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2 **Supplementary Information for**
3 **Increasing valence pushes DNA nanostar networks to the**
4 **isostatic point**

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8 **This PDF file includes:**

- 9 Supplementary text
10 Figs. S1 to S13
11 Tables S1 to S2
12 References for SI reference citations

13 Supporting Information Text

14 S1. DNAns-specific information

15 This section includes information regarding the DNA sequences used for each DNA Nanostars (DNAns)
16 design, DNA concentration measurements, binding probability calculations, and hysteresis of structure
17 formation:

18 **S1a. DNA sequences.** Each of the f single-stranded (ss) DNA segments required to make $f = 3, 4, 5$, and
19 6 DNAns are 49 nucleotides long. In each sequence, 40 of the nucleotides are designed to form portions of
20 an arm, 6 constitute the sticky-end, and 3 are unbound (Fig. S1).

21 **S1b. Binding probability and mean valence calculations.** We can get a sense of the magntiude of the
22 change in connectivity from calculating the binding probability, p_b , between two DNAns arms. p_b was
23 calculated in a mean-field approximation that ignores connectivity using the DNA thermodynamics package
24 NUPACK (1). For our overhang sequence 5'-ACGTACG-3', and an oligo concentration of $f \cdot [\text{DNAns}]$, we
25 find that $p_b > 0.84$ at $T = 35^\circ\text{C}$, the highest temperature explored, and grows slowly to $p_b \approx 1$ at $T = 5^\circ\text{C}$
26 (see Fig. S2). This leads to a predicted mean-valence, $\langle z \rangle \approx p_b f \approx 0.98 f$ at $T_{\text{ref}} = 20^\circ\text{C}$. The use of such a
27 model to calculate p_b for the DNA nanostar system was quantitatively validated by Rovigatti *et. al.* (2) in
28 comparison to DNAns simulations that incorporate connectivity.

29 **S1c. Concentration measurements.** We used a droplet spectrophotometer (Nanodrop 2000c; Fisher Sci-
30 entific) to measure the absorption at 260 nm (A260) and derive the DNAns concentration, $[\text{DNAns}]$, of
31 f -armed DNAns solutions (Fig. S3 and Fig. S4). The A260 of an f -armed DNAns solution is related to its
32 $[\text{DNAns}]$ by

$$33 \quad [\text{DNAns}] = \text{A260} / (\sum_i \epsilon_i) \quad [1]$$

34 where ϵ_i is the extinction coefficient of the i^{th} ss DNA segment composing the f -armed DNAns (see Fig.
35 S1 for the extinction coefficients). Note that equation 1 assumes all DNAns in solution have completely
36 denatured into their single stranded components.

37 The DNAns concentrations of solutions prepared for the rheometer were so large that their A260 could
38 not be measured directly. The A260 of solutions diluted 100-fold, $\text{A260}_{1\%}$, in pure water was measurable,
39 but we questioned whether this dilution was sufficient to completely denature the DNAns and permit the
40 use of equation 1. To check, we compared the A260 of a 100-fold dilution of the most concentrated sample
41 of a given valence, $\text{A260}_{1\%}^*$, to the A260 of a 10,000-fold dilution made in pure water from that same
42 100-fold dilution, $\text{A260}_{0.01\%}^*$. In all cases, $100 \cdot \text{A260}_{0.01\%}^* \geq \text{A260}_{1\%}^*$, consistent with incomplete denaturation
43 at 100-fold dilution (Table S1).

The proportion of single-stranded DNA present in the 100-fold diluted solutions, α , can be calculated
based on the extinction coefficients of single-stranded (ss) and double-stranded (ds) DNA (3)

$$44 \quad \epsilon_{\text{ss}} = [37 \mu\text{g}/(\text{mL} \cdot \text{cm})]^{-1} \quad \text{and} \quad \epsilon_{\text{ds}} = [50 \mu\text{g}/(\text{mL} \cdot \text{cm})]^{-1}$$

given that the absorbance of the 1% solution is due to both ss and ds DNA

$$45 \quad \text{A260}_{1\%}^* = \alpha \frac{c}{37} + (1 - \alpha) \frac{c}{50} \quad [2]$$

46 and assuming the absorbance of the 0.01% solution is due to only ss DNA

$$47 \quad \text{A260}_{0.01\%}^* = 0.01 \frac{c}{37}, \quad [3]$$

48 where c is the total mass concentration of DNA present in the 1% solution. Dividing equation 2 by 3, we
49 can derive an expression for α :

$$50 \quad \alpha = \left(\frac{\text{A260}_{1\%}^*}{\text{A260}_{0.01\%}^*} - 74 \right) / 26 \quad [4]$$

51 We found α ranged from $\approx 30\%$ for the most concentrated 3-arm nanostar solutions, to $\approx 96\%$ for the
52 most concentrated 6-arm nanostar solutions (Table S1). We note that we assumed α was concentration
53 independent and only measured the α of the most concentrated solution at a given f in order to interpret
54 the $A260_{1\%}$ of all (less concentrated) solutions of a given f -armed nanostar and, thus, avoid an excessive
55 amount of multi-step dilutions.

56 Since α tracks the correct portion of ss and ds DNA present at the 1% dilutions, we used it to correct
57 the measured [DNAns] from the respective $A260_{1\%}$ measurement because using equation 1 assumes the
58 1% solution contains only ss DNA (*i.e.*, equation 1 assumes $A260_{1\%} = c/37$). The correction is found by
59 dividing the ideal A260 for ss DNA solutions, $A260_{ss} = c/37$, by equation 1:

$$60 \quad \text{correction} = A260_{ss}/A260_{1\%}^* = 1/(0.24 \cdot \alpha + 0.76) \quad [5]$$

61 For the lowest and highest measured α , the correction to the concentration was $\approx 23\%$ and $\approx 1\%$, respectively
62 (Table S1). We note that the error introduced to the [DNAns] by assuming a concentration-independent α
63 has little to no effect ($\ll 1\%$ error) on the scaling exponents measured from the G'_p ([DNAns]) data because
64 the relative ratio of concentrations (*i.e.*, equation 1) is kept the same when using a concentration-independent correction.

65 **S1d. Hysteresis measurements.** We tested for hysteresis of structure formation by repeatedly measuring
66 the storage, G' , and loss, G'' , modulus as a function of temperature at $\gamma = 5\%$, $\omega = 10$ rad/s, and at two
67 different rates of temperature change, $|dT/dt| = 1.5$ and $15^\circ\text{C}/\text{min}$. G' and G'' were independent of the
68 rate of cooling and heating for rates $|dT/dt| \leq 15^\circ\text{C}/\text{min}$ for all f and [DNAns] (Fig. S5). The rate at
69 which temperature changed between the different temperature measurements of $G'(\omega, T)$ and $G''(\omega, T)$
70 contributing to the master curve was $|dT/dt| \leq 15^\circ\text{C}/\text{min}$. Thus, all DNAns networks studied here form
71 reversibly and are always in equilibrium.

72 S2. Frequency sweep technical information

73 This section includes information regarding time-temperature superposition (*i.e.*, master curve construction)
74 and network relaxation time:

75 **S2a. Time-temperature superposition.** Time-temperature superposition (TTS) constructs a material's
76 viscoelastic response over a wide range of frequencies by superimposing frequency sweeps made at different
77 temperatures onto one another. The resulting superimposed curve is known as the "master curve". We used
78 the TTS protocol provided in the rheometer's proprietary software (TRIOS v.4.0.1, TA Instruments) to
79 construct master curves at a reference temperature of $T_{\text{ref}} = 20^\circ\text{C}$, allowing for both vertical and horizontal
80 shifts of the raw frequency sweeps. An example of the TTS process is illustrated in Figure S6. Figures
81 S7A and S7B show all vertical and horizontal shift factors from the TTS procedure vs. temperature for all
82 valences.

83 The vertical shift factor accounts for slight changes in connectivity (Fig. S2) that result in an $\approx 40 - 80\%$
84 increase in G' (Fig. S7B). The horizontal shift factor accounts for changes in DNAns binding dynamics and
85 exhibits an Arrhenius temperature-dependence (see section S2b, Fig. S8).

86 **S2b. Relaxation times and Arrhenius dependence of relaxation times.** The horizontal shift factor is the
87 ratio between the DNAns network relaxation times at two different temperatures. For all f and [DNAns],
88 the relaxation times show an Arrhenius temperature behavior, $\tau_c \propto \exp(E_A/RT)$ (Fig. S8), with an
89 activation energy, E_a , that is independent of [DNAns] and f . Averaged over all [DNAns] and valences,
90 $\langle E_a \rangle = 210 \pm 10$ kJ/mol.

