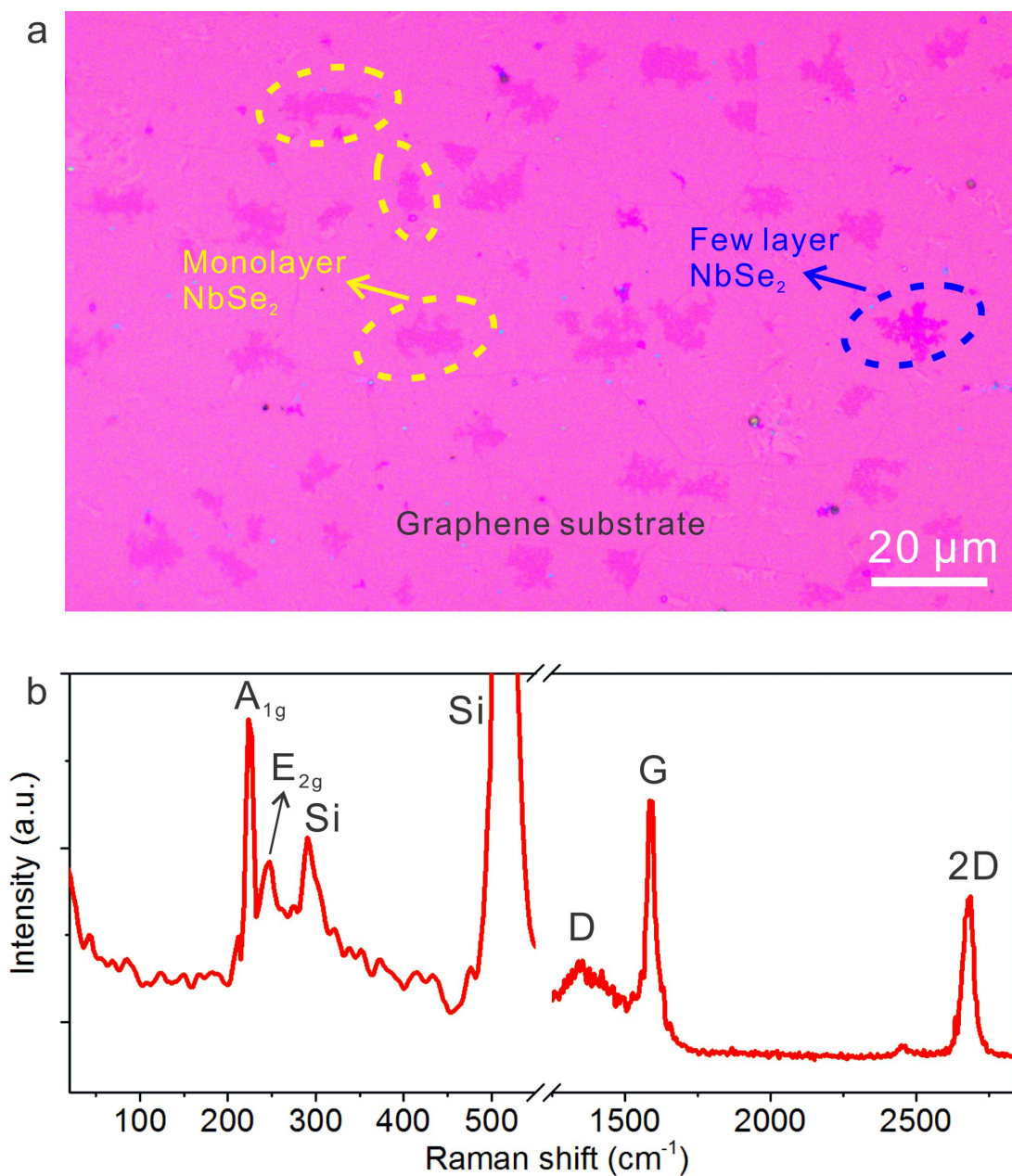
**Supplementary Figure 1** Illustration of the reaction chamber for NbSe<sub>2</sub> synthesis



