Supplementary Information

Directly visualizing the sign change of *d*-wave superconducting gap in

$Bi_2Sr_2CaCu_2O_{8+\delta}$ by phase-referenced quasiparticle interference

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Supplementary Figure 1 Measured QPI images and corresponding FT-QPI patterns at different energies on Bi-2212. $V_{set} = -100 \text{ mV}$, $I_{set} = 100 \text{ pA}$.



Supplementary Figure 2 Theoretical calculated tunneling spectra and PR-QPI results in a *d*-wave superconductor. **a** Fermi surface and gap function calculated for a *d*-wave superconductor. The inset in a shows the calculated Fermi surface and the definition of the parameter angle θ . The *d*-wave gap function $\Delta(\mathbf{k}) = \Delta_0(\cos k_x - \cos k_y)$ with $\Delta_0 = 23$ meV is assigned to the Fermi surface, and the resultant angle θ dependent superconducting gap along the Fermi surface is shown in a. The gap function along the Fermi surface (open symbols) is consistent with a *d*-wave function $\Delta(\theta) = 44.5 \cos 2\theta$ meV. **b** The tunneling spectra in an impurity-free area and on site of the non-magnetic impurity. The scattering scalar potential is set to $V_{\rm s}=20$ meV for the non-magnetic impurity. c Simulated LDOS at -15 meV around the single non-magnetic impurity with dimensions of 512×512 atom lattice. The impurity is set to be at the center of the image. **d**, **e** FT-QPI patterns for the single impurity at ± 15 meV. **f** The resultant PR-QPI image $g_r(\mathbf{q}, -15 \text{ meV})$ calculated by DBS-QPI method from the simulation results in d and e. g Average intensity per pixel for the characteristic scattering spots in **f**. One can find that the $g_r(\mathbf{q}, -15 \text{ meV})$ values near gap-sign-preserved scattering vectors \mathbf{q}_1 , \mathbf{q}_4 , and \mathbf{q}_5 are positive, while the values near gap-sign-reversed scattering vectors \mathbf{q}_2 , \mathbf{q}_3 , \mathbf{q}_6 , and \mathbf{q}_7 are negative.



Supplementary Figure 3 Simulated results around a single non-magnetic impurity with different positive scattering potentials in a *d*-wave superconductor.



Supplementary Figure 4 Simulated results around a single non-magnetic impurity with different negative scattering potentials in a *d*-wave superconductor.



Supplementary Figure 5 Topographic image ($V_{\text{bias}} = -200 \text{ mV}$, $I_{\text{t}} = 50 \text{ pA}$) and tunneling spectra measured on Bi-2212. The spectra are measured along different arrowed lines in the topographic image.



Supplementary Figure 6 Control experiments of FT-QPI patterns and corresponding PR-QPI patterns at different energies measured in area-1 on sample 2. The left two columns show the FT-QPI intensity maps. The 3rd column shows the PR-QPI patterns out of the corresponding data. The 4th column shows the PR-QPI intensity of the seven characteristic spots derived from the 3rd column. The PR-QPI pattern at -20 meV is the same as the one in Fig. 6a. $V_{set} = -100$ mV, $I_{set} = 100$ pA.



Supplementary Figure 7 Control experiments of FT-QPI patterns and corresponding PR-QPI patterns at ±20 meV measured in area-2 of sample 2 and in an area of sample 3. The left two columns show the FT-QPI intensity maps. The 3rd column shows the PR-QPI patterns out of the corresponding data. The 4th column shows the PR-QPI intensity of the seven characteristic spots derived from the 3rd column. $V_{set} = -100$ mV, $I_{set} = 100$ pA.



Supplementary Figure 8 Theoretical simulated PR-QPI results for the case of multiple non-magnetic impurities with a *d*-wave superconducting gap function. **a**-**e** Simulation results for the case of 60 same impurities with same scattering potential $V_s = 20$ meV. These impurities are randomly distributed in an area with dimensions of 512×512 atom lattice as shown in **a**. **f**-**t** Simulation results for the cases of 60 (**f**-**j**), 100 (**k**-**o**) and 500 (**p**-**t**) randomly distributed impurities with the random scattering potential values from 0 to 100 meV. The areas are also with dimensions of 400×400 atom lattice. One can see that the signs of PR-QPI signal for the characteristic scattering spots are the same as the ones of the single impurity situation.



Supplementary Figure 9 Theoretical simulated PR-QPI results with a sign-preserved nodal gap function. a The tunneling spectra in an impurity-free area and on site of the non-magnetic impurity for a superconductor with a nodal but sign-preserved gap. The Fermi surface is the same as the one in Supplementary Figure 2, and the superconducting gap function used in the calculation is $\Delta(\mathbf{k}) = \Delta_0 |(\cos k_x - \Delta_0)| |(\cos k_$ $\cos k_y$) with $\Delta_0 = 23$ meV. The scattering scalar potential is set to $V_s = 20$ meV for the non-magnetic impurity. **b** Simulated LDOS around a single non-magnetic impurity at -15 meV with dimensions of 512×512 atom lattice. The impurity is set to be at the center of the image. c,d FT-QPI patterns at ± 15 meV for the single impurity. e The resultant PR-QPI image $g_r(\mathbf{q}, -15 \text{ meV})$ calculated by DBS-QPI method from the simulation results in c and d. f Average intensity per pixel for the characteristic scattering spots in e. One can see that the $g_r(q, -15 \text{ meV})$ values near all characteristic scattering vectors are positive. **g-k** The corresponding simulation results for the case of multiple impurities. The 60 impurities are the same as the one in **b**, and they are randomly distributed in an area with dimensions of 512×512 atom lattice as shown in g. The locations of impurities are the same as the ones in Supplementary Figure 8a.



Supplementary Figure 10 Theoretical simulated PR-QPI results with a sign-preserved nodeless gap function. a The tunneling spectra in an impurity-free area and on site of the non-magnetic impurity for a superconductor with a sign-preserved nodeless gap. The Fermi surface is the same as the one in Supplementary Figure 2, and the superconducting gap function used in the calculation is $\Delta(\mathbf{k}) = \Delta_0 |\cos k_x - \Delta_0| \cos k_x |\mathbf{k}|$ $\cos k_y ig|$ + $\Delta_{
m m}$ with $\Delta_0 = 23$ meV, $\Delta_{
m m} = 2\,$ meV. The scattering scalar potential is set to $V_s = 20$ meV for the non-magnetic impurity. **b** Simulated LDOS around a single non-magnetic impurity at -15 meV with dimensions of 60×60 atom lattice. The impurity is set to be at the center of the image. c,d FT-QPI patterns for the single impurity at ± 15 meV. **e** The resultant PR-QPI image $g_r(\mathbf{q}, -15 \text{ meV})$ calculated by DBS-QPI method from the simulation results in c and d. f Average intensity per pixel for the characteristic scattering spots in **e**. One can see that the $q_r(\mathbf{q}, -15 \text{ meV})$ values near all characteristic scattering vectors are positive. g-k The corresponding simulation results for the case of multiple impurities. The 60 impurities are the same as the one in a, and they are randomly distributed in an area with dimensions of 512×512 atom lattice as shown in g. The locations of impurities are the same as the ones in Supplementary Figure 8a.



Supplementary Figure 11 Simulated results around a single magnetic impurity with different magnetic scattering potentials in a *d*-wave superconductor.