91 This value is consistent with the enthalpy of binding of the overhang, $|\Delta H| = 239$ kJ/mol, which was
92 calculated using the two-state melting application on the DINAmelt web server for overhang concentrations
93 between $600\mu\text{M}$ and $3600\mu\text{M}$ at $T = 20^\circ\text{C}$ and 150mM NaCl (4, 5). Thus, for all f and [DNAns], the
94 relaxation time τ_c of the network is determined by the kinetics of a single DNAns-DNAns bond.

95 **S2c. Numerical estimate of solvent-drag-induced crossover frequency ('non-affine-to-affine model').** It
 96 has been shown that a Maxwellian crossover in G' and G'' can arise from solvent drag inducing a non-affine
 97 to affine transition in networks (6–8). Following ref. (6), a crossover induced by solvent drag should occur
 98 at a frequency $\omega_c \approx K_c/\alpha$, where K_c is the network chain's spring constant, and α is the drag coefficient
 99 of a node. We showed in the main text that $K_c \approx 10^{-4}$ N/m for a simple FJC model of DNAns arms;
 100 this was validated by the similarity of this value to estimated cluster spring constants. We then estimate
 101 the drag from a Stoke's picture: the drag of a sphere is $\alpha = 6\pi\eta a \approx 1.5 \times 10^{-10}$ N s/m, where $\eta \approx 1$
 102 mPa s is the water viscosity, and the sphere radius $a \approx 8$ nm is set to the DNAns arm length. Note that
 103 this overestimates the drag, since the DNAns does not occupy all volume within the sphere, and leads
 104 to an underestimate of ω_c . Regardless, the crossover frequency is then $\omega_c \approx 7 \times 10^5$ Hz, which, while an
 105 underestimate, exceeds the measured crossover frequencies by five orders of magnitude. This indicates that
 106 the measured crossover is not caused by solvent drag, and supports the contention that it is instead caused
 107 by the kinetics of bond breakage (Section S2b).

108 S3. Strain sweep specific information

109 This section includes information regarding the strain sweep measurement, characteristic and yield strain
 110 measurements, and the uniaxial extension ratio calculation:

111 **S3a. Linear and non-linear elasticity of DNAns networks.** To investigate the linear and non-linear response
 112 of DNAns networks, we measured the storage, G' , and loss, G'' , modulus over a range of strains $1\% \leq$
 113 $\gamma \leq 200\%$ at a single frequency, $\omega > 62$ rad/s, and temperature, $T = 20^\circ\text{C}$, using a discrete logarithmic
 114 sampling distribution of 10 points per decade. For all DNAns networks, G' and G'' were strain-independent
 115 for $\gamma \leq 10\%$, consistent with measuring a linear elastic response to applied deformation (Fig. S9). After the
 116 strain-independent regime, all networks, except for the two highest concentrated $f = 6$ networks, displayed
 117 resolvable strain hardening until yielding (Fig. S9). The hardening response of such networks are well fit by
 118 $G' \sim \exp((\gamma/\gamma^*)^2)$, where γ^* is the characteristic strain at which stiffening occurs (9–11) (see section S3b).

119 We observed that both the torque amplitude and radial displacement (*i.e.*, strain) amplitude time-
 120 trace were continuous at yielding (Fig. S10), consistent with rheometer plates maintaining contact to an
 121 elastic material. If the plate lost contact with the material at yielding, both traces would appear highly
 122 discontinuous and noisy.

123 **S3b. Characteristic and yield strain.** The non-linear elastic response of nearly all DNAns networks followed
 124 the relation $G'(\gamma) \sim \exp((\gamma/\gamma^*)^2)$, where γ^* is the characteristic strain at which strain hardening occurs.
 125 At a given f , γ^* remained either independent of or slightly increased with [DNAns] (Fig. S11A). On the
 126 other hand, as a function of f , γ^* decreased from $\approx 200\%$ at $f = 3$ to $\approx 110\%$ at $f = 6$ (Fig. S11A).

127 Given the digital nature of the non-linearity measurement, we defined the maximum strain, γ_m , as the
 128 strain at which $G'(\gamma)$ maximized. Moreover, we defined the error in maximum strain, γ_m , as the resolution of
 129 the data sampling. Note that a logarithmic sampling decreases in resolution as strain increases. Accordingly,
 130 we report larger errors for strains of greater values. Despite the large uncertainty in our measurement of
 131 γ_m , we observed that γ_m remained roughly independent of or slightly decreased with [DNAns] (Fig. S11B).
 132 As a function of f , γ_m decreased from $\approx 170\%$ at $f = 3$ to $\approx 20\%$ at $f = 6$ (Fig. S11B).

133 **S3c. Uniaxial extension ratio calculation.** The uniaxial extension ratio, λ_{\max} , of a polymer chain in a
 134 network is related to the characteristic strain of the network by (9):

$$135 \quad \gamma^* = \lambda_{\max} - \lambda_{\max}^{-1} \quad [6]$$

136 At a given f , λ_{\max} remained roughly independent of or slightly increased with [DNAns] (Fig. S13). As a
 137 function of f , λ_{\max} decreased as valence grew: for $f = 3, 4$, and 5 , $\lambda_{\max} = 2.3 \pm 0.1, 1.7 \pm 0.1$, and 1.9 ± 0.2 ,
 138 respectively (Fig. S13). For the lowest concentrated $f = 6$ DNAns network, $\lambda_{\max} = 1.7 \pm 0.2$.

Arm number	Strand	Sequence (5' to 3')	Extinction Coefficient (L/(mole*cm))
f = 3	1	CTACTATGGCGGGTGATAAAAACGGGAAGAGCATGCCCATCCACGATCG	484100
f = 3	2	GGATGGGCATGCTCTTCCCGAACTCAACTGCCGTGGTATACGACGATCG	458900
f = 3	3	CGTATCACCAGGCAGTTGAGAAATTTATCACCCGCCATAGTAGACGATCG	480500
f = 4	1	same as f = 3 strand 1	484100
f = 4	2	same as f = 3 strand 2	458900
f = 4	3	CGTATCACCAGGCAGTTGAGAACATGCGAGGGTCCAATACGACGATCG	483300
f = 4	4	CGGTATTGGACCCCTCGCATGAATTTATCACCCGCCATAGTAGACGATCG	469000
f = 5	1	same as f = 4 strand 1	484100
f = 5	2	same as f = 4 strand 2	458900
f = 5	3	same as f = 4 strand 3	483300
f = 5	4	CGGTATTGGACCCCTCGCATGAACCATGCTGGACTCAACTGACACGATCG	462500
f = 5	5	GTCAGTTGAGTCCAGCATGGAATTTATCACCCGCCATAGTAGACGATCG	479800
f = 6	1	same as f = 5 strand 1	484100
f = 6	2	same as f = 5 strand 2	458900
f = 6	3	same as f = 5 strand 3	483300
f = 6	4	same as f = 5 strand 4	462500
f = 6	5	GTCAGTTGAGTCCAGCATGGAACGCATCAGTTGCGGCGCCGACGATCG	465500
f = 6	6	GCGGCGCCGCAACTGATGCGAATTTATCACCCGCCATAGTAGACGATCG	467200

Fig. S1. Table of sequences used for each DNAs design and their respective extinction coefficients (taken from the oligo specification sheets sent by manufacturer). Complementary portions are colored identically, the sticky-end is highlighted in blue, and unbound nucleotides are highlighted in red.