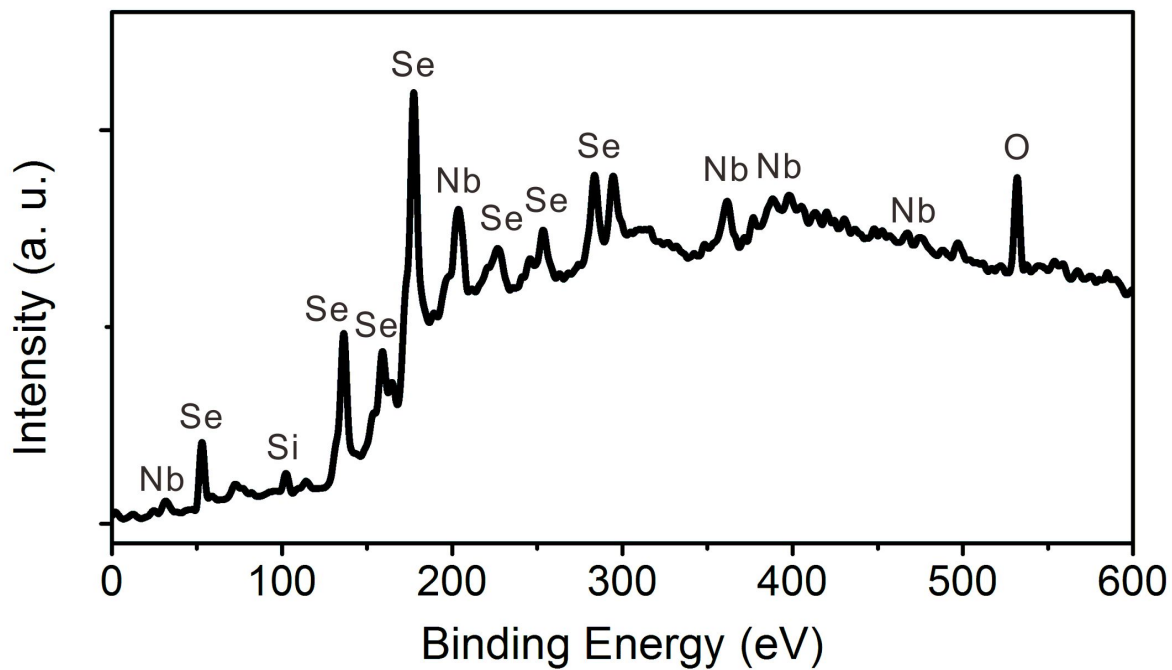
**Supplementary Figure 2** XRD pattern of the partially oxidized niobium powder NbO<sub>x</sub> ( $x \leq 2.5$ ). After oxidation, the powder contains Nb, Nb<sub>2</sub>O<sub>5</sub> and NbO.



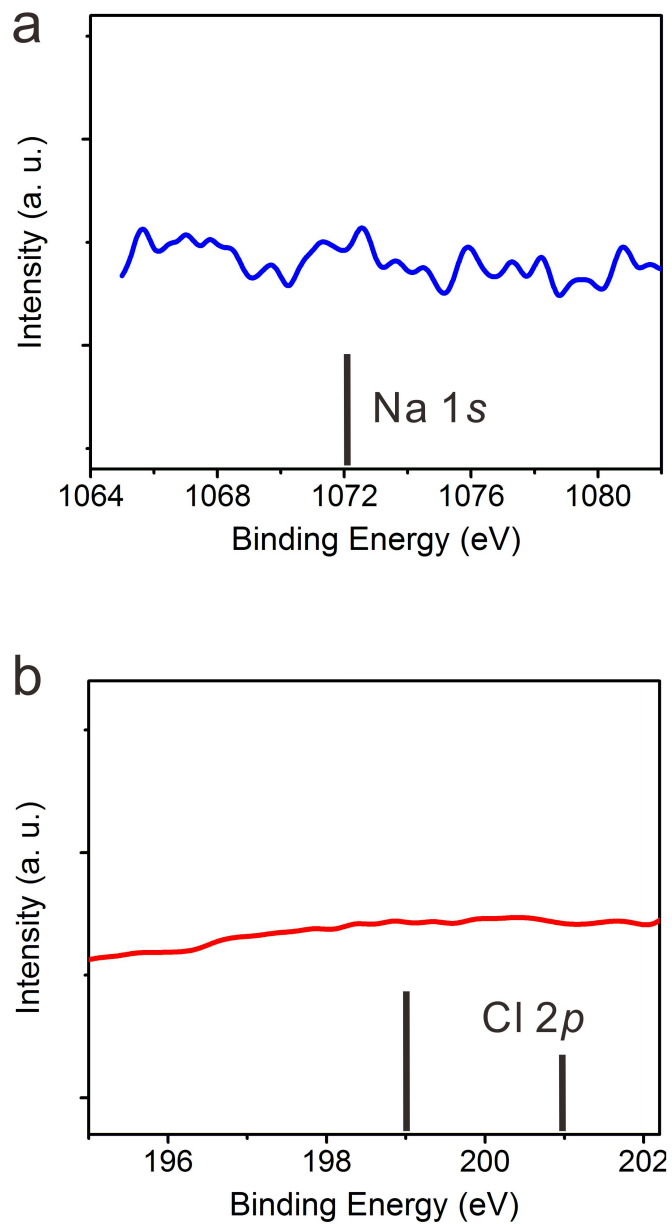
**Supplementary Figure 3.** SEM morphology of ultrathin NbSe<sub>2</sub> layers grown on diverse substrates (a) quartz, (b) Si(100) (c) CVD graphene transferred to SiO<sub>2</sub>/Si. This result indicates that NbSe<sub>2</sub> layers can be grown on various amorphous or crystalline substrates, different substrates may result in NbSe<sub>2</sub> grains with different morphology. The ability to grow on diverse substrates makes it feasible to study the NbSe<sub>2</sub>-substrate interaction and related properties.



