Science Advances

advances.sciencemag.org/cgi/content/full/4/4/eaaq1636/DC1

NAAAS

Supplementary Materials for

Readily accessible shape-memory effect in a porous interpenetrated coordination network

Mohana Shivanna, Qing-Yuan Yang, Alankriti Bajpai, Susan Sen, Nobuhiko Hosono, Shinpei Kusaka, Tony Pham, Katherine A. Forrest, Brian Space, Susumu Kitagawa, Michael J. Zaworotko

> Published 27 April 2018, *Sci. Adv.* **4**, eaaq1636 (2018) DOI: 10.1126/sciadv.aaq1636

This PDF file includes:

- fig. S1. Structure of the three phases.
- fig. S2. Diffraction pattern and single-crystal images.
- fig. S3. Low-pressure gas sorption.
- fig. S4. Recyclability of shape-memory phase at 195 K $CO₂$.
- fig. S5. High-pressure $CO₂$ sorption.
- fig. S6. In situ variable temperature PXRD.
- fig. S7. Solvent-induced phase change.
- fig. S8. PXRD data for **X-pcu-3-Zn-3i-γ** obtained from different experiments.
- fig. S9. Phase change from γ to α phase.
- fig. S10. Variable temperature PXRD of shape-memory phase.
- fig. S11. Distortion of paddlewheel MBB and orientation of X-ligands in **X-pcu-3-Zn-3i-α**.
- fig. S12. Distortion of paddlewheel MBB and orientation of X ligands in **X-pcu-3-Zn-3i-β**.
- fig. S13. Distortion of paddlewheel MBB and orientation of X ligands in **X-pcu-3-Zn-3i-γ**.
- fig. S14. Thermogravimetric analysis (TGA) profiles.
- fig. S15. Comparison of modeling and experimental isotherm.
- fig. S16. Free-energy difference.
- fig. S17. Chemically distinct atoms in **X-pcu-3-Zn-3i-α**.
- fig. S18. The chemically distinct atoms in **X-pcu-3-Zn-3i-β**.
- fig. S19. The chemically distinct atoms in **X-pcu-3-Zn-3i-γ**.
- table S1. Examples of previously reported FMOMs with large hysteresis.
- table S2. Crystal data and refinement parameters.
- table S3. Numerical labeling of atoms corresponds to fig. S17.
- table S4. The crystallographic distances (in Å) between various atoms in **X-pcu-3-Zn-3i-α**.
- table S5. The partial charges (in e^-) for the chemically distinct atoms in **X-pcu-3 Zn-3i-β**.
- table S6. The crystallographic distances (in Å) between various atoms in **X-pcu-3-Zn-3i-β**.
- table S7. Numerical labeling of atoms corresponds to fig. S18.
- table S8. Labels of atoms correspond to fig. S18.
- table S9. Calculated energies for three phases.
- References $(41–62)$

fig. S1. Structure of the three phases. Relative orientation of the three interpenetrated networks in **X-pcu-3-Zn-3i** shown in red, blue and green.

In-situ Single Crystal-to-Single Crystal (SCSC) transformation studies

fig. S2. Diffraction pattern and single-crystal images. A single crystal undergoes transformation from **X-pcu-3-Zn-3i-α** to **X-pcu-3-Zn-3i-β** at 120 °C in SC-SC fashion.

Gas sorption studies

fig. S3. Low-pressure gas sorption. Low temperature gas sorption on **X-pcu-3-Zn-3i.** First cycle of **X-pcu-3-Zn-3i-β** exhibits stepped isotherm and then it transforms to shape memory phase, **Xpcu-3-Zn-3i-γ** followed of N₂ adsorption isotherms showed a type I profile. a) 195 K CO₂, adsorption (closed blue circles) and desorption (open blue circles), 77 K N_2 adsorption (closed red circles) and desorption (open red circles). The BET surface area (S*BET*) of X-pcu-3-Zn-3i-γ, the S_{BET} was found to be 800 m^2g^{-1} calculated from 77 K N₂ sorption.

fig. S4. Recyclability of shape-memory phase at 195 K CO2. (**A**) First cycle of 195 K CO² sorption. (**B**) Sample obtained after the first cycle was regenerated (130 °C, 2 h *in vacuo*) and second cycle of 195 K CO² sorption was measured**.**

fig. S5. High-pressure CO² sorption. Three consecutive cycles at 288 K, 298 K and 305 K after first cycle**.**

In-situ variable temperature PXRD

fig. S6. In situ variable temperature PXRD. X-pcu-3-Zn-3i-α undergoes phase transformation at 120 °C.