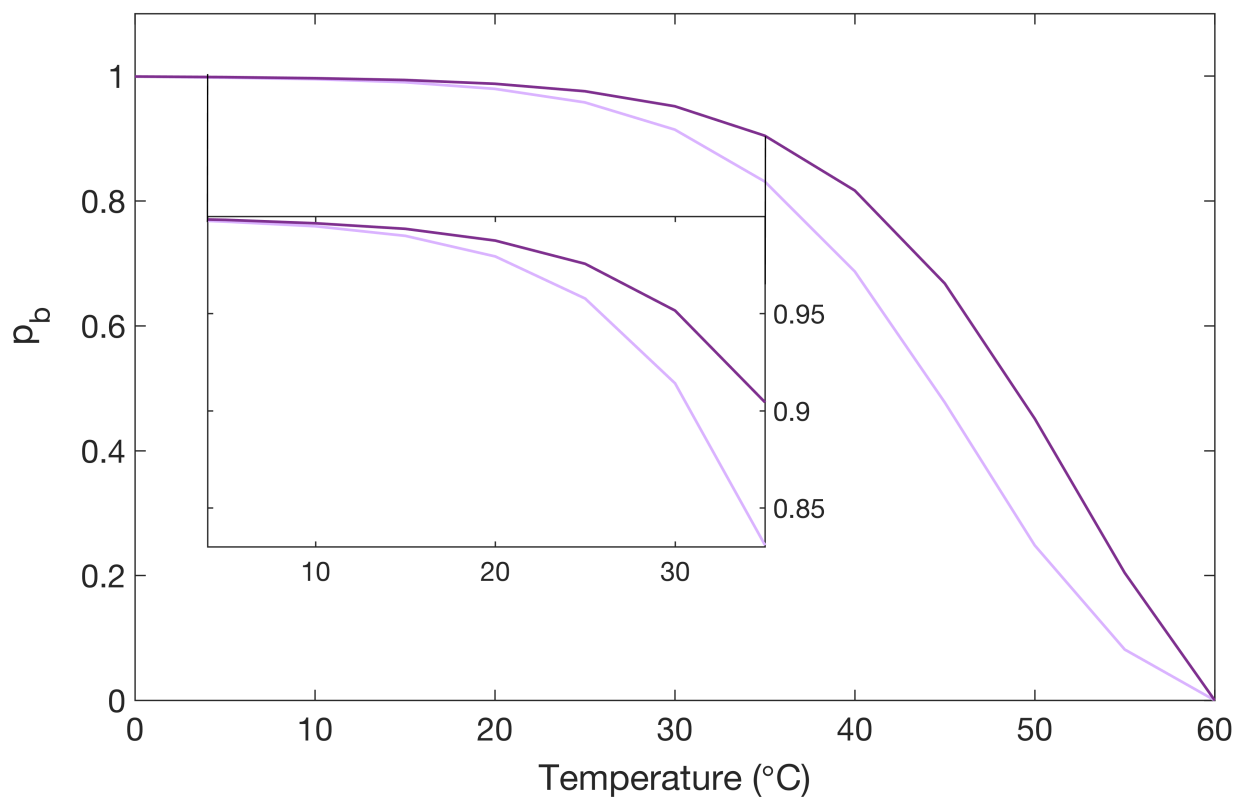


Fig. S2. Binding probability, p_b , of overhang sequence (5'-ACGATCG-3') duplex formation versus temperature at 150 mM NaCl. The light and dark purple line denote $[\text{DNANs}] = 600 \mu\text{M}/f$ and $[\text{DNANs}] = 3600 \mu\text{M}/f$, respectively. We note that the p_b reported is normalized to account for the two unbound A's in the DNANs-DNANs overhand bond, which were incorporated in the NUPACK calculation to account for their steric effects on duplex formation.

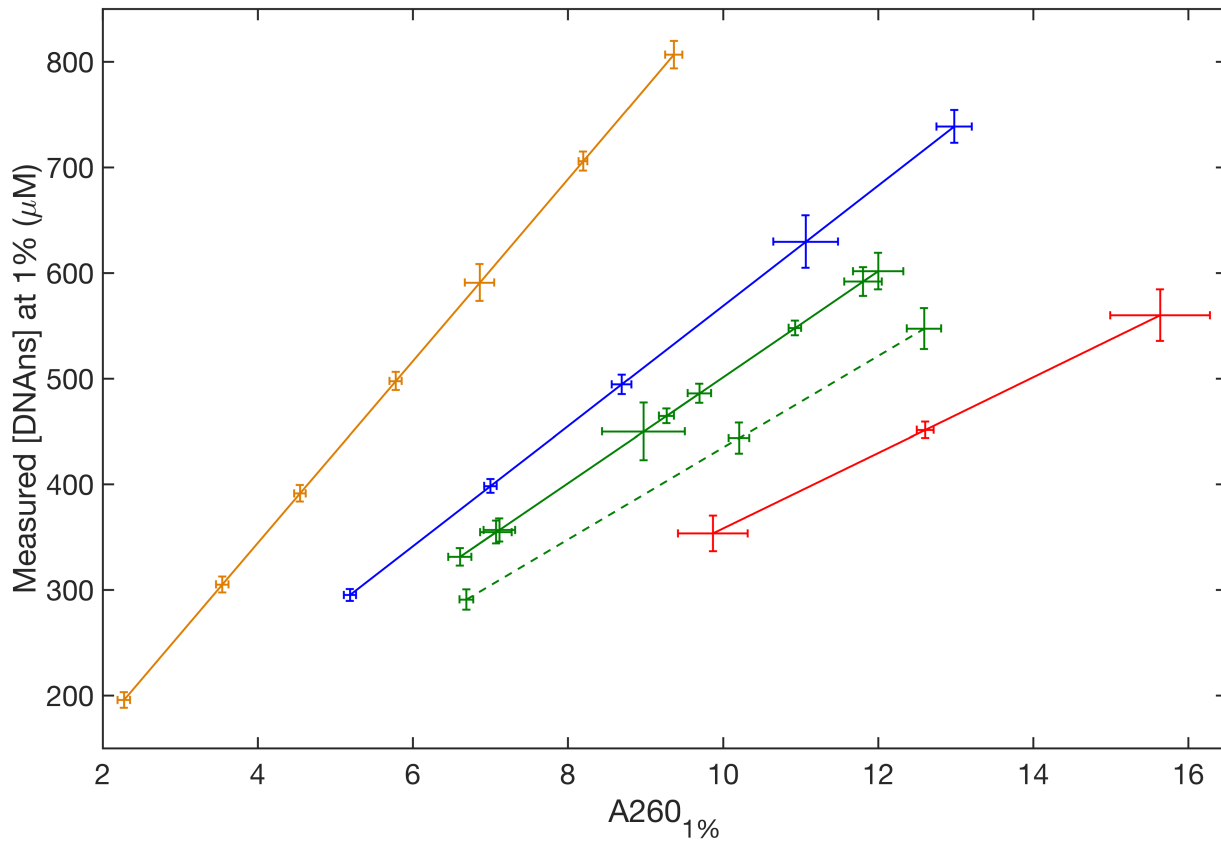


Fig. S3. Plot of measured [DNAs], with applied correction factors shown in Table S1, as a function of the measured A260 in the 1% dilutions, $A260_{1\%}$. The orange, blue, green, and red denote $f = 3, 4, 5,$ and $6,$ respectively. The lines are not fits, but an interpolated line connecting the data points meant as a guide for the eye. The dashed line denotes the second batch of $f = 5$ samples that was made several months after the first batch of $f = 5$ samples.