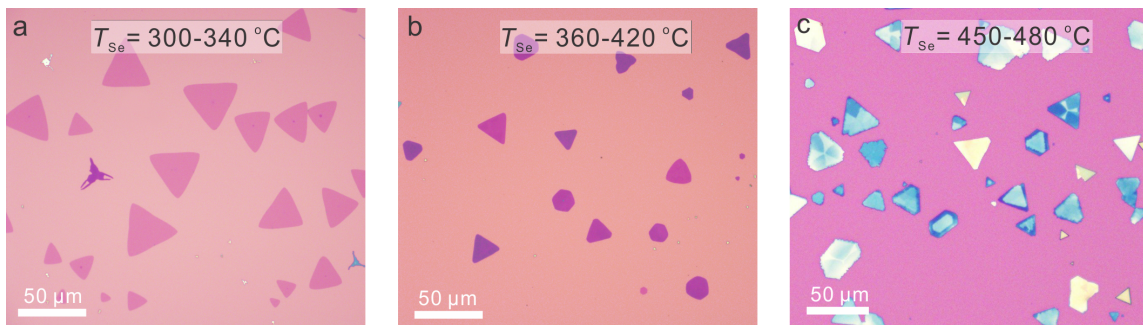
**Supplementary Figure 4.** Morphology and Raman spectrum of NbSe<sub>2</sub> grown on graphene substrates. (a) Optical image of monolayer and few-layer NbSe<sub>2</sub> grown on monolayer CVD graphene/Si substrate. NbSe<sub>2</sub> grows into irregular shapes on graphene. (b) Raman spectrum of monolayer NbSe<sub>2</sub> on monolayer graphene, showing A<sub>1g</sub> and E<sub>2g</sub> modes of NbSe<sub>2</sub>.



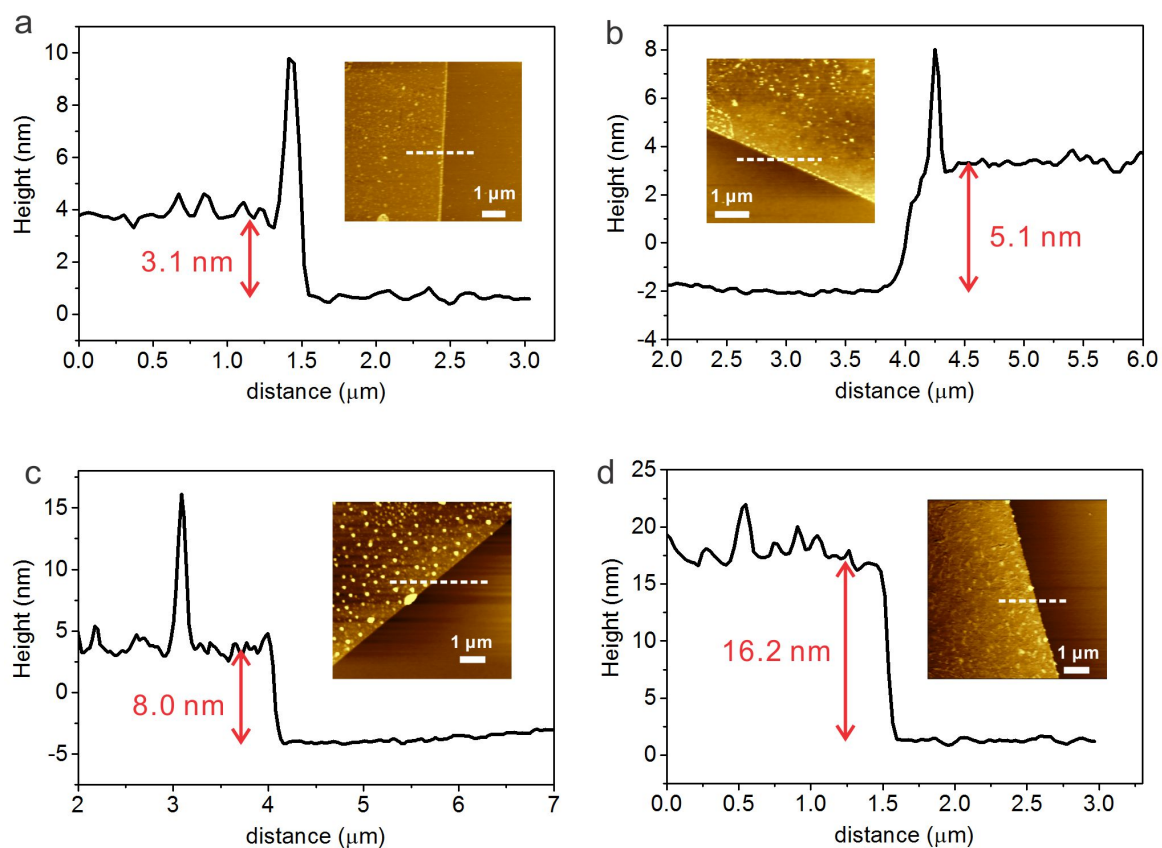
**Supplementary Figure 5.** XPS survey spectrum of few-layer NbSe<sub>2</sub> deposited on SiO<sub>2</sub>/Si substrates



