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# Supplemental Information

# High-Throughput Screening of Solid-State

Li-Ion Conductors

## Using Lattice-Dynamics Descriptors

Sokseiha Muy, Johannes Voss, Roman Schlem, Raimund Koerver, Stefan J. Sedlmaier, Filippo Maglia, Peter Lamp, Wolfgang G. Zeier, and Yang Shao-Horn



**Figure S1**. The mean absolute error of lithium-phonon band centers of the 53 compounds in Figure 2 as a function of the scaling parameter  $N_1$  and  $N_2$ . Related to Figure 2.



**Figure S2**. Comparison between the Li-phonon band centers of 71 compounds computed with *Quickhess* and full DFT calculations available in the Materials Project database. Related to Figure 2



**Figure S3.** The difference between phonon band centers computed with two different sets of scaling parameters (N<sub>1</sub> = 2; N<sub>2</sub> = 10 and N<sub>1</sub> = 5; N<sub>2</sub> = 10) for all of the ~1200 compounds retained after pre-screening. Related to Figure 2 and 3.



**Figure S4.** Histograms of a) Total-, b) Anion- and c) Lithium-phonon band centers for all chemistries excluding oxides. Related to Figure 3.



**Figure S5.** Correlation between the Li-phonon band centers and the average Li-anion bond length. Related to Figure 4.



**Figure S6**. Correlation between a) Total- and b) Anion-phonon band center with the anion pband centers. Hydrogen-containing compounds (hydrides, hydroxides and amides) were excluded due to the particular low mass of hydrogen and the difference in the nature of bonding which involves the s-orbital instead of p-orbitals. The inset shows the schematic of the anionprojected electronic DOS. The p-band centers were computed by integrating over the occupied states with respect to the top of the valence band indicated by the dashed line. Related to Figure 5.



**Figure S7**. Absence of correlation between Li-phonon band center and a) Total electronic band centers b) anion p-band center and c) Lithium electronic band center. Related to Figure 5 and 6.



**Figure S8.** Correlation between oxidation potential and a). Total-phonon band centers b). Lithium-phonon band centers. Correlation between reduction potential and c). Total-phonon band centers d). Anion-phonon band centers e). Lithium-phonon band centers. f). Correlation between stability window and band gaps. Related to Figure 6.



**Figure S9.** All compounds having stability window smaller than 3V and Li-phonon band center smaller than 35 meV along with their energy above hull. Compounds having non-zero energy above hull are thermodynamically unstable at 0K although they can be stabilized at higher temperature by entropic contributions. LGePS, LSiPS and LSnPS stand respectively for  $Li<sub>10</sub>GeP<sub>2</sub>S<sub>12</sub>$ ,  $Li<sub>10</sub>SiP<sub>2</sub>S<sub>12</sub>$  and  $Li<sub>10</sub>SnP<sub>2</sub>S<sub>12</sub>$ . Related to Figure 7.

**Table S1.** Compounds with Li-phonon band centers smaller than 35 meV along with their stability window, energy above hull and the material id as appeared in the Materials Project database. Related to Figure 7.

|                    |            | E hull      | Li-band centers | <b>Stability Window</b> |
|--------------------|------------|-------------|-----------------|-------------------------|
| Name               | MP id      | (eV/atom)   | (meV)           | (V)                     |
| Li2PbO3            | mp-22450   | 0           | 29.76           | 0.76                    |
| LiNbO3             | mp-755559  | 0.0466      | 32.297          | 2.14                    |
| LiTa3O8            | mp-559908  | 0.0139      | 30.778          | 2.45                    |
| LiGdO2             | mp-754204  | 0.0402      | 28.342          | 2.95                    |
| LiErO <sub>2</sub> | mp-10971   | 0.0026      | 33.518          | 3                       |
| LiTa3O8            | mp-7638    | $\mathbf 0$ | 30.288          | 2.44                    |
| LiNO3              | mp-8180    | $\Omega$    | 34.795          | 2.55                    |
| LiIO <sub>3</sub>  | mp-22955   | 0.0091      | 34.595          | 1.93                    |
| LiNb3O8            | mp-3368    | 0           | 31.571          | 1.66                    |
| LiNbO3             | mp-1078377 | 0.0203      | 28.772          | 2.13                    |
| LiTa3O8            | mp-757158  | 0.0384      | 24.783          | 2.46                    |
| LiNb(TeO4)3        | mp-758389  | 0.0015      | 33.295          | 0.33                    |
| LiGe2(PO4)3        | mp-541272  | 0.0015      | 27.54           | 1.58                    |
| LiP3W2O13          