

Ohmic transition at contacts key to maximizing fill factor and performance of organic solar cells

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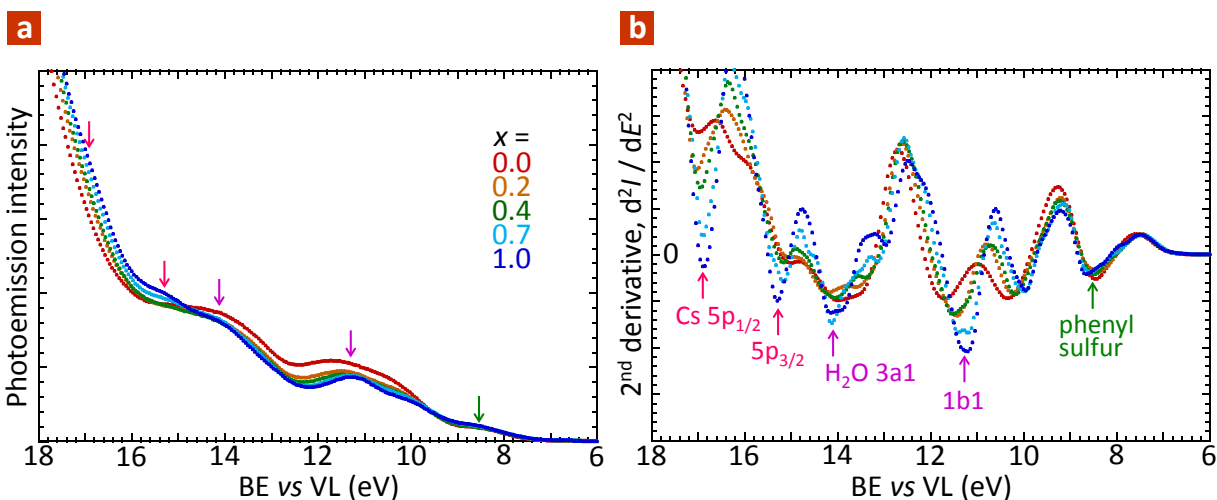
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Supplementary Information

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Supplementary Figure 1



Supplementary Figure 1. Valence band spectrum of PEDT: PSSCs_xH_{1-x} films. (a) Valence band spectrum, plotted against binding energy (BE) relative to vacuum level (VL) for film. Arrows indicate peak positions determined from second-derivative spectrum. (b) Second-derivative spectrum, obtained by digital derivatization of the spectrum in (a) with appropriate smoothing. Spectral assignments are marked. Phenyl refers to the localized phenyl MO of PSS; sulfur refers to the localized (lone-pair) sulfur MO of PEDT. Molecular H₂O is bound to the -SO₃⁻...Cs⁺ ion clusters, but not -SO₃H. The constant energies of these features reveal the absence of any contribution from a varying surface-dipole layer to the work function shift. Assignments are based on Chia et al¹ for phenyl and sulfur and Winter et al for the others.²

Supplementary Table 1

Parameter		ϕ (eV)					
		5.17	5.10	5.03	4.94	4.88	4.72
Electron mobility	μ_n			$9 \times 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-2}$			
Hole mobility	μ_p			$1 \times 10^{-4} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-2}$			
Electron density, at electron contact	N_n			$2.0 \times 10^{17} \text{ cm}^{-3}$			
Hole density, at hole contact	P_p	2.0	1.7	1.4	0.9	0.6	0.1
				$\times 10^{17} \text{ cm}^{-3}$			
Recombination factor	ξ			0.2			
Zero-K built-in potential	V_o			0.73 V			
Dielectric constant	ϵ_r			2.0			

Supplementary Table 1. Device parameter set employed for simulation of bulk *JV* characteristics for PEDT: PSSC_{S_x}H_{1-x}/ P3HT: PCBM/ Ca solar cells. The parameters have been validated in crosslinked donor polymer network: fullerene solar cells with well controlled donor–acceptor morphology over wide thickness range.³ P_p is a variable scaled by $(\phi - \phi_{\text{pin}})$.