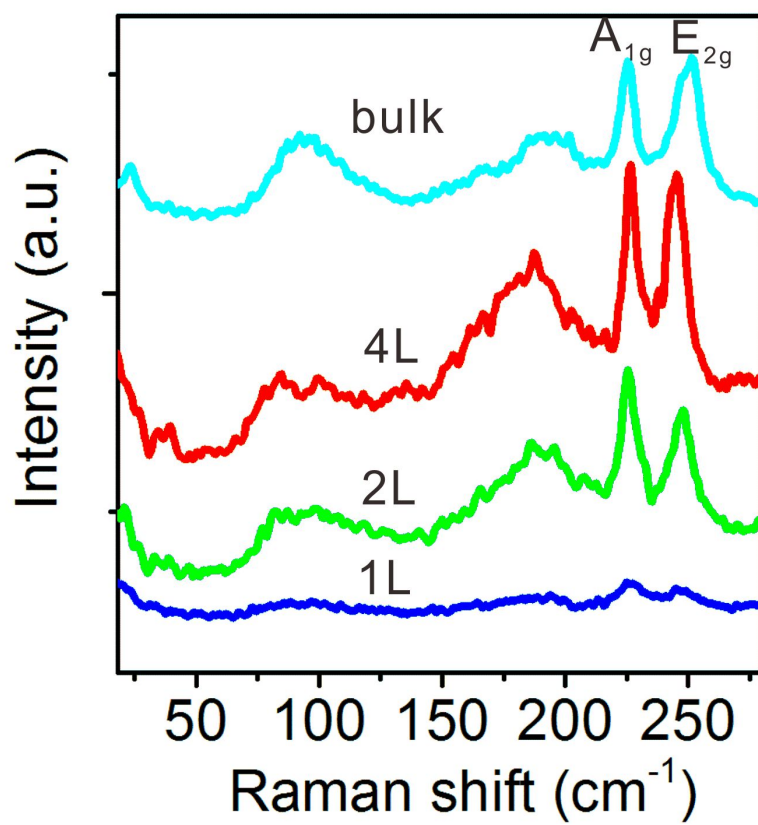
**Supplementary Figure 6** XPS characterization of Na and Cl residues in the as-grown NbSe<sub>2</sub>. (a) Na 1s and (b) Cl 2p XPS spectra collected on NbSe<sub>2</sub> samples. For comparison, standard Na 1s and Cl 2p line positions are shown in the figure. No Na 1s and Cl 2p peaks are observed in the XPS spectra, indicating as-grown NbSe<sub>2</sub> sample are free of Na and Cl contamination.



**Supplementary Figure 7.** Representative morphologies of NbSe<sub>2</sub> crystals synthesized with  $T_{Se}$  setting at (a) 300-340 °C, (b) 360-420 °C and (c) 450-480 °C, respectively.

