## High-Performance Reconstruction of Microscopic Force Fields from Brownian Trajectories

Pérez García et al.



Supplementary Figure 1. Holographic optical tweezers setup. An expanded laser beam (wavelenght 532 nm) is reflected by a spatial light modulator (SLM, Hamamatsu X10468-04), through a telescope (lens L1, f = 400 mm, and lens L2, f = 175 mm), where the diffracted field of interest is discriminated by the iris diaphragm, a half-wave plate (HWP), a polarising beam splitter (PBS), and a second telescope (T2, magnification M = 3); it is then reflected by a dichroic mirror (DM), and finally goes through a quarter wave plate (QWP) to be focused by an objective (Carl Zeiss,  $100 \times$ , oil immersion, 1.30 NA) within the sample chamber. The trapped particle is visualised through an imaging system consisting of a diode-based white-light illumination, a condenser ( $20 \times$ , 0.75 NA), the  $100 \times$  objective, a tube lens (f = 200 mm), and a CMOS Camera (Basler acA800-510um with a ROI of  $256 \times 40$  pixels and exposure time of  $100 \,\mu$ s in the acquisition configuration). The shape of the beam is controlled by means of the SLM: (a) blazed grating phase modulation used to generate a single optical tweezers (Fig. 2); (b) blazed grating plus a l = 1 azimuthal phase to generate a LG<sub>1</sub> beam (Fig. 3); (c) two identical tweezers with a phase shift  $\pi$  to generate a multiwell potential (Fig. 4). (d) The QWP permits us to switch the polarisation state of the beam between linearly polarised ( $\beta = 0$ , Figs. 3d and 3g), circularly (+) polarised ( $\beta = +\pi/4$ , Figs. 3e and 3h), and circularly (-) polarised ( $\beta = -\pi/4$ , Figs. 3f and 3i).



Supplementary Figure 2. Performance of FORMA as a function of sampling time. Values obtained by FORMA from simulated timeseries of an optically trapped particle for (a) the trap stiffness k and (b) the diffusion coefficient D as function of the sampling frequency f. The vertical blue and red lines respectively indicate the characteristic frequency in the trap  $(f_c = k/\gamma = 60.3 \text{ s}^{-1})$  and the sampling frequency  $(f_s = 4504.5 \text{ s}^{-1})$  of the results reported in Figure 2. The shaded areas represent the regions where the error in the determinations is greater than 10%. In all cases, we have simulated 24 trajectories of the motion of a spherical microparticle with radius  $R = 0.5 \,\mu\text{m}$  in an aqueous medium of viscosity  $\eta = 0.0011 \text{ Pas}$ . The time for each simulations is fixed to 22 s.



Supplementary Figure 3. Examples of force fields. Examples of force fields and their decomposition into conservative and non-conservative components (a) Isotropic harmonic potential where  $k_1 = k_2 > 0$ . (b) Elliptical harmonic potential where  $0 < k_1 < k_2$ . (c) Saddle point where  $k_1 < 0$  and  $k_2 > 0$ . (d) Unstable point where  $k_1 < k_2 < 0$  and the axes are titled by  $\theta = 45^{\circ}$ . The full and empty circles indicate the stable and unstable equilibrium points respectively.



Supplementary Figure 4. Speckle optical tweezers setup. Optical micromanipulation with an image speckle. An expanded laser beam (wavelength 532 nm) is reflected by the spatial light modulator (SLM, Hamamatsu X10468-04), goes through a telescope (lens L1, f = 400 mm, and lens L2, f = 100 mm) where it is spatially filtered by an on-axis iris diaphragm (aperture  $D_i = 1.54$  mm), passes through a dichroic mirror (DM), is focused by lens L3 (f = 150 mm) on the back-focal plane of an objective (Olympus,  $40 \times$ , 0.65 NA). The particle is visualsed through an imaging system similar to that described in Supplementary Figure 1. (a) A uniform-distributed random phase mask is projected on the SLM to generate the speckle pattern (c), and (b) a ring-like dynamic phase distribution is used to set the initial position of the particle within the area of interest (d). Once the particle is in the desired initial position, the SLM projects the random phase distribution to generate the speckle pattern, and the particle position is recorded as it diffuses for about 6 seconds, then the speckle pattern is changed for the ring-like trap to set the initial position, this process is repeated 500 times. The polystyrene particle is trapped in 2D as it is pushed towards the upper coverslip by the beam radiation pressure.

$\operatorname{coord.}$	$k (\mathrm{pN}\mu\mathrm{m}^{-1})$	$k^* (\mathrm{pN}\mu\mathrm{m}^{-1})$	$\Omega(s^{-1})$	$\Omega^* \ (s^{-1})$	$D \ (\mu m^2 s^{-1})$	$D^* \ (\mu m^2 s^{-1})$
x	0.6	$0.60\pm0.02$	2	$2.27 \pm 1.15$	0.392	$0.387 \pm 0.001$
y	0.5	$0.49\pm0.02$	1.5	$1.57\pm0.81$	0.392	$0.387 \pm 0.002$
z	0.2	$0.20\pm0.01$	3	$2.86\pm0.78$	0.392	$0.389 \pm 0.001$

Supplementary Table 1. Estimated parameters in a non-conservative 3D trap using FORMA. Results of the analysis with FORMA of a 3D Brownian simulation, assuming a harmonic trap with different k's and  $\Omega$ 's along the three coordinate axes. We simulated a spherical particle with radius  $R = 0.5 \,\mu$ m in an aqueous medium of viscosity  $\eta = 0.0011 \,\text{Pas}$  at a sampling frequency  $f_s = 4504.5 \,\text{s}$ , and we used 24 windows of  $2 \cdot 10^5$  samples for the averaging.