Supplementary Note 1: Fermi-Level Pinning and Effective Work Function

We use ϕ to denote the vacuum work function of the electrode, in this case PEDT: PSS($\text{Cs}_x\text{H}_{1-x}$), which is the energy difference between its Fermi level (FL) and the outside vacuum level (VL). This can be measured, for example by ultraviolet photoemission spectroscopy (UPS).⁴⁻⁶ We denote the effective work function of the electrode in contact with the semiconductor by ϕ_{eff} , which is the energy difference between FL of the electrode and the effective VL of the bulk of the semiconductor.⁷ The effective VL is what pertains when the semiconductor is notionally split to create a vacuum interface but without creating a dipole layer at the vacuum interface. ϕ_{eff} rather than ϕ is the quantity that is relevant to energy-level alignment in devices. It can be inferred from V_{bi} measurements in diodes, assuming ϕ_{eff} at one of the contacts is known *a priori*, for example through systematic correlation.⁷ Surveys of polymer organic semiconductor diodes suggest that ϕ_{eff} is a well-defined and portable under some conditions, in particular in the absence of Fermi-level (FL) pinning.⁷ When ϕ exceeds the work function for the onset of FL pinning ϕ_{pin} , charge transfer to the semiconductor band edge occurs, which opens a VL offset that greatly slows down the creep of FL up the density-of-states of the semiconductor.^{8,9} As a consequence ϕ_{eff} appears to be pinned at ϕ_{pin} . The location of ϕ_{pin} can thus be obtained from the interpolated knee in the plot of ϕ_{eff} against ϕ . The resultant accumulation of carrier density in an ultrathin layer with diffused tail at the semiconductor side of the contact, which we refer to variously as the ‘ δ -doped layer’ or ‘ δ -density’, can be measured and quantified by subgap electroabsorption spectroscopy.¹⁰

Supplementary Note 2: Fermi-Level Pinning Calculations

Calculations were performed at 290 K with the following parameters:

Parameter	Symbol	Value
Gaussian density-of-states	σ_G	$3 \times 10^{13} \text{ cm}^{-2}$
Static dielectric constant	ϵ_r	2.0 ^a
Effective dielectric thickness	d	2.5 nm
Contact double-layer capacitance	C_{dl}	$0.7 \mu\text{F cm}^{-2}$

Footnotes:

^a We used a ϵ_r lower than 3 because of the absence of the conjugated polymer in the local ion cluster.

In a typical calculation, a certain density-of-states (DOS) function for the semiconductor is assumed. Then for each assumed effective work function of the hole contact ϕ_{eff} given by the work function for the semiconductor, we computed the vacuum-level offset at the contact $\Delta\phi_{\text{dl}}$ and the associated work function of the electrode ϕ , according to: $\phi = \phi_{\text{eff}} + \Delta\phi_{\text{dl}}$. $\Delta\phi_{\text{dl}}$ is given by the voltage drop associated with the double-layer capacitance due to the accumulated hole density and its image, which in the zero-K limit is given by: $\Delta\phi_{\text{dl}} = \frac{e\sigma}{C_{\text{dl}}}$, where σ is the accumulated hole density given by $\sigma = \int_{\epsilon_F}^{\infty} N_i(\epsilon_i)(1-f_i(\epsilon_i))d\epsilon_i$, where $N_i(\epsilon_i)$ is the DOS and $f_i(\epsilon_i)$ is the Fermi-Dirac function, and C_{dl} is the double-layer capacitance given by $C_{\text{dl}} = \frac{\epsilon_0\epsilon_r}{d}$, where ϵ_0 is the vacuum permittivity, ϵ_r is the local dielectric constant, and d is the effective double-layer thickness. For an interface with a perfect conductor, d is given by half the distance to the image charge. We have checked that a variation in $\Delta\phi_{\text{dl}}$ by up to a factor of three, due to uncertainty in C_{dl} , does not change the form of the plots. The FL pinning onset work function ϕ_{pin} is established from the kink in the ϕ_{eff} vs ϕ plot. This was computationally obtained by intersecting the unity-slope

straight line in the $\phi < \phi_{\text{pin}}$ regime, with a straight line fitted over $\phi_{\text{pin}} + 0.1 < \phi < \phi_{\text{pin}} + 0.4$ eV regime, following the usual experimental practice.

For a DOS with a hemi-gaussian frontier region with integrated hemi-gaussian DOS of 4×10^{13} cm^{-2} , our simulation gives $\phi_{\text{pin}} \approx (2.35 \pm 0.02) * \sigma_G$, where σ_G is the gaussian (standard deviation) width. For comparison, the ionization energy I_E conventionally defined by extrapolated onset of the highest-occupied molecular orbital (HOMO) band edge is at $2.00 * \sigma_G$. Therefore in the absence of polarization-induced band-bending, Fermi-level pinning is expected to occur very close to the conventionally-defined band edge, at *ca.* $0.35 * \sigma_G$ into the gap from the band edge. This amounts to 63 meV for a typical σ_G of 0.18 eV. Furthermore, we found for σ_G of 0.18 eV that contact resistivity scales as: $\rho_c \sim \sigma^{-0.82}$ for ϕ in the range of $\phi_{\text{pin}} \pm 0.2$ eV.

Supplementary References

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