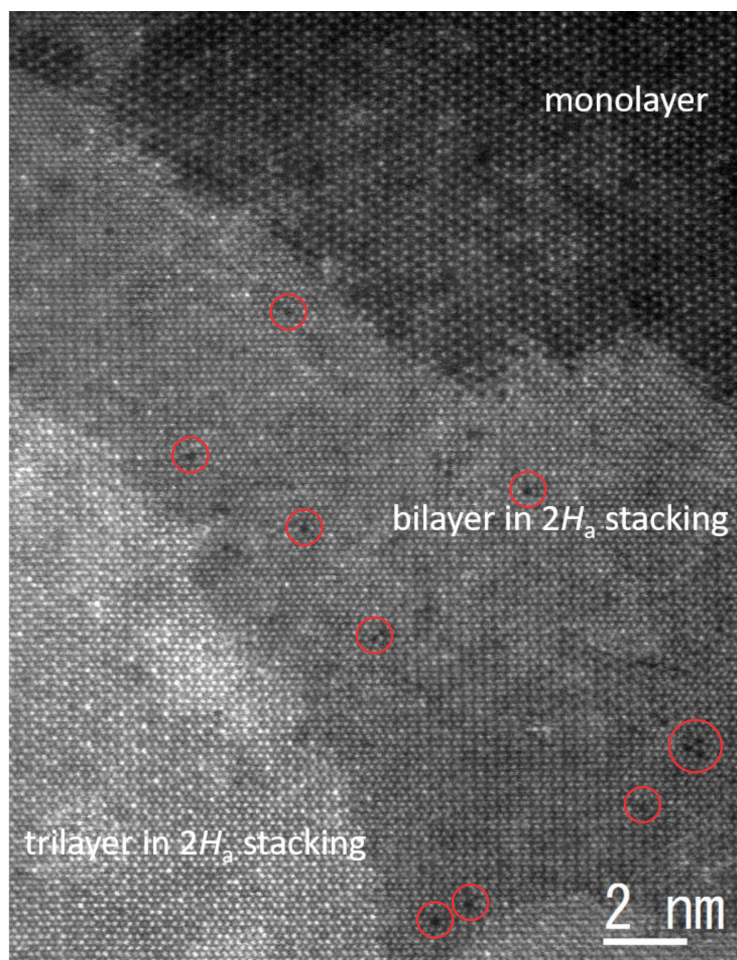


**Supplementary Figure 8.** AFM height profiles and images (inset) of NbSe<sub>2</sub> with different thickness.

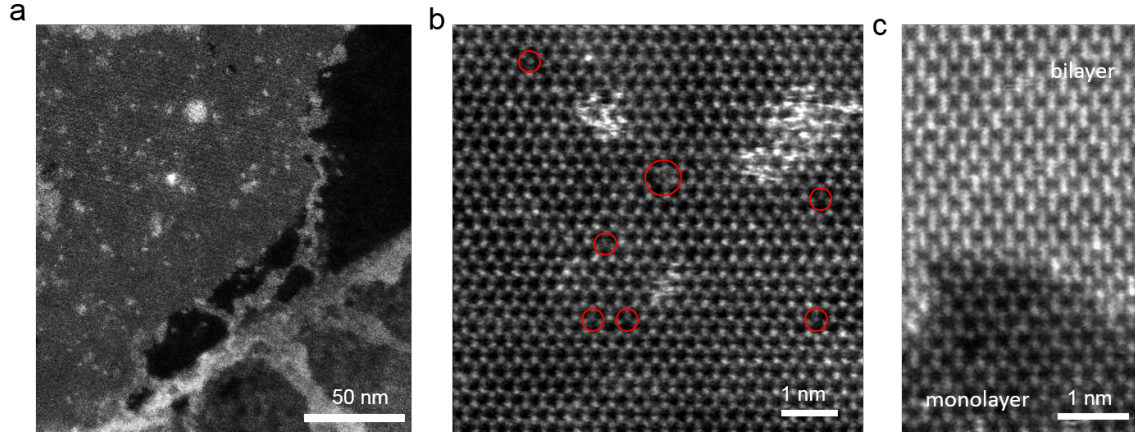


**Supplementary Figure 9.** Room-temperature Raman spectra of NbSe<sub>2</sub> with different thickness.





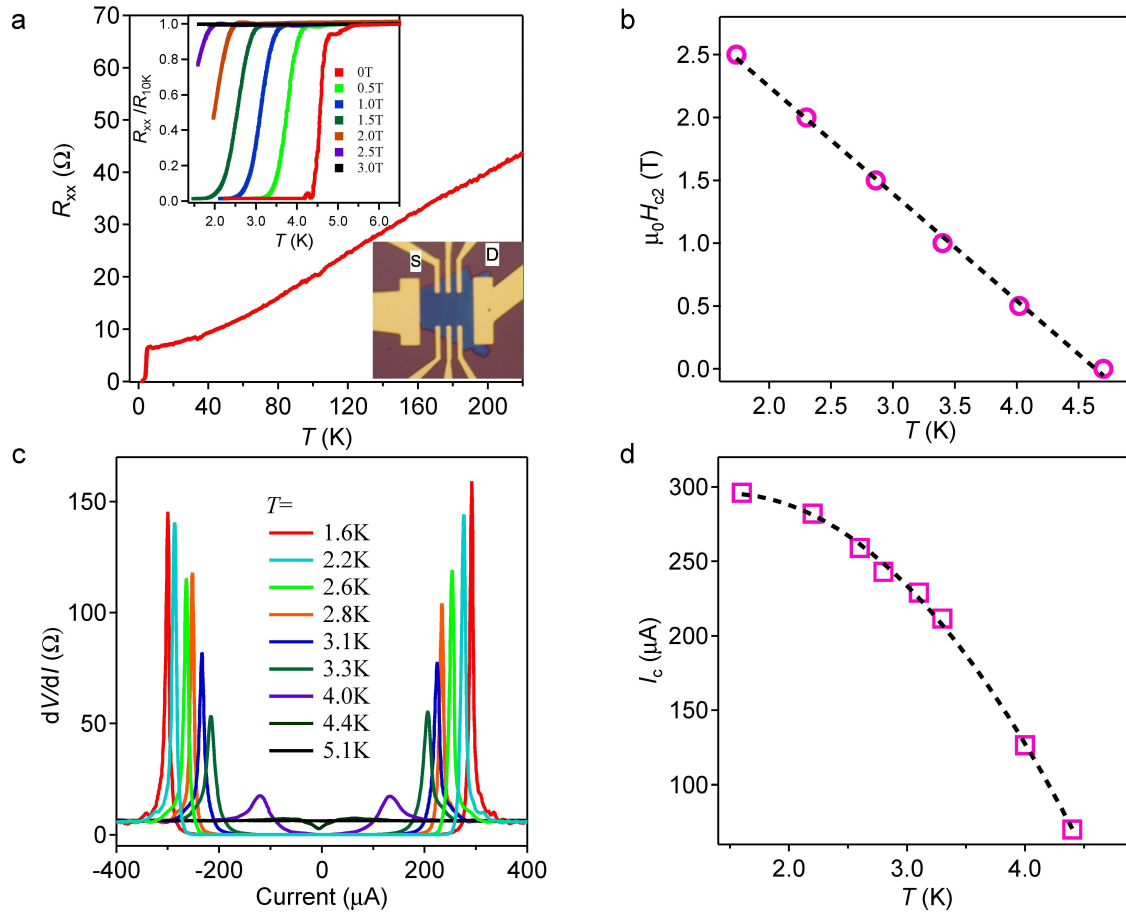
**Supplementary Figure 10.** ADF-STEM image of the multiple layer region of NbSe<sub>2</sub>. The stacking of the CVD-grown sample is found to be mostly in 2H<sub>a</sub> stacking, while a few regions show 2H<sub>c</sub> stacking. The major defect types in multiple layer NbSe<sub>2</sub> are Se vacancies (indicated by red circles), similar to the case in monolayer NbSe<sub>2</sub>.



**Supplementary Figure 11.** ADF-STEM images of monolayer NbSe<sub>2</sub> grown directly on graphene. (a) Low magnified ADF-STEM image showing a large region of monolayer NbSe<sub>2</sub>. (b) Atomic resolution ADF-STEM image of the high-quality hexagonal NbSe<sub>2</sub> lattice. Point defects are indicated by the red circles. Most of the defects are Se vacancy, similar to the case grown on SiO<sub>2</sub>/Si substrate. (c) Atomic resolution ADF-STEM image showing a step edge of the bilayer NbSe<sub>2</sub> film with the 2H<sub>c</sub> stacking.