PXRD studies: interconversion of phases

fig. S7. Solvent-induced phase change. PXRD: conversion of **X-pcu-3-Zn-3i-β** to **X-pcu-3-Zn-3i-α** after soaking in DMF.

fig. S8. PXRD data for X-pcu-3-Zn-3i-γ obtained from different experiments.

fig. S9. Phase change from γ to α phase. Comparison of PXRD patterns shows that **X-pcu-3- Zn-3i-γ** reverts to **X-pcu-3-Zn-3i-α** only in soaking in DMF**.**

fig. S10. Variable temperature PXRD of shape-memory phase. PXRD patterns show that **Xpcu-3-Zn-3i-γ** retained on heating up to 240°C.

Distortion of paddlewheel MBB: bond angles

A centroid (C) was created for the paddlewheel MBB. Bond angles were measured for each linker through this centroid.

fig. S11. Distortion of paddlewheel MBB and orientation of X-ligands in X-pcu-3-Zn-3i-α.

fig. S12. Distortion of paddlewheel MBB and orientation of X ligands in X-pcu-3-Zn-3i-β.

fig. S13. Distortion of paddlewheel MBB and orientation of X ligands in X-pcu-3-Zn-3i-γ.

Thermogravimetric analyses (TGA)

fig. S14. Thermogravimetric analysis (TGA) profiles. (**A**) as-synthesized **X-pcu-3-Zn-3i-α**, 24% weight loss corresponds to ~6 DMF. (**B**) **X-pcu-3-Zn-3i-γ** obtained after 195 K CO² sorption. (**C**) **X-pcu-3-Zn-3i-β** obtained after heating **X-pcu-3-Zn-3i-α** at 130 °C for 12 h. (**D**) **Xpcu-3-Zn-3i-β** obtained after washing (solvent exchange) **X-pcu-3-Zn-3i-α** with MeCN.

Molecular modelling

fig. S15. Comparison of modeling and experimental isotherm. Modelling experiments for both beta and gamma at 195K CO₂ is very well agreement with experimental gas sorption data.

Estimation of free-energy difference (0*F***host) between X-pcu-3-Zn-3i-α and X-pcu-3-Zn-3i-β for CO² adsorption and desorption.**

$$
\Delta F = \begin{vmatrix} 31.75893195 & kJ/mol \end{vmatrix}
$$

fig. S16. Free-energy difference. Langmuir-Freundlich (LF) fittings of the CO₂ adsorption (A) and desorption (**B**) isotherms for **X-pcu-3-Zn-3i-β** at 195 K.

A. Grand Canonical Monte Carlo

Simulations of CO₂ adsorption were performed in the three different phases (α , **β**, and γ) of **Xpcu-3-Zn-3i** using grand canonical Monte Carlo (GCMC) methods *(50)*. The single X-ray crystallographic structures that were obtained for all three phases herein were used for the parametrizations and simulations. All GCMC simulations were performed within the rigid crystal structures of the three individual phases. System cell dimension lengths of $2 \times 2 \times 1$, $2 \times 2 \times 2$, and $2 \times 1 \times 2$ were used for **X-pcu-3-Zn-3i-α**, **X-pcu-3-Zn-3i-β**, and **X-pcu-3-Zn-3i-γ**, respectively. A spherical cut-off distance corresponding to half the shortest system cell dimension length was used for the simulations in all three phases.

A five-site polarizable potential that was developed previously for $CO₂$ was used for the simulations in this work (51) . The total potential energy of the MOM–CO₂ system was calculated through the sum of the repulsion and dispersion, permanent electrostatic, and polarization energies. These were calculated using the Lennard-Jones 12–6 potential, partial charges with Ewald summation, and a Thole-Applequist type model *(52)*, respectively. The chemical potential for CO² was determined for a range of pressures and temperatures through the Peng-Robinson equation of state *(53)*. All simulations were performed using the Massively Parallel Monte Carlo (MPMC) code (54). For all state points considered, the simulations consisted of 5.0×10^6 Monte Carlo steps to guarantee equilibration, followed by an additional 5.0×10^6 steps to sample the desired thermodynamic properties.

B. Parametrization

All atoms within the three different phases of **X-pcu-3-Zn-3i** were treated with Lennard-Jones 12–6 parameters, atomic point partial charges, and scalar point polarizabilities to model repulsion and dispersion, permanent electrostatic, and polarization interactions, respectively. The Lennard-Jones parameters (ε and σ) for all aromatic C, H, and N atoms were taken from the Optimized Potentials For Liquid Simulations – All Atom (OPLS-AA) force field *(55)*, while such parameters for Zn and O were taken from the Universal Force Field (UFF) *(56)*.

