

Supplementary Information for  
Discovery of coexisting Dirac and triply degenerate magnons in a  
three-dimensional antiferromagnet

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### Supplementary Note 1: Effective Hamiltonian of triply degenerate points

Due to the PT symmetry, all magnon bands in  $\text{Cu}_3\text{TeO}_6$  have twofold degeneracy. The magnonic Hamiltonian can be written as

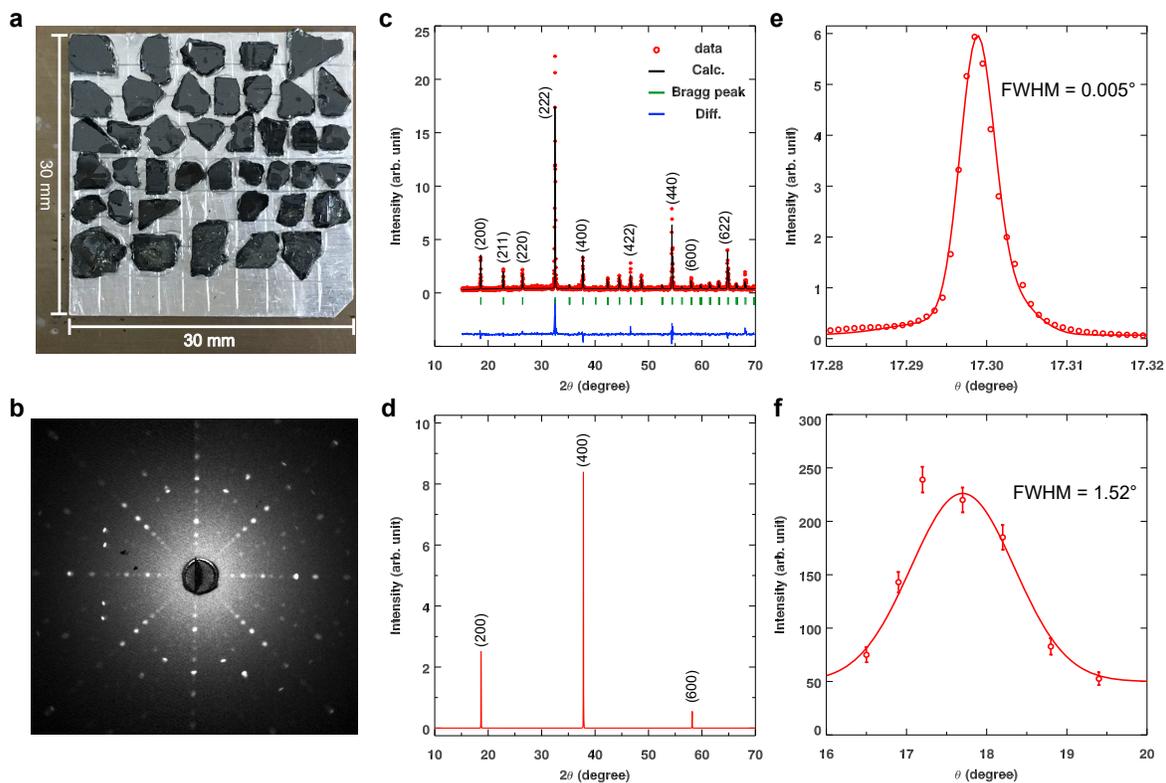
$$H = H_1 \oplus H_1^*. \quad (1)$$

Near the triply-degenerate nodal point, the linearized  $\mathbf{k} \cdot \mathbf{p}$  Hamiltonian takes the form