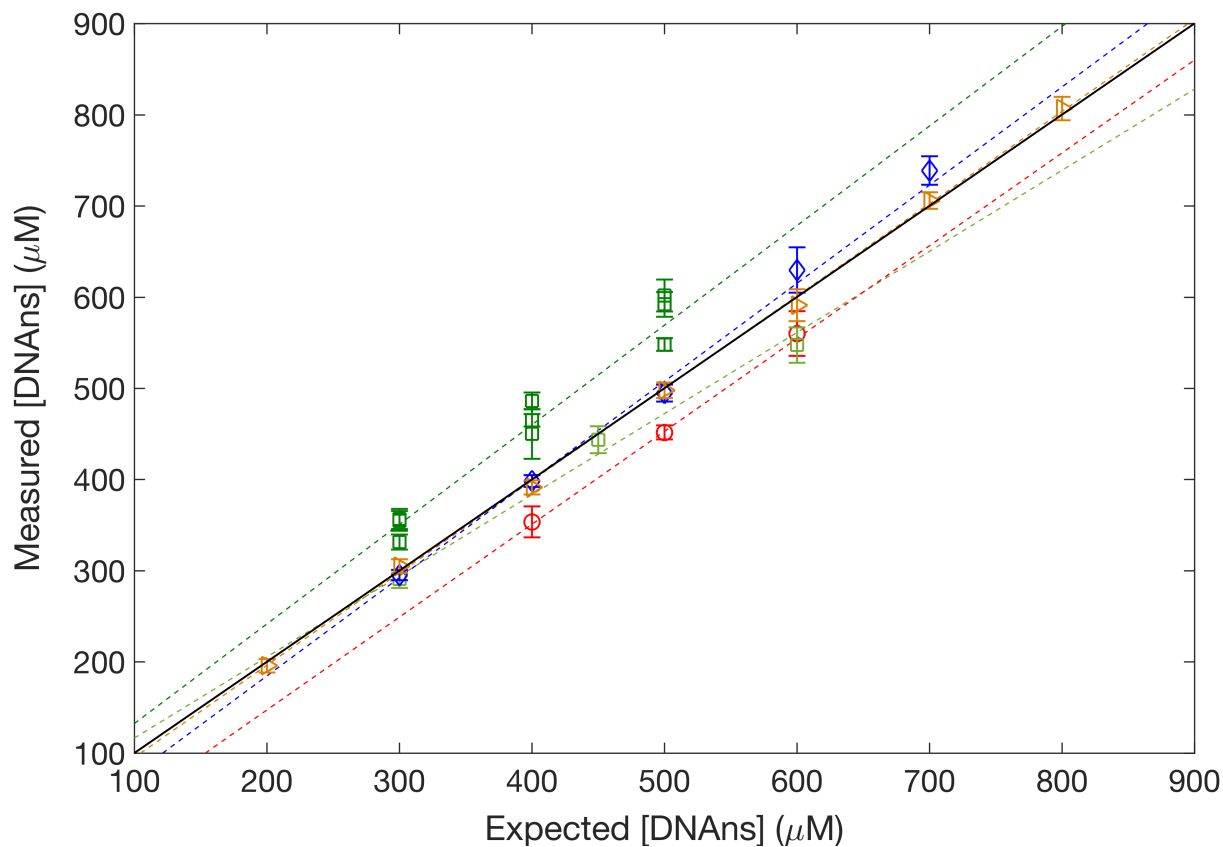


Fig. S4. Plot of measured [DNAns], with applied correction factors shown in Table S1, as a function of expected [DNAns]. The orange, blue, green, and red denote $f = 3, 4, 5,$ and $6,$ respectively. The $f = 5$ samples colored in the lighter shade of green were made several months after the first batch of $f = 5$ samples (dark green) in order to reproduce mechanical measurements. The dashed lines denote linear fits to the data. The solid, black line denotes the relationship Measured [DNAns] = Expected [DNAns] (*i.e.*, $f(x) = x$).

f	$A260_{1\%}^*$	$100 \cdot A260_{0.01\%}^*$	α	Correction to [DNAns]
3	9.4 ± 0.1	11.5 ± 0.3	$30 \pm 6\%$	1.23 ± 0.02
4	13.0 ± 0.2	14.0 ± 0.4	$72 \pm 7\%$	1.09 ± 0.02
5	7.1 ± 0.2	8.5 ± 0.5	$39 \pm 10\%$	1.19 ± 0.04
$\times 5$	12.6 ± 0.1	13.0 ± 0.1	$89 \pm 20\%$	1.03 ± 0.05
6	15.6 ± 0.6	15.8 ± 0.3	$96 \pm 10\%$	1.01 ± 0.03

Table S1. Table of $A260_{1\%}^*$ and $100 \cdot A260_{0.01\%}^*$ values as a function of f , along with the calculated percentage of ssDNA present at 1%, α , and the subsequent correction to [DNAns] at 1%. \times This batch of $f = 5$ samples was made several months after the first batch of $f = 5$ samples in order to reproduce mechanical measurements.

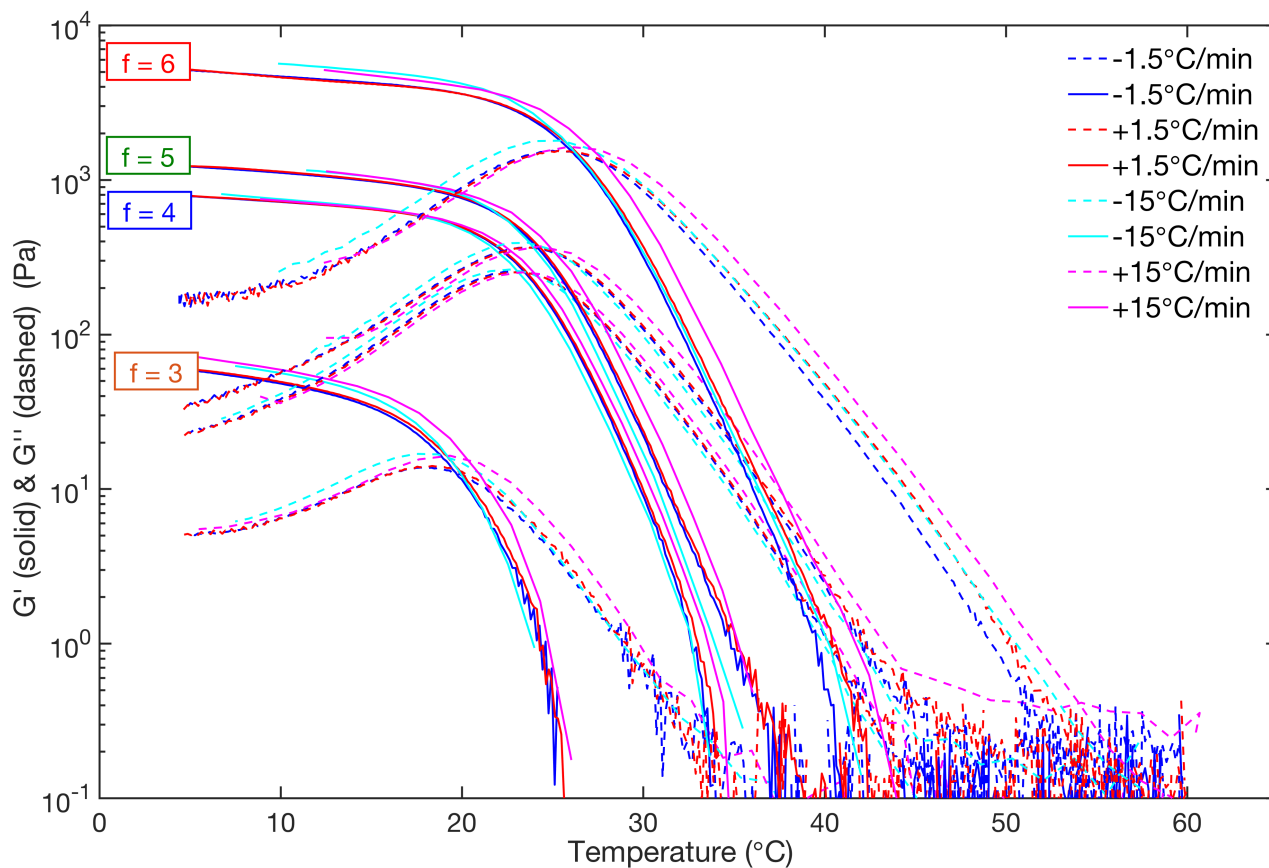


Fig. ss. Log-lin plot of G' and G'' vs. temperature at $\omega = 10$ rad/s and $\gamma = 5\%$ for different rates of cooling and heating (labeled in plot) of a (190 ± 10) μM 3-arm, (290 ± 10) μM 4-arm, (290 ± 10) μM 5-arm, and (360 ± 10) μM 6-arm DNAs solution. The overlap of the curves at different heating and cooling rates show that the viscoelastic response of the DNAs solution is independent of the rate of cooling and heating.