Examination of the crystal structure of the **α**, **β**, and **γ** phases of **X-pcu-3-Zn-3i** revealed 126, 42, and 42 atoms in chemically distinct environments, respectively (**Figures S17–S19**). The partial charges for the unique atoms in these phases were determined through electronic structure calculations on a variety of clusters that were extracted from the crystal structure of the respective phases. For these calculations, all light atoms (C, H, N, and O) were treated with the 6-31G^{*} basis set, while the LANL2DZ ECP basis set (57) was used for Zn^{2+} . The NWChem *ab initio* software *(58)* was used to calculate the electrostatic potential surface for each fragment and the partial charges were subsequently fitted onto the atomic positions of the fragments using the CHELPG method *(59)*. For each chemically distinct atom, the partial charges were averaged between the fragments. The partial charges were then adjusted such that the total charge of the system was equal to zero. The final tabulated partial charges for each chemically distinct atom in **X-pcu-3- Zn-3i-α**, **X-pcu-3-Zn-3i-β**, and **X-pcu-3-Zn-3i-γ** are provided in **Tables S3**, **S5**, and **S7**, respectively. Note, the crystallographic distances between various chemically distinct atoms for the individual phases are provided in **Tables S4**, **S6**, and **S8**.

The exponential damping-type atomic point polarizabilities for all C, H, N, and O atoms were taken from a carefully parametrized set provided by the work of van Duijnen and Swart ($C =$ 1.28860 \AA^3 , H = 0.41380 \AA^3 , N = 0.97157 \AA^3 , O = 0.85200 \AA^3) (60). The polarizability parameter for Zn^{2+} was determined in previous work ($\text{Zn}^{2+} = 1.98870 \text{ Å}^3$) *(61)* and was used herein. These polarizability values were assigned to the nuclear centers of all atoms of the individual phases to model explicit polarization.

C. Density Functional Theory

The relative energies of the three different phases of **X-pcu-3-Zn-3i** were evaluated through periodic density functional theory (DFT) calculations that were implemented with the Vienna *ab initio* Simulation Package (VASP) *(62)*. These calculations were performed using the projector augmented wave (PAW) method with the Perdew–Burke–Ernzerhof (PBE) functional. In addition, dispersion effects were treated using the DFT-D3 correction method of Grimme *et al*.

Dispersion effects are expected to be important in the crystal structure of these phases since nonbonded portions of these structures are in close proximity to each other. The crystal structure of each phase of **X-pcu-3-Zn-3i** has varying number of atoms within its unit cell due to increasing or decreasing symmetry after each transition. Therefore, the calculated energies for the unit cell of the **β** and **γ** phases were subjected to a scalar (3 and 3/4, respectively) in order for these energies to be properly compared with that for the **α** phase.

The calculated energies for the three different phases of **X-pcu-3-Zn-3i** are displayed in **Table S7**. The calculations suggest that the **α** phase is the most stable phase of **X-pcu-3-Zn-3i** even though it is only representative of the material prior to activation. After activation, the material transitions to the **β** phase, which is the next most stable phase according to our calculations. After **X-pcu-3-Zn-3i-β** achieves CO₂ saturation, the material then transitions to the γ phase, which has the lowest energy of the three phases. Since the α phase is not accessed prior to or during $CO₂$ adsorption, it can be speculated that a prohibitive energy barrier exists between the **α** and **β** phases, which is catalyzed by the activation process to achieve the **β** phase, but not surmountable under successively applied conditions.

fig. S17. Chemically distinct atoms in X-pcu-3-Zn-3i-α. defining the numbering system corresponding to **Tables S3** and **S4**. Atom colors: $C = cyan$, $H = white$, $N = blue$, $O = red$, $Zn =$ silver.

fig. S18. The chemically distinct atoms in X-pcu-3-Zn-3i-β. defining the numbering system corresponding to **Tables S5** and **S6**. Atom colors: $C = cyan$, $H = white$, $N = blue$, $O = red$, $Zn =$ silver.

fig. S19. The chemically distinct atoms in X-pcu-3-Zn-3i-γ. Defining the numbering system corresponding to **Tables S7** and **S8**. Atom colors: $C = cyan$, $H = white$, $N = blue$, $O = red$, $Zn =$ silver.

table S1. Examples of previously reported FMOMs with large hysteresis.

table S2. Crystal data and refinement parameters.

 \mathcal{L}

table S5. The partial charges (in *e* **−) for the chemically distinct atoms in X-pcu-3-Zn-3i-β.** Label of atoms correspond to **fig. S18**.

table S8. Labels of atoms correspond to fig. S18. The crystallographic distances (in Å) between various atoms in **X-pcu-3-Zn-3i-γ**. **Atom Pair Distance (Å)** 1–2 2.03074 $1-17$ 2.00319 1–18 2.06310

table S9. Calculated energies for three phases. Calculated energies (in eV) for the crystal structure of the three different phases (**α**, **β**, and **γ**) of **X-pcu-3-Zn-3i** as determined through periodic DFT calculations using VASP.

The values shown for the **β** and **γ** phases are the result of multiplying the original calculated energy of the unit cell of these phases by 3 and 3/4, respectively.