

Supporting Information for

**Weak Shape Anisotropy Leads to a  
Non-monotonic Contribution to Crowding  
Impacting Protein Dynamics under  
Physiologically Relevant Conditions**

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## Deformation of ellipsoids

The deformation of ellipsoids is calculated by the difference between the measured radius and the initial radius of ellipsoids. (Figure S1).

## Intermediate scattering function

The short-time collective diffusion coefficient  $D_s(q)$  (cf. Figure 2) is extracted from a single exponential fit to the short-time part of the intermediate scattering function

$$S(q, t) \approx S(q, 0)e^{-q^2 D_s(q)t} \quad (\text{S1})$$

where the intermediate scattering function (Figure S2A) is calculated by

$$S(q, t) = \left\langle \frac{1}{N} \sum_{j,k} e^{i\mathbf{q} \cdot (\mathbf{r}_j(0) - \mathbf{r}_k(t))} \right\rangle. \quad (\text{S2})$$

Note that only the short-time part of  $S(q, t)$  is used for the fit ( $t \leq 0.2\tau$ , where  $\tau = R_h^2/D_0$  is the characteristic time of colloids).

## Mean-square-displacement

The short-time self-diffusion coefficient  $D_s^s$  (cf. Figure 3) is calculated from the mean-square-displacement  $\langle \Delta r^2 \rangle = 6D_s^s t$  of colloids (Figure S2B). Note that only the short-time linear part of the MSD is used for the fit ( $t \leq 0.6\tau$ ).

## Number of connected neighbors

The local environment of individual colloids is analyzed by the average number of connected neighbors  $N_{nbr}$  (i.e. belonging to the same cluster) within a cutoff radius  $R_{cut} = 7R_h$  of attractive spheres and ellipsoids as a function of volume fraction (Figure S3), which is normalized by the number density,  $N_{nbr}^* = N_{nbr}/[(N/L_s^3) \times (\frac{4}{3}\pi R_{cut}^3) - 1]$ . Here, two colloids are considered as connected neighbors when their bead separation is smaller than the distance at the first minimum of the pair correlation function ( $r \leq 1.4\sigma$ ), and the cutoff radius  $R_{cut}$  is set to eliminate the effect of the simulation box size in the analysis.

### **Orientational angle distribution**

The orientational ordering of ellipsoids is analyzed by the distributions  $P(\cos(\theta))$  of angles  $\theta$  between neighboring ellipsoids (Figure S4). The orientation of an ellipsoid is defined by its semi-principal axis  $r_a$ , where the angle  $\cos(\theta)$  of two ellipsoids varies from 0 for perpendicular orientation to 1 for parallel orientation.

### **Structure factor**

The characteristic of the structure can be analyzed by the structure factor (cf. Figure 5)

$$S(q) = 1 + \frac{N}{V} \int_0^\infty 4\pi r^2 [g(r) - 1] \frac{\sin(qr)}{qr} dr \quad (\text{S3})$$

where the the pair-correlation function

$$g(r) = \frac{V}{N^2} \left\langle \sum_i \sum_{j \neq i} \delta(\mathbf{r} - \mathbf{r}_{ij}) \right\rangle \quad (\text{S4})$$

is the probability to find a pair of colloids a distance  $r$  apart.

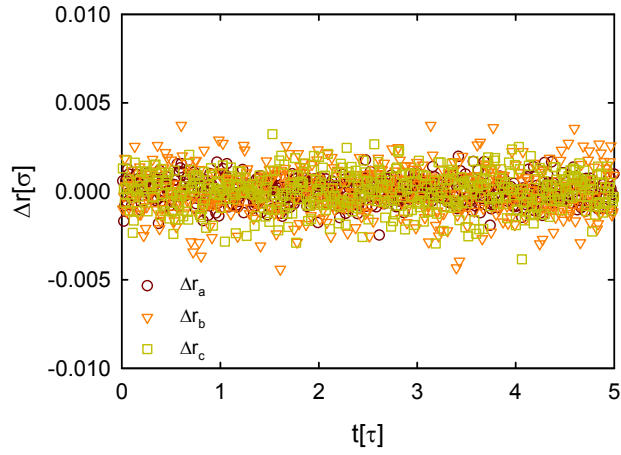


Figure S1: Deformation of ellipsoids.