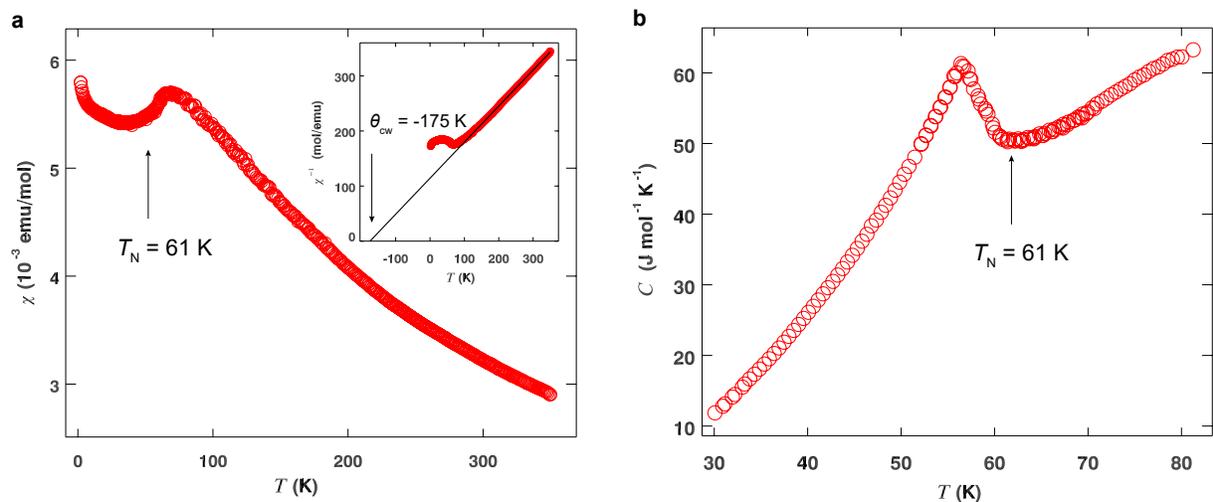
$$H_1(\mathbf{q}) = E_0 \mathbf{I}_{3 \times 3} + \sum_{\alpha=x,y,z} v_\alpha q_\alpha S_\alpha, \quad (2)$$

where  $\mathbf{q} = \mathbf{k} - \mathbf{k}_0$  is the momentum deviation from the crossing point at  $\mathbf{k}_0$ ,  $E_0$  is the energy of the nodal point,  $v_\alpha$  is the group velocity, and  $S_\alpha$  are the  $3 \times 3$  matrix representation of the spin-1 operators. Equation 2 gives an effective description of the magnonic dispersions around the H point shown in Supplementary Figure 6.