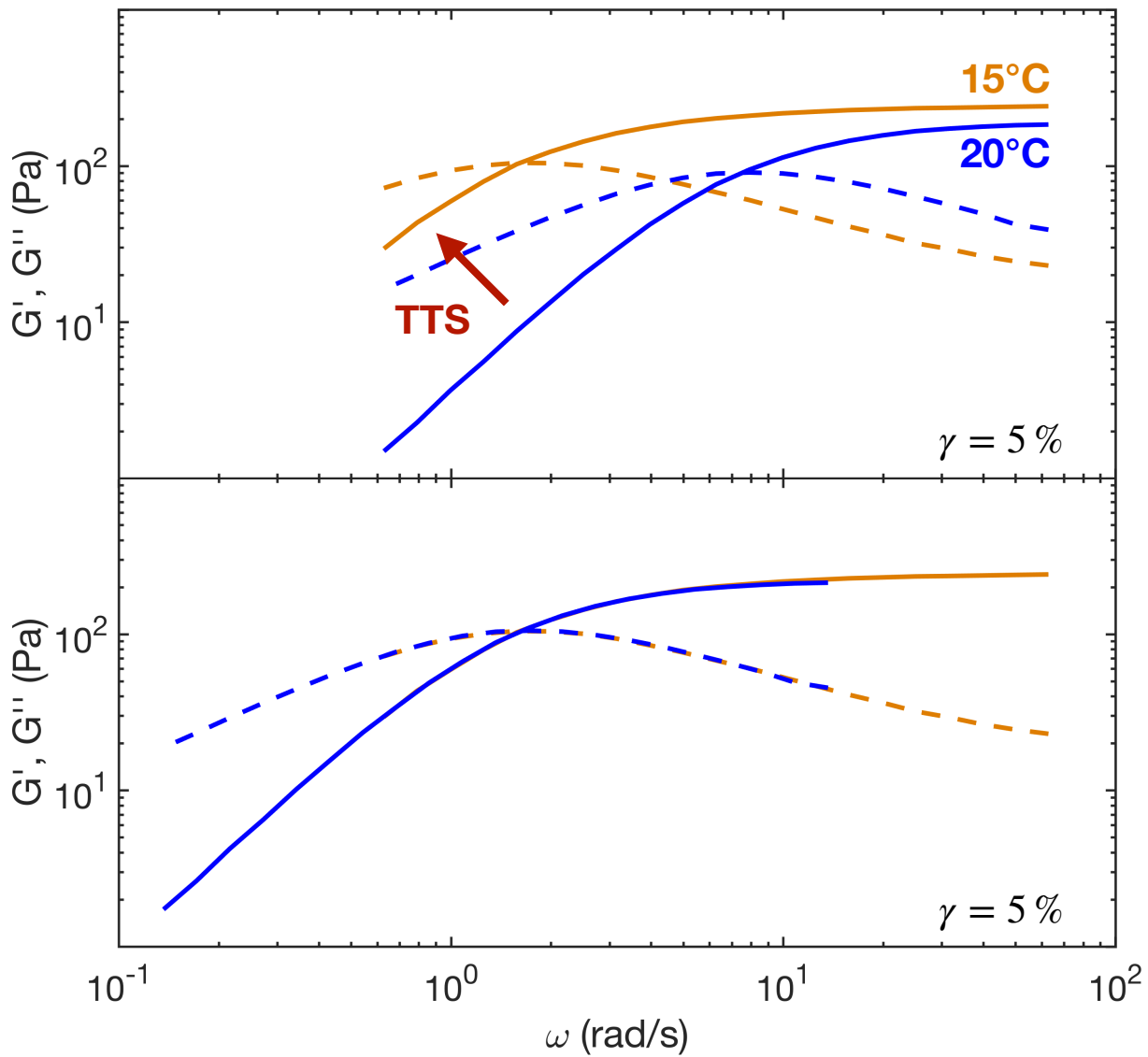


Fig. S6. The top and bottom panels are a log-log plot of $G'(\omega)$ and $G''(\omega)$ pre- and post-TTS, respectively. In this example, the reference temperature is $T_{ref} = 15^\circ C$. As a note, the data shown is a subset of the $f = 3$ DNAs $G'(\omega)$ and $G''(\omega)$ data.

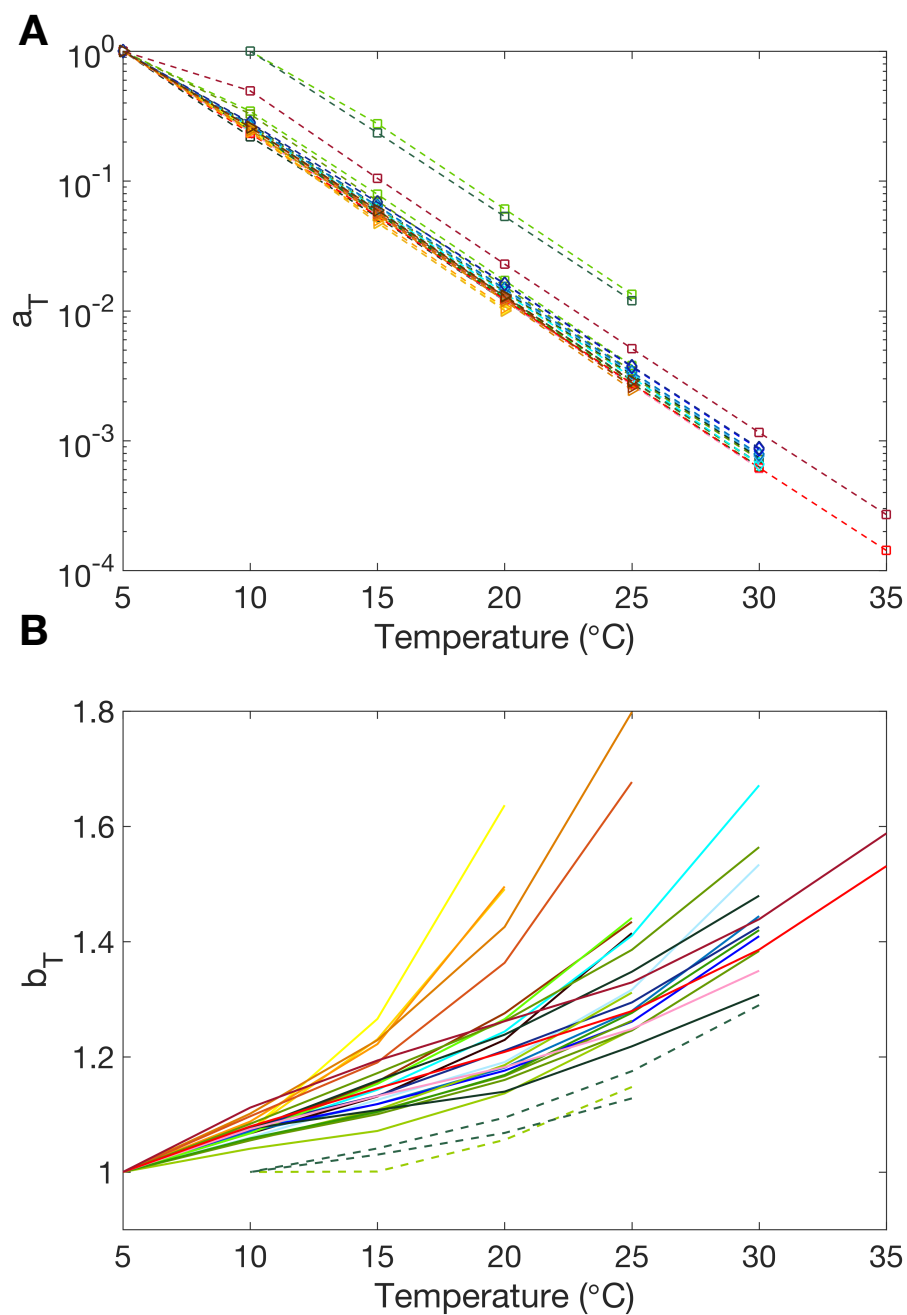


Fig. S7. Log-lin plot of the $T_{\text{ref}} = 5^\circ\text{C}$ (A) horizontal shift factors, a_T , and (B) vertical shift factors, b_T , as a function of temperature. $f = 3, 4, 5$ and 6 are denoted by yellow triangles, blue diamonds, green squares, and red circles, respectively. The lines are guides for the eyes. The darker colored lines signify a higher [DNAs]. The dashed lines denote $f = 5$ samples that were not cooled to $T = 5^\circ\text{C}$.