Eq.	point	$x_{\rm eq}^*$ (µm)	$y_{\rm eq}^*$ (µm)	$k_1^*  (pN  \mu m^{-1})$	$k_2^* (\mathrm{pN}\mu\mathrm{m}^{-1})$	$\theta^*$ (deg)	$\Omega^* \ ({\rm s}^{-1})$	$N_{\rm s}$
	1	1.79	8.08	0.14	0.07	358	0.47	7486
	2	4.41	7.66	0.21	0.04	37	-0.07	175406
	3	1.41	7.38	0.20	-0.09	15	0.67	830
	4	2.11	7.08	0.03	-0.08	116	-0.69	562
	5	5.61	6.72	0.11	0.07	105	-4.64	1240
	6	5.21	6.58	0.04	-0.10	108	1.58	1084
	7	1.91	6.58	0.04	-0.11	63	-0.14	841
	8	0.75	6.22	0.10	0.08	56	-0.25	41667
	9	4.71	6.08	0.06	0.00	358	1.24	13144
	10	3.11	6.02	0.07	0.01	49	-1.43	5411
	11	3.81	5.98	-0.05	-0.06	32	-4.87	2101
	12	5.61	5.78	0.07	-0.05	124	0.59	626
	13	1.51	5.78	0.03	-0.12	43	1.09	1146
	14	3.21	5.58	-0.01	-0.05	12	2.10	2369
	15	3.61	5.38	-0.03	-0.07	11	-2.82	1829
	16	6.27	5.12	0.21	0.18	113	-0.00	123870
	17	4.85	5.08	0.02	-0.04	3	0.17	6154
	18	1.11	5.08	0.08	0.03	352	0.10	1340
	19	2.27	4.74	0.16	0.14	5	-0.07	486547
	20	3.43	4.60	0.07	-0.03	89	-0.86	4819
	21	5.43	4.44	0.02	-0.09	126	3.14	5379
	22	4.71	4.38	0.08	0.00	99	-0.32	26013
	23	1.41	4.38	0.11	-0.06	62	0.55	14948
	24	3.51	4.28	-0.03	-0.04	339	0.89	3920
	25	1.05	4.20	0.12	0.06	94	-0.53	32006
	26	6.01	3.98	0.02	-0.06	53	0.14	1188
	27	3.95	3.98	0.02	-0.12	330	-0.58	4590
	28	5.27	3.96	0.08	0.04	125	-0.74	17540
	29	2.61	3.78	0.04	-0.13	35	-2.89	2561
	30	2.41	3.48	-0.01	-0.11	55	-0.07	3546
	31	5.67	3.42	0.05	-0.04	332	-0.80	3564
	32	1.11	3.38	-0.03	-0.08	17	-2.75	1220
	33	4.31	3.34	-0.03	-0.08	356	-2.03	2505
	34	7.23	3.32	0.22	0.14	27	-0.30	42399
	35	3.19	3.24	0.10	0.05	73	-0.17	100763
	36	2.31	2.98	0.07	-0.05	78	0.18	3440
	37	5.07	2.92	0.13	-0.02	80	-3.78	1729
	38	5.91	2.88	0.02	-0.04	71	4.53	1411
	39	4.81	2.68	-0.00	-0.10	352	1.58	1407
	40	5.71	2.58	0.03	-0.07	121	2.57	1449
	41	3.71	2.48	0.05	-0.01	356	-1.62	7907
	42	1.45	2.44	0.14	0.13	21	-0.64	57803
	43	3.01	2.08	0.05	-0.10	106	-4.62	1529
	44	1.01	0.76	0.16	0.07	76	-0.27	18292
	45	4.69	0.72	0.10	-0.03	105	-0.22	16629
	46	4.05	0.72	0.09	0.02	39	0.21	34384
	47	5.15	0.56	0.14	0.06	111	-0.49	27232

Supplementary Table 2. Equilibrium points identified by FORMA in a speckle parttern. For each equilibrium point identified in the data shown in Figure 5, FORMA provides the position  $(x_{eq}^*, y_{eq}^*)$  (from the lower left corner of Figure 5a), its stiffnesses  $k_1^*$  and  $k_2^*$  along the principle axes, the orientation  $\theta^*$  of the principle axes with respect to the Cartesian axes, and the frequency  $\Omega^*$  associated to the non-conservative (rotational) component of the force field.  $N_s$  is the number of measurements of the particle displacements used by FORMA for the estimation.