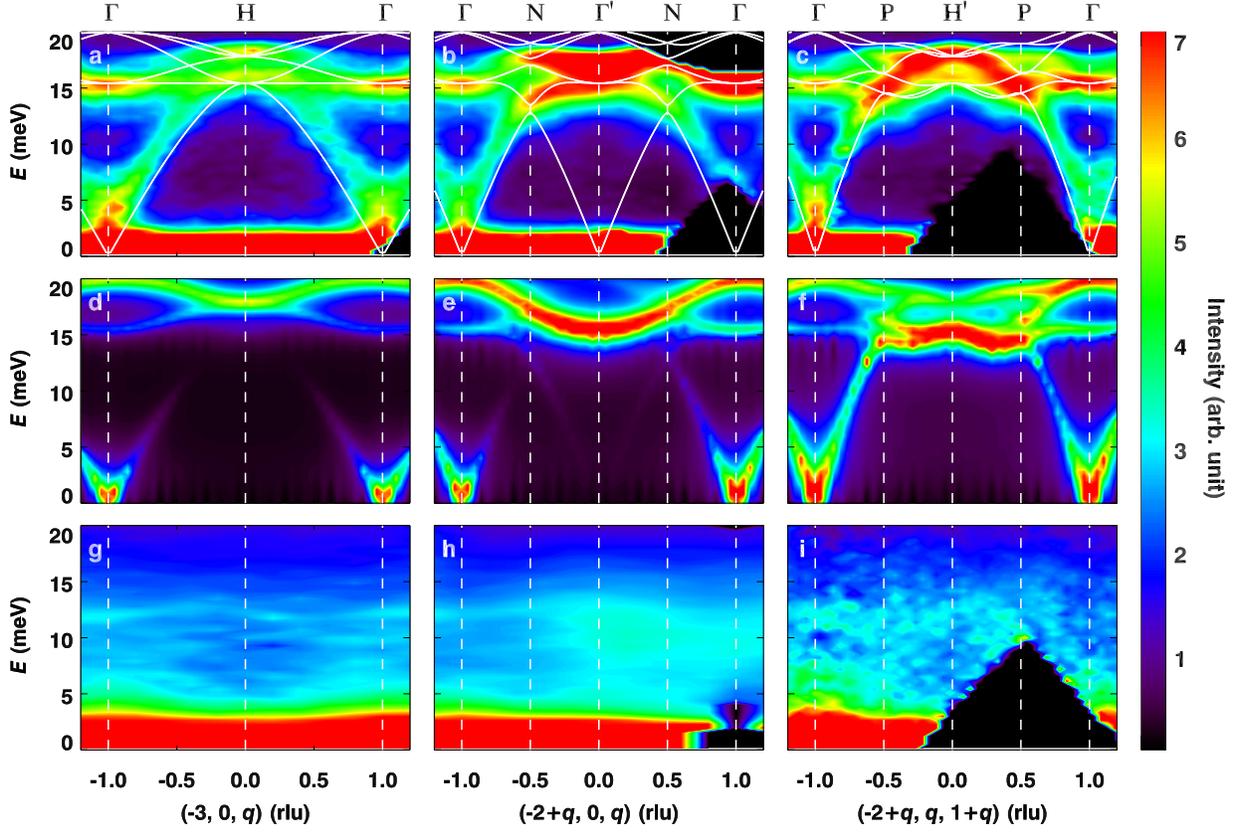
### Supplementary Figures



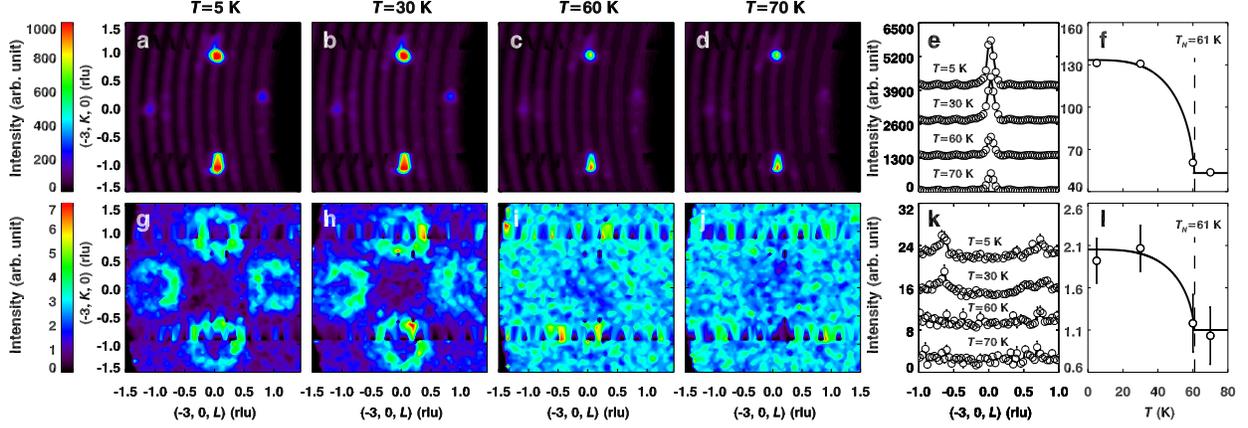
**Supplementary Figure 1.** Single crystals and structural characterisations. **a**, A photograph of the 40 coaligned single crystals of  $\text{Cu}_3\text{TeO}_6$  with a total mass about 3 g on an aluminum plate. The natural cleavage plane is the (100) plane. **b**, Laue X-ray pattern of a single crystal with the beam along the [100] direction. **c**, Powder X-ray diffraction pattern with indices for major Bragg peaks. **d**, Single-crystal X-ray diffraction pattern for the (100) plane. **e** and **f**, Rocking curves of single crystals for the (200) peak measured with X-ray and neutrons, respectively. FWHM represents full width at half maximum. Errors represent one standard deviation throughout the paper.