### Supplementary Note 1: Superconductivity of few-layer NbSe<sub>2</sub> crystals

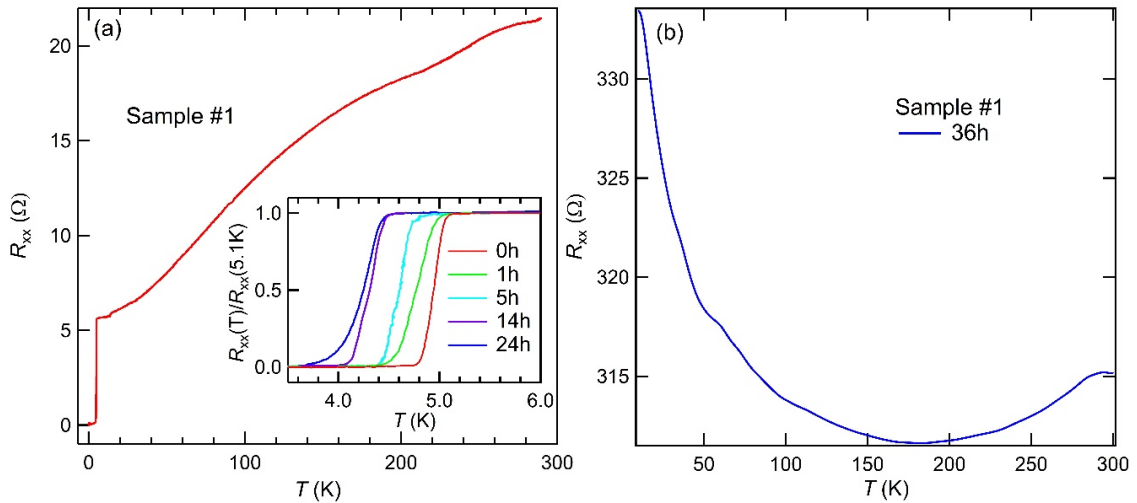
The superconductivity was also observed in few-layer NbSe<sub>2</sub> crystals. Supplementary Figure 12a shows the superconductivity observed on Sample B with the 10-layer NbSe<sub>2</sub>. Similar to that observed in Sample A in the main text, the sample shows a metallic behavior at high temperatures. With the temperature further reduced,  $R_{xx}$  drops at onset temperature  $T_{\text{onset}} = 5.5\text{K}$  and reaches zero resistance at  $T_{\text{zero}} = 4.2\text{K}$ . Compared to the thinner sample (Sample A),  $T_{\text{zero}}$  increases from 0.8K to 4.2K. From the  $H_{c2}$ - $T_c$  phase diagram shown in Supplementary Figure 12b, the Ginzburg-Landau coherence length  $\xi_{\text{GL}}$  is estimated to be about 18nm. Supplementary Figure 12c shows the differential resistance  $dV/dI$  in zero magnetic field as a function of the bias current at different temperatures. The peaks in the  $dV/dI$  curve indicate the transition from the superconducting to the normal state. The peaks shift to zero with increasing temperature, and finally disappear at 5.1K. Supplementary Figure 12d shows the temperature dependence of the critical current  $I_c$  derived from Supplementary Figure 12c.



**Supplementary Figure 12.** Transport measurements on Sample B of a 10-layer NbSe<sub>2</sub> device. (a) Temperature dependence of the longitudinal resistance  $R_{xx}$  in zero magnetic field from 300K to 0.26K. Upper left inset: Superconductivity in different magnetic fields. Lower right inset: Optical image of a typical 10-layer NbSe<sub>2</sub> device. (b) Temperature dependence of the upper critical field  $H_{c2}$ . The dashed line is the linear fit to  $H_{c2}$ . (c) Differential resistance in zero magnetic field as a function of the bias current at different temperatures. (d) Temperature dependence of the critical current. The dashed line is a guide to the eye.

## Supplementary Note 2: Influence of oxygen and moisture on the superconductivity of few-layer NbSe<sub>2</sub>

The influence of oxygen and moisture on the CVD-grown NbSe<sub>2</sub> was carried out by exposing samples (#1 and #2) to ambient conditions with different time followed immediately by measuring their transport properties. Supplementary Figure 13 shows the typical experimental results collected on Sample #1. To characterize the changes of sample quality, the residual resistance ratio  $RRR$ , the critical temperature  $T_c$  ( $0.5R_N$ ), and the superconducting transition width  $\Delta T_c$  are summarized in Supplementary Table 1. As the exposure time increases from 0h to 24h,  $RRR$  decreases from 3.82 to 3.56 and  $T_c$  ( $0.5R_N$ ) reduces from 4.94K to 4.21K. Correspondingly, the superconducting transition broadens gradually from 0.37K to 1.05K. For Sample #1, the superconductivity can still be observed when it was exposed to air for 24 hours. However, when the exposure time increases to 36 hours, no superconductivity can be detected and a metal-insulator transition emerges around 170K as shown in Supplementary Figure 13b. No conductive behavior was found when the sample was further exposed to air. From Supplementary Table 1, Sample #2 seems to be much more stable and the superconductivity can survive with 36 hours exposure. These data indicate that, even the NbSe<sub>2</sub> samples are capped with graphene, they are still sensitive to oxygen and moisture. The ambient degrades the samples with increasing exposure time.