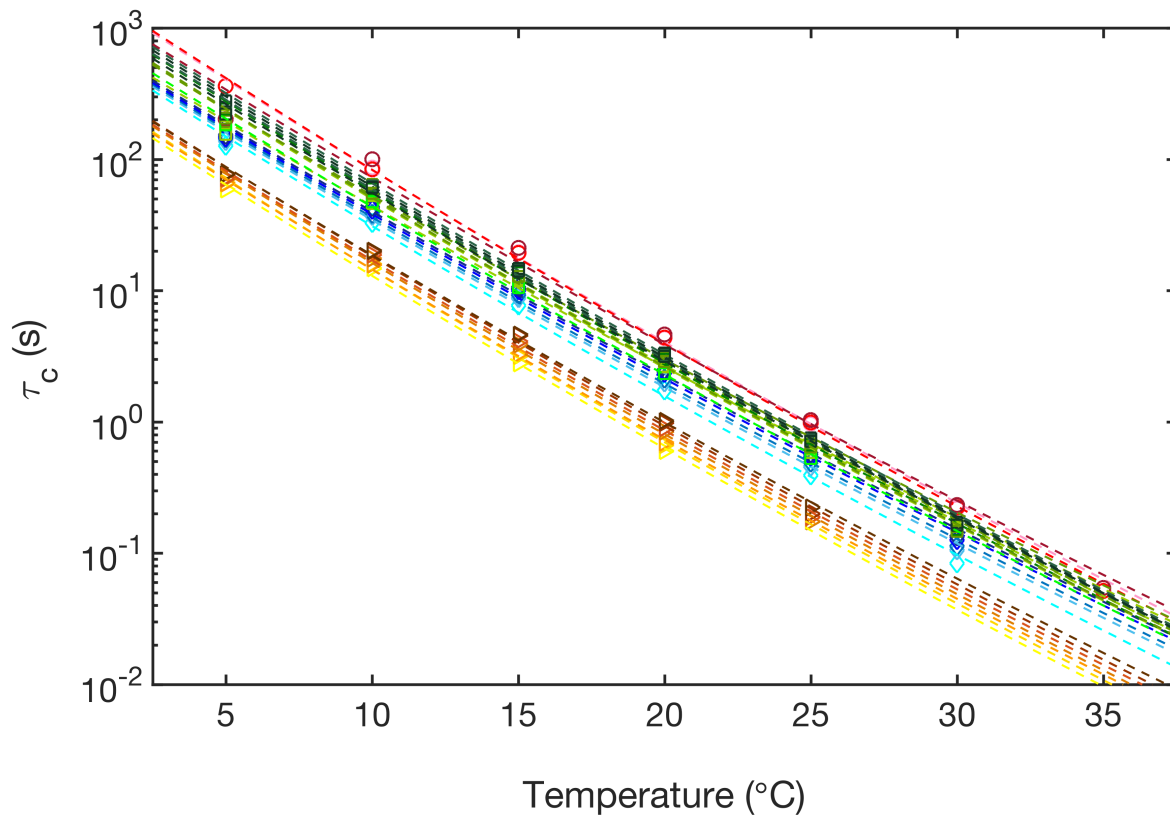


Fig. s8. Log-lin plot of the relaxation time, τ_c , as a function of temperature at $\gamma = 5\%$ for $f = 3$ (yellow triangles), 4 (blue diamonds), 5 (green squares), and 6 (red circles). Darker colors signify a higher [DNAns]. The dashed lines are the Arrhenius fits, $\tau_c \propto \exp(E_A/RT)$.

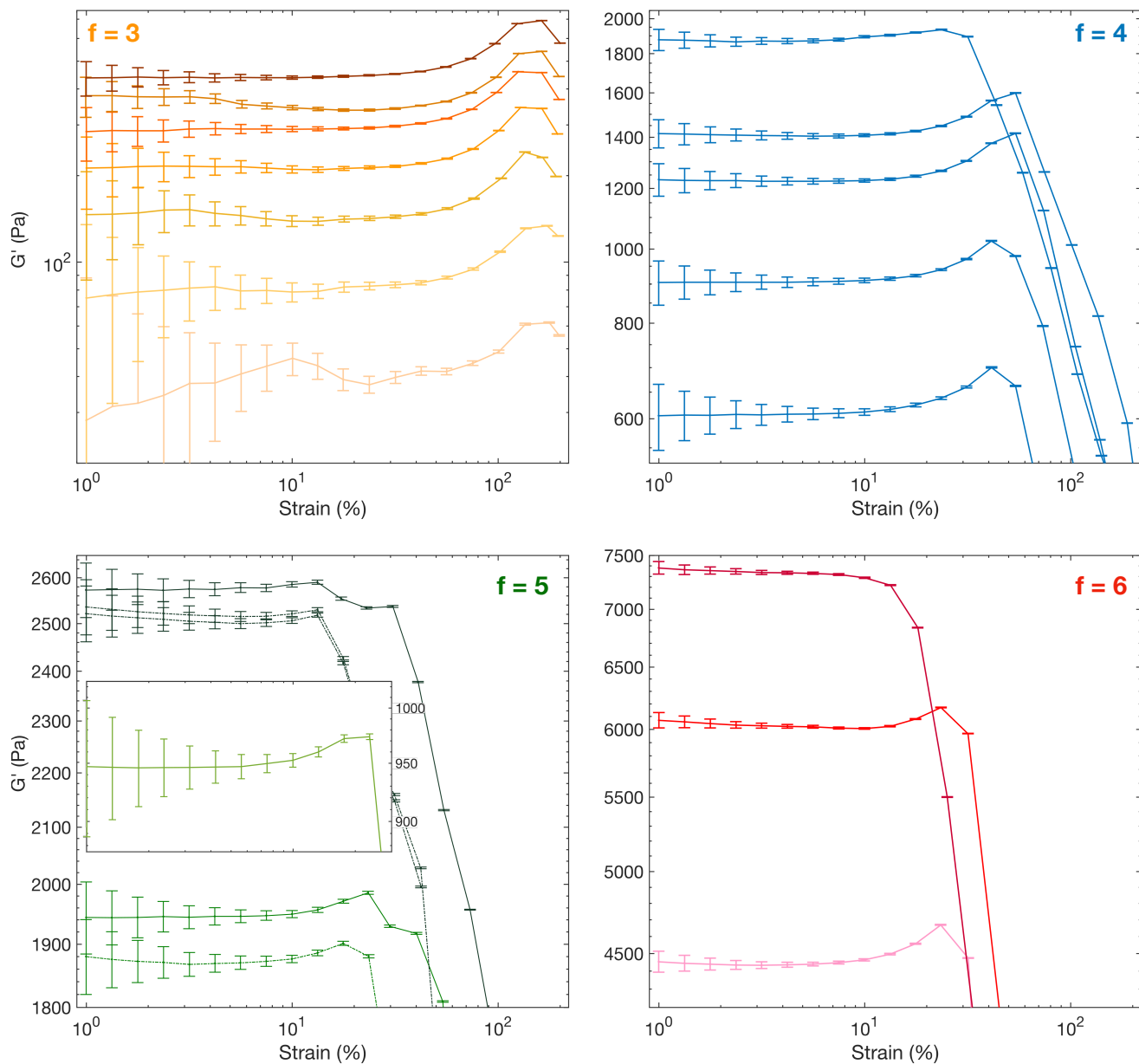


Fig. s9. Log-log plot of the storage modulus, G' , as a function of strain, γ , at a temperature of $T = 20^\circ\text{C}$ for $f = 3$ (top left panel), 4 (top right panel), 5 (bottom left panel), and 6 (bottom right panel). The curves with darker colors signify a higher [DNAs]. The solid and dash-dot lines denote data taken at a frequency of $\omega = 62.8$ rad/s and $\omega = 125.6$ rad/s, respectively. The inset of the $f = 5$ panel is the lowest concentrated $f = 5$ DNAs network, with $G' \approx 950$ Pa in the linear regime (scale on right of inset).