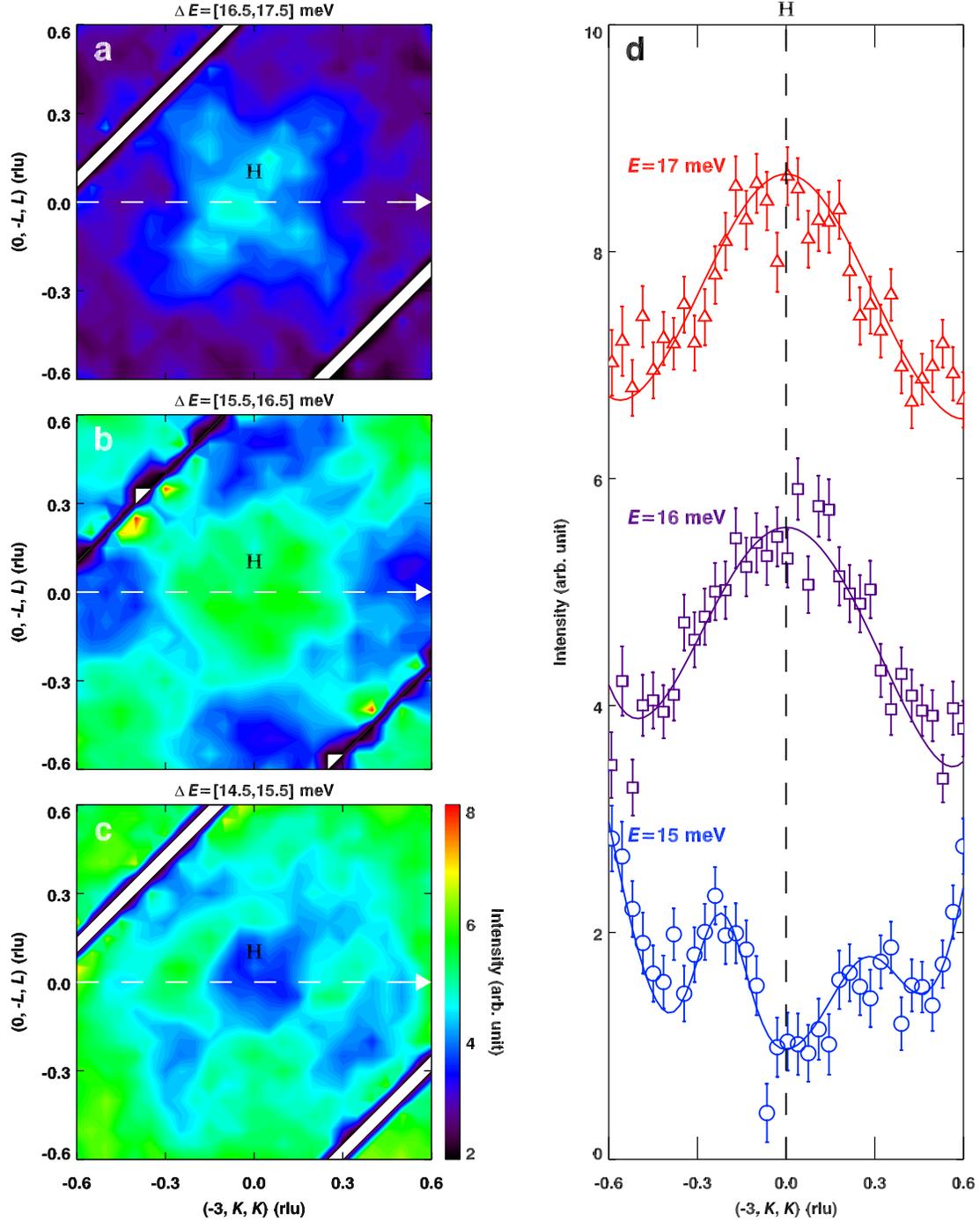
**Supplementary Figure 2.** Magnetic susceptibility and specific heat. **a**, Magnetic susceptibility of a single crystal measured with a magnetic field of 1 T applied along the [100] direction. The inset shows the inverse susceptibility with a Curie-Weiss fit for temperatures ranging from 150 to 350 K, indicated by the solid line.  $\theta_{CW}$  is the Curie-Weiss temperature resulting from the fit. **b**, Specific heat as a function of temperature. The magnetic transition temperature  $T_N \sim 61$  K is indicated in both **a** and **b**.