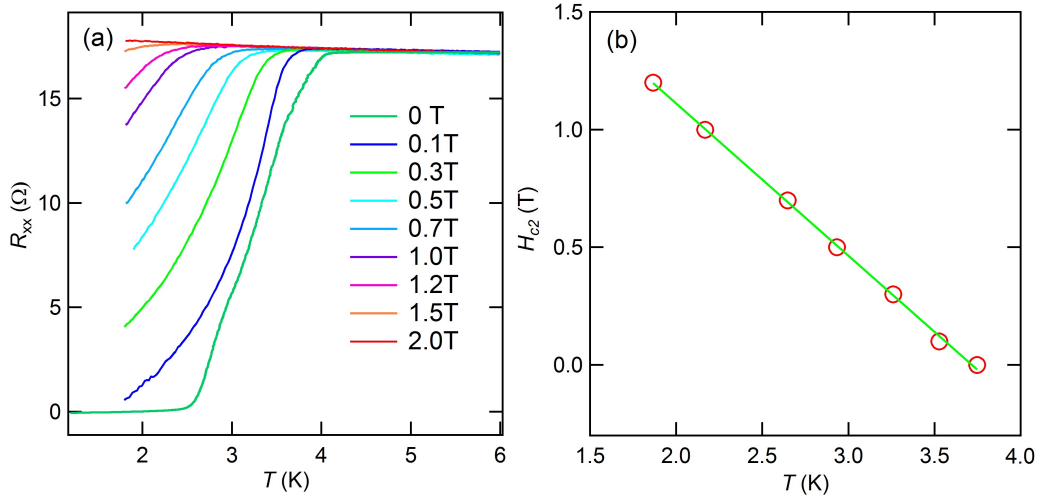


**Supplementary Figure 13.** Ambient-exposure effects of superconductivity in CVD-grown NbSe<sub>2</sub>. (a) Temperature dependence of resistivity for NbSe<sub>2</sub> exposed with different time from 0h to 24h (indicated by different colors). (b)  $R-T$  plot for the same NbSe<sub>2</sub> sample after 36h exposure to ambient.

**Supplementary Table 1.** Residual resistance ratio  $RRR$ , critical temperature  $T_c$  ( $0.5R_N$ ), and superconducting transition width  $\Delta T_c$  for CVD-grown few-layer NbSe<sub>2</sub> samples exposed to ambient

Sample	Ambient exposure time	$RRR$	$T_c$ (K)	$\Delta T_c$ (K)
Sample #1	0h	3.82	4.94	0.37
	1h	3.63	4.74	0.64
	5h	-	4.60	0.58
	14h	3.46	4.30	0.67
	24h	3.56	4.21	1.05
Sample #2	0h	3.15	3.25	0.70
	1h	3.14	3.25	0.65
	5h	3.15	3.24	0.50
	14h	2.95	3.0	0.42
	19h	2.81	2.82	1.42
	36h	2.57	2.73	1.73

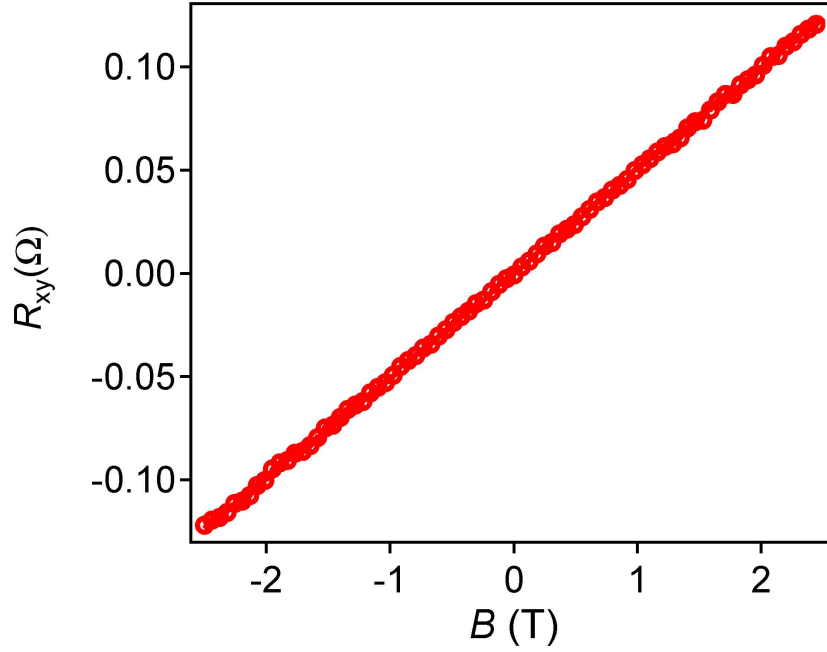
### Supplementary Note 3: Superconductivity of few-layer NbSe<sub>2</sub> grown on graphene substrate



**Supplementary Figure 14.** Superconductivity in 8-layer thick NbSe<sub>2</sub> grown on Graphene substrate. (a) Temperature dependence of the resistance under different magnetic fields. (b) Temperature dependence of the upper critical field  $H_{c2}$ . The solid line is a linear fit to  $H_{c2}$ .

The superconductivity of few-layer NbSe<sub>2</sub> grown on graphene substrate was also investigated. The transport data of an 8-layer NbSe<sub>2</sub> sample grown on graphene is shown in Supplementary Figure 14. At zero magnetic field, the superconducting transition critical temperature  $T_c$  of an 8-layer NbSe<sub>2</sub> on graphene is  $\sim 3.2$ K. The superconducting transition  $\Delta T_c$  is about 1.5K, which is wider than the few-layer (5-layer and 10-layer) NbSe<sub>2</sub> grown on SiO<sub>2</sub>/Si substrate. These observations could be caused by the presence of PMMA residues on transferred CVD graphene film. The residues on graphene may act as nucleation centers for NbSe<sub>2</sub> growth. Therefore, more defects may exist in the NbSe<sub>2</sub> crystals grown on graphene substrates, resulting in poorer superconducting properties.