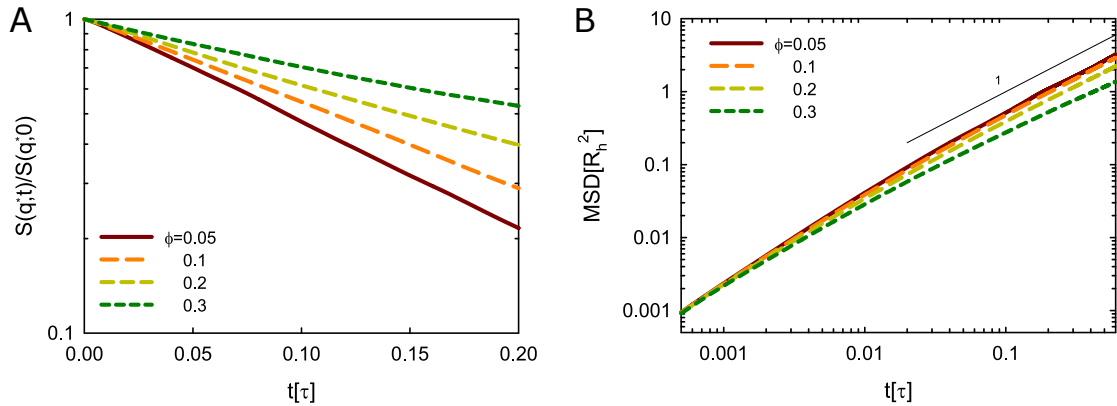


Figure S2: (A) Intermediate scattering function (ISF) for different volume fractions of hard ellipsoid. The ISF of  $0.02\tau \leq t \leq 0.2\tau$  is used for the fit to obtain the short-time collective diffusion coefficient  $D_s(q)$ . (B) Mean-square displacement (MSD) for different volume fractions of hard ellipsoid. The MSD of  $0.02\tau \leq t \leq 0.6\tau$  is used for the fit to obtain the short-time self-diffusion coefficient  $D_s^s$ .

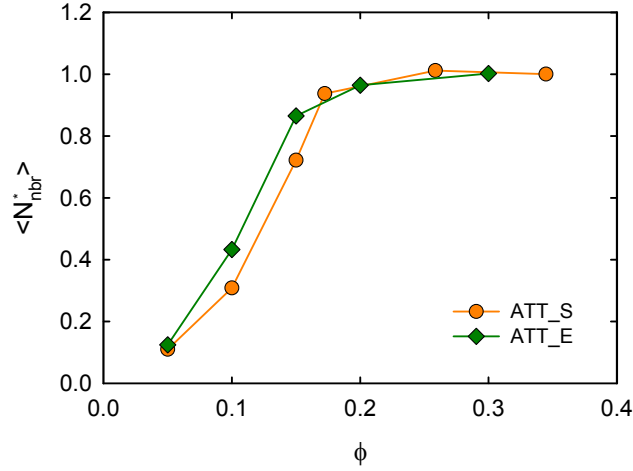


Figure S3: Normalized average number of connected neighbors  $N_{nbr}^*$  within a cutoff radius  $R_{cut} = 7R_h$  of attractive spheres (ATT\_S) and ellipsoids (ATT\_E) as a function of volume fraction.

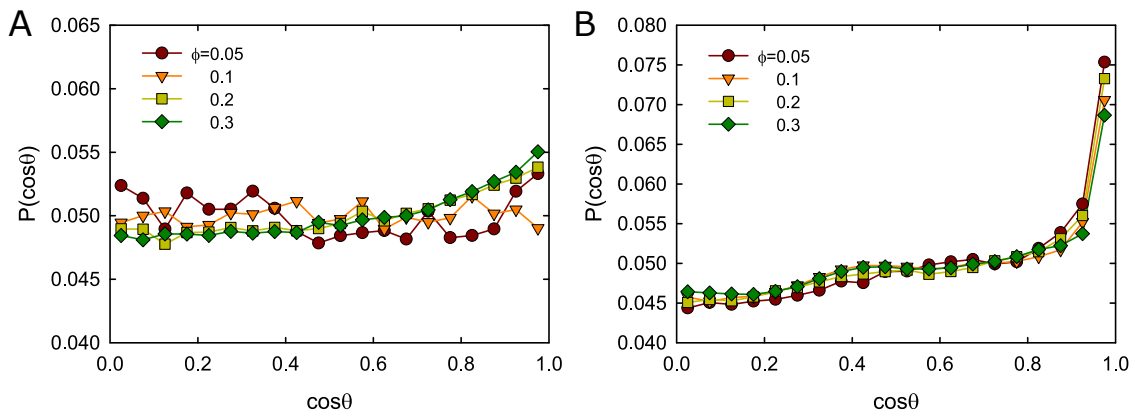


Figure S4: Distributions of angles  $\theta$  between neighboring (A) hard and (B) attractive ellipsoids for different volume fractions.