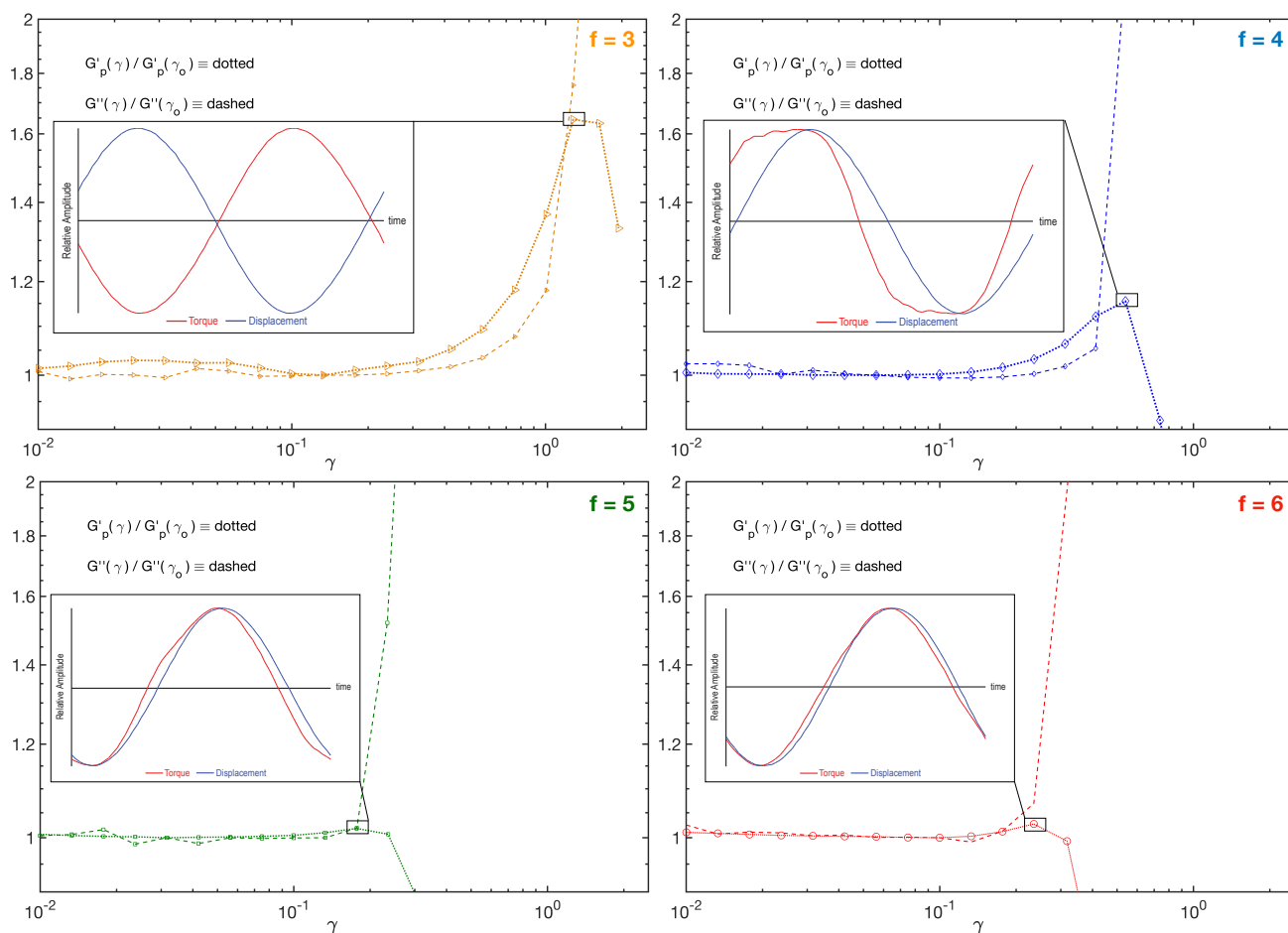


Fig. S10. Log-log plot of the normalized storage, $G'(\gamma)/G'(\gamma_0)$, and loss modulus, $G''(\gamma)/G''(\gamma_0)$ as a function of strain, γ , at a temperature of $T = 20^\circ\text{C}$ of $f = 3$ at $[\text{DNANs}] \approx 500 \mu\text{M}$ (top left panel), $f = 4$ at $[\text{DNANs}] \approx 500 \mu\text{M}$ (top right panel), the $f = 5$ gel at $[\text{DNANs}] \approx 450 \mu\text{M}$ (bottom left panel), and $f = 6$ gel $[\text{DNANs}] \approx 450 \mu\text{M}$ (bottom right panel). The insets of each panel show the measured relative amplitude of the torque (red) and radial displacement (blue) wave-forms as a function of time at the point of yielding.

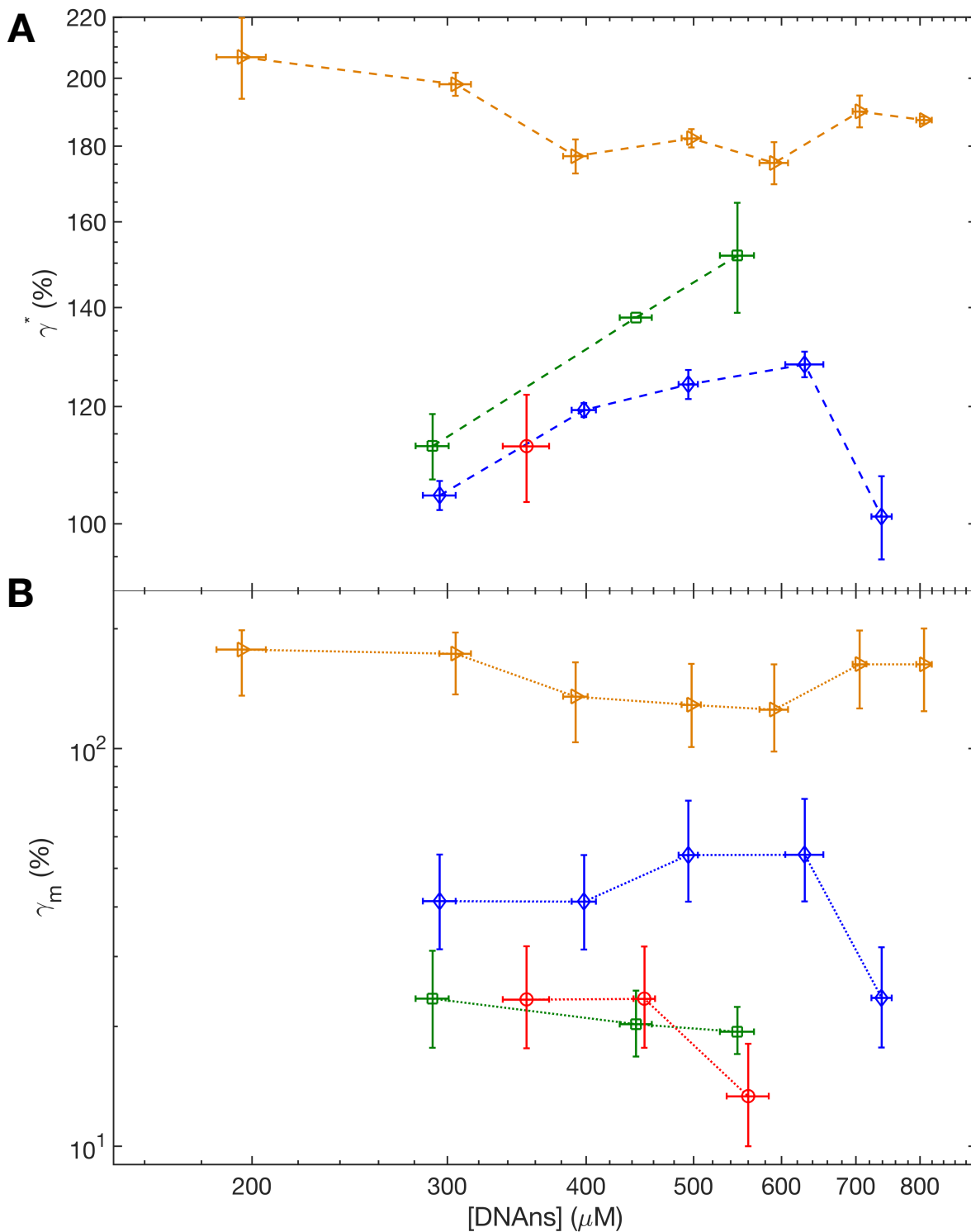


Fig. S11. **(A)** Log-log plot of γ^* , as a function of [DNAns] for $f=3$ (yellow triangles), 4 (blue diamonds), 5 (green squares), and 6 (red circles) at $T = 20^\circ\text{C}$ and $\omega \geq 62.8$ rad/s. **(B)** Log-log plot of γ_m , as a function of [DNAns] for $f=3$ (yellow triangles), 4 (blue diamonds), 5 (green squares), and 6 (red circles) at $T = 20^\circ\text{C}$ and $\omega \geq 62.8$ rad/s. In **(A)** and **(B)**, the dashed and dotted lines are guides for the eye.

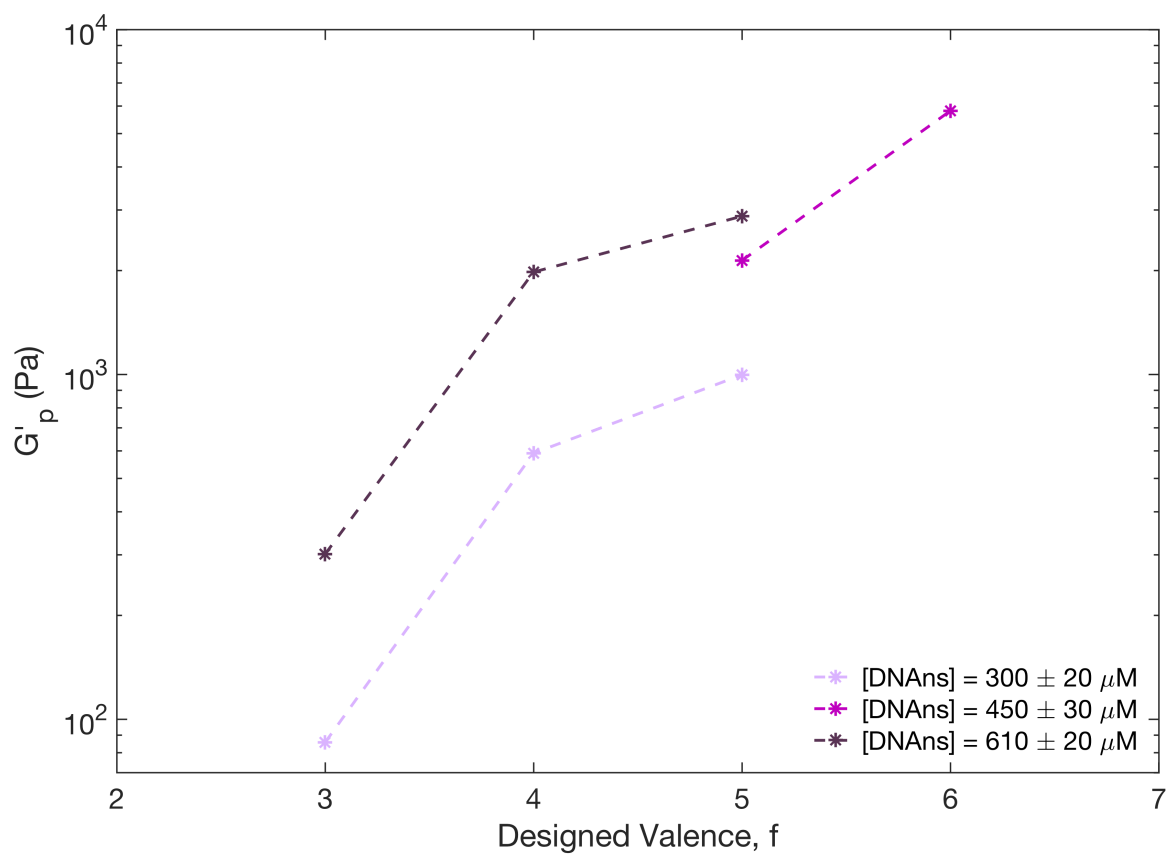


Fig. S12. Log-lin plot of the plateau modulus, G'_p , as a function of designed valence, f , at $T = 20^\circ\text{C}$ and $\gamma = 5\%$. The DNAns concentration of each data set is given in the legend.