**Supplementary Note 4: Estimation of carrier density and mean free path for monolayer superconducting NbSe<sub>2</sub>**



**Supplementary Figure 15.** Hall resistance  $R_{xy}$  measured at 2.0 K for a monolayer NbSe<sub>2</sub> crystal. Measurements were performed in the normal state with a current bias of  $1 \mu\text{A}$  on the crystals while the magnetic field was applied perpendicularly to the sample plane. The carrier density and mean free path were estimated to be  $1.25 \times 10^{16} \text{ cm}^{-2}$  and 1.3 nm by using the following procedure.

**1) Estimation of carrier density  $n_s$**

By fitting the experimental data in Supplementary Figure 15, we can calculate the carrier density  $n_s$  via the following formula:

$$n_s = \frac{I/e}{dV_H/dB} = \frac{1/e}{dR_H/dB} = 1.25 \times 10^{16} \text{ cm}^{-2} \quad (1)$$

**2) Estimation of sheet resistance  $R_s$**

The width  $w$  and length  $l$  for the measured sample can be determined by the optical image. The normal state resistance  $R_N = 242.3 \text{ } \Omega$  was taken at  $T = 1.5 \text{ K}$ . Then the sheet resistance  $R_s$  can be calculated via the following formula:

$$R_s = \frac{R_N w}{l} = 727 \text{ } \Omega \text{ } \square^{-1} \quad (2)$$

### 3) Estimation of mobility $\mu$

According to Drude model, the mobility  $\mu$  can be calculated via the following formula:

$$\mu = \frac{1}{e R_s n_s} = 0.69 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1} \quad (3)$$

### 4) Estimation of the Fermi velocity $v_F$

In two-dimensional systems, the Fermi velocity  $v_F$  can be calculated via the following formula:

$$v_F = \frac{\hbar}{m} \sqrt{2\pi n_s} = 3.25 \times 10^6 \text{ m s}^{-1} \quad (4)$$

### 5) Estimation of the momentum scattering time $\tau_m$

$$\tau_m = \frac{\mu m}{e} = 3.9 \times 10^{-16} \text{ s} \quad (5)$$

### 6) Estimation of the mean free path $l_m$

$$l_m = v_F \tau_m = 1.3 \text{ nm} \quad (6)$$

### 7) Estimation of the Fermi wave vector $k_F$

$$k_F = \sqrt{2\pi n_s} = 2.8 \times 10^8 \text{ cm}^{-1} \quad (7)$$

### 8) Estimation of the Ioffe-Regel criterion

$$k_F l_m = 37 \quad (8)$$



## Supplementary Note 5: Comparison of NbSe<sub>2</sub> and some high- $T_c$ superconductors

It is instructive to consider together NbSe<sub>2</sub> and some high- $T_c$  cuprates and Fe-based superconductors (FeSCs)<sup>1-5</sup>. Structurally, while they all crystallize in layered structures, they possess different superconductively active layers including NbSe<sub>2</sub> layer in NbSe<sub>2</sub>, CuO<sub>2</sub> planes in cuprates and Fe<sub>2</sub>As<sub>2</sub>/Fe<sub>2</sub>Se<sub>2</sub> layers in FeSCs. While un-doped NbSe<sub>2</sub> is a superconductor, the superconductivity in cuprates and FeSCs emerges upon doping carriers into antiferromagnetic parent compounds<sup>2</sup>. Despite their diversity in structure, the phase diagrams of these materials share the common characteristic that superconductivity exists near the boundary of an ordered phase (CDW in NbSe<sub>2</sub><sup>6</sup>, spin/charge density wave in cuprates and iron arsenides<sup>1, 3, 7</sup>). The presence of various nearby magnetic and other competing orders in cuprates and FeSCs makes it difficult to understand the relationship between ordered phases and superconductivity<sup>7, 8</sup>. By comparison, NbSe<sub>2</sub> is a much simple model, in which the competitive CDW and superconductivity order is believed to be induced by pure electron–phonon coupling<sup>8</sup>. Therefore, the research on NbSe<sub>2</sub> may provide insight into the understanding of high- $T_c$  superconductivity.

On the other hand, the 2D superconductivity in high- $T_c$  superconductors is seldom studied due to the difficulty in ultrathin sample preparation. FeSe as the simplest FeSCs has been successfully synthesized by MBE. The most striking finding in monolayer FeSe is the significant enhancement of  $T_c$  when grown on a SrTiO<sub>3</sub> substrate<sup>9, 10</sup>. However, the superconductor-substrate interaction in NbSe<sub>2</sub> system has not been investigated. With our proposed CVD technology, we grew NbSe<sub>2</sub> layers on graphene substrates and the transport measurement indicates a decrease of  $T_c$  compared with SiO<sub>2</sub>/Si substrate (see Supplementary Note 3). This phenomenon indicates that the substrate also plays important role in determining the superconductivity of ultrathin NbSe<sub>2</sub>. Based on the developed CVD method, more investigations could be done to understand the 2D superconductor-substrate interplay, as well as to explore the unexpected and novel superconductivity of NbSe<sub>2</sub> on specific substrates.

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