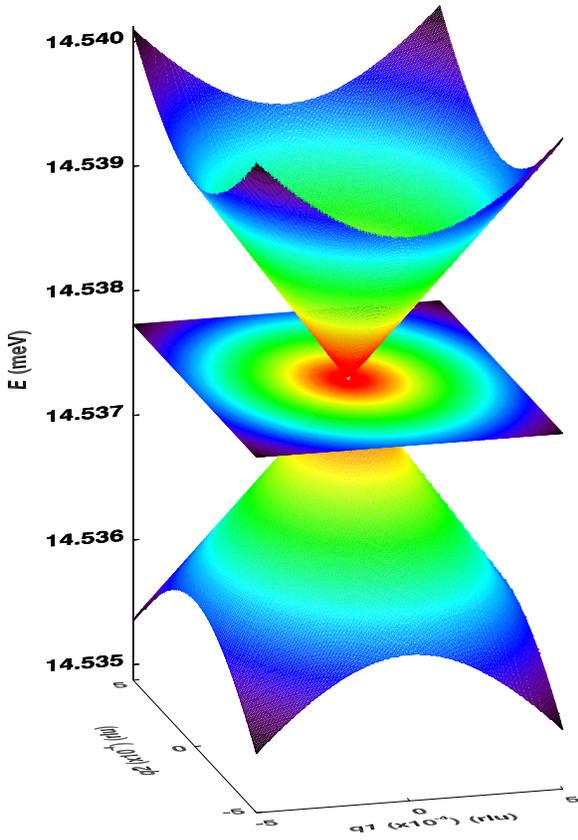
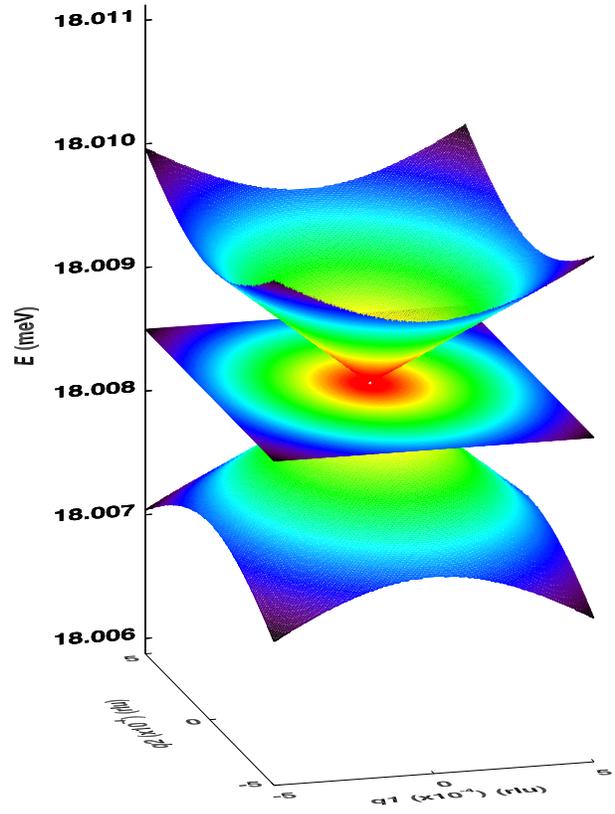
**Supplementary Figure 3.** Experimental and calculated spin-wave excitations. **a-c**, Inelastic neutron scattering results of the spin excitation spectra measured at  $T = 5$  K along [001] (**a**), [101] (**b**), and [111] (**c**) directions, respectively. **d-f**, Calculated magnetic spectra using the linear-spin-wave theory based on a Hamiltonian with nearest-neighbour (NN) exchange interaction  $J_1$  and next NN exchange interaction  $J_2$ . The calculated dispersions are plotted as solid lines in **a-c**. Here,  $J_1$  and  $J_2$  are 10.99 and 1.01 meV, respectively. Vertical dashed lines indicate the  $\mathbf{Q}$  (wave vector) positions illustrated in Fig. 1b in the main manuscript. **g-i**, Same as in **a-c** but measured at  $T = 70$  K.



**Supplementary Figure 4.** Temperature dependence of the magnetic Bragg peaks and spin-wave excitations. **a, b, c,** and **d**, Constant-energy cuts integrated over the energy range  $[-2, 2]$  meV at 5, 30, 60, and 70 K, respectively. **g-j**, Similar as in **a-d** but with energies integrated over  $[8, 10]$  meV. **e**, Constant- $\mathbf{Q}$  cuts along the  $[001]$  direction through  $K = [0.9, 1.1]$  rlu in **a-d**. **k**, Similar as in **e** but through  $K = [-0.1, 0.1]$  rlu in **g-j**. The intensities of the cuts in **e** and **k** are offset for a better visualisation. **f** and **l**, Integrated intensities over  $L = [-1, 1]$  rlu in **e** and **k**, respectively. Solid curves in **f** and **l** are guides to the eye. Dashed lines denote the magnetic transition temperature  $T_N \sim 61$  K. Data in **a-d** and **g-j** are integrated over  $[H, 0, 0] = [-3.2, -2.8]$  rlu. Intensities of elastic scattering at  $T = 70$  K in **d-f** are due to nuclear scattering, which remain finite above  $T_N$ .



**Supplementary Figure 5.** A triply degenerate node at a lower energy of 16 meV. **a-c**, Contours plotted against two orthogonal axes  $[\bar{3}11]$  and  $[0\bar{1}1]$  with different energy intervals. Dashed arrows indicate the trajectory of the cuts plotted in **d**. The vertical dashed line in **d** denotes the position of the triply degenerate node. Lines through data are fits with Gaussian functions. Intensities in **d** are offset according to the energy. These results illustrate that the bands cross each other at the H point.

**a****b**

**Supplementary Figure 6.** Calculated dispersions near triply degenerate points. Calculated dispersions near the H point based on Eq. 1 in the main manuscript. Three magnon bands cross at the H point at two different energies. The magnon bands can be described by Supplementary Equation 2 effectively.