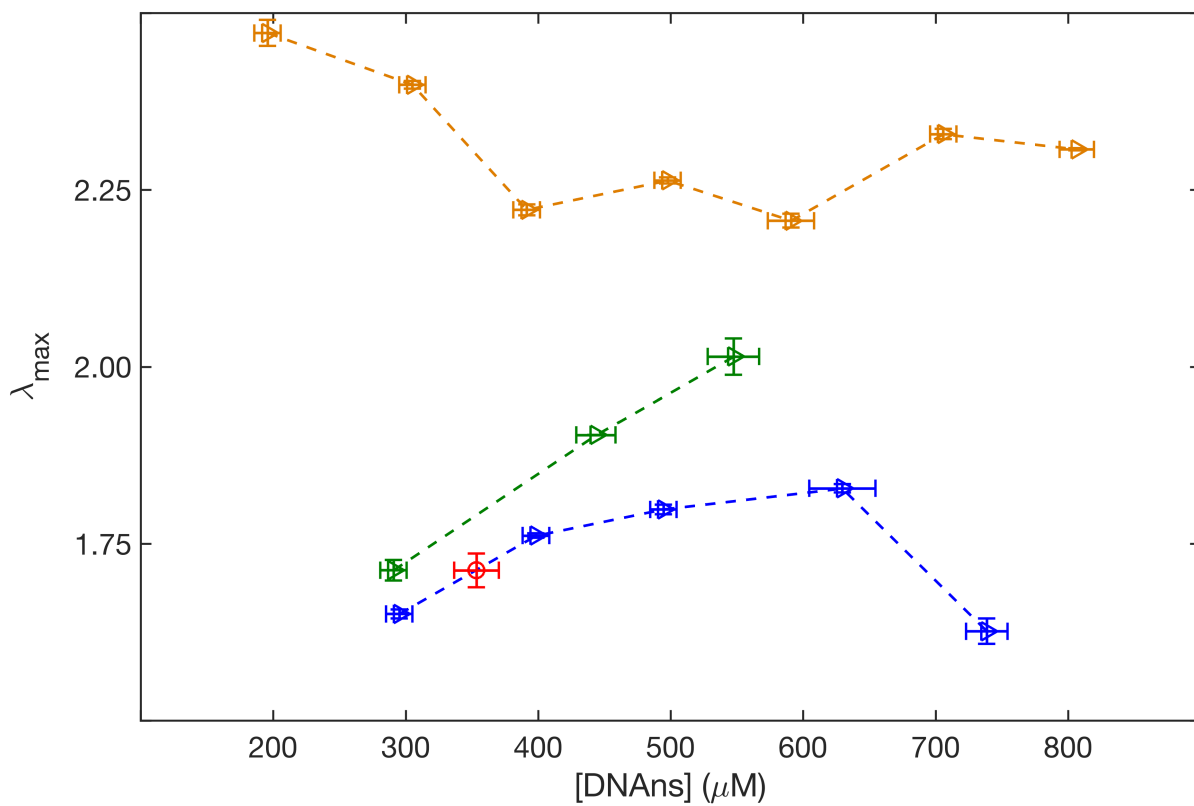


Fig. S13. Plot of λ_{max} , as a function of [DNAs] for $f = 3$ (yellow triangles), 4 (blue diamonds), 5 (green squares), and 6 (red circles) at $T = 20^\circ\text{C}$. The dashed lines are guides for the eye.

f	[DNAns] (μM)	ϕ_{rods}	ϕ_{sphere}	G'_p (Pa)
3	195	0.009	0.30	39.5
	305	0.015	0.47	85.6
	390	0.019	0.61	127
	500	0.024	0.77	215
	590	0.029	0.92	301
	705	0.034	1.09	388
	810	0.039	1.25	492
4	295	0.019	0.46	590
	400	0.026	0.62	915
	495	0.032	0.77	1240
	630	0.041	0.98	1460
	740	0.048	1.14	1980
5	290	0.023	0.45	996
	330	0.027	0.51	1370
	355	0.029	0.55	1380
	355	0.029	0.55	1400
	445	0.036	0.69	1940
	450	0.036	0.70	2140
	465	0.037	0.72	2185
	485	0.039	0.75	2140
	550	0.044	0.85	2630
	550	0.044	0.85	2800
	590	0.048	0.91	2820
	600	0.048	0.92	2870
6	350	0.034	0.55	4620
	450	0.044	0.70	5810
	560	0.054	0.87	7450

Table S2. Correspondence between [DNAns], DNAns volume fraction ϕ , and G'_p . ϕ is calculated in two ways: (i) ϕ_{rods} assumes the DNAns is composed of f rods, with each rod having a length given by the dsDNA arm length (8.5 nm, which includes half the sticky end) and radius (1 nm); and (ii) ϕ_{sphere} gives the volume fraction of a monodisperse system of spheres of concentration [DNAns], with the sphere radius set to the DNAns arm length (8.5 nm). ϕ_{sphere} is useful in gaining intuition on system structure through comparison to volume fractions of, e.g., FCC or random-sphere packings (12, 13).

139 **References**

- 140 1. Zadeh JN, et al. (2011) Nupack: analysis and design of nucleic acid systems. *Journal of Computational*
141 *Chemistry* 32(1):170–173.
- 142 2. Rovigatti L, Bomboi F, Sciortino F (2014) Accurate phase diagram of tetravalent dna nanostars. *The*
143 *Journal of Chemical Physics* 140(15):154903.
- 144 3. Nwokeoji AO, Kilby PM, Portwood DE, Dickman MJ (2017) Accurate quantification of nucleic acids
145 using hypochromicity measurements in conjunction with uv spectrophotometry. *Analytical Chemistry*
146 89(13567):13567–13574.
- 147 4. Markham NR, Zuker M (2005) Dinamelt web server for nucleic acid melting prediction. *Nucleic Acids*
148 *Research* 33(Web Server Issue):W577–W581.
- 149 5. Markham NR, Zuker M (2008) “UNAFold: software for nucleic acid folding and hybridization” in
150 *Bioinformatics, Volume II Structure, Function and Applications*. (Humana Press, New Jersey).
- 151 6. Yucht MG, Sheinmanb M, Broedersz CP (2013) Dynamical behavior of disordered spring networks.
152 *Soft Matter* 9(7000):7000–7006.
- 153 7. Zhang L, Mao X (2016) Finite-temperature mechanical instability in disordered lattices. *Physical*
154 *Review E* 93(022110):1–8.
- 155 8. Rocklin DZ, Hsiao LC, Szakasits M, Solomon MJ, Mao X (2018) Elasticity of colloidal gels: structural
156 heterogeneity, floppy modes, and rigidity. *arXiv* 1808.01533v1:1–7.
- 157 9. Seitz ME, et al. (2008) Fracture and large strain behavior of self-assembled triblock copolymer gels.
158 *Soft Matter* 5(2):447–456.
- 159 10. Erk KA, Henderson KJ, Shull KR (2010) Strain stiffening in sythetic and biopolymer networks.
160 *Biomacromolecules* 11(5):1358–1363.
- 161 11. Douglas JF (2009) Elasticity of networks with permanent and thermoreversible cross-links. *MRS*
162 *Proceedings* 1234:1234–QQ04–01.
- 163 12. Liu W, Jin Y, Chen S, Maksec HA, Li S (2018) Equation of state for random sphere packings with
164 arbitrary adhesion and friction. *Soft Matter* 13(421):421–427.
- 165 13. Ziff RM, Torquato S (2017) Percolation of disordered jammed sphere packings. *Journal of Physics A:*
166 *Mathematical and General* 50(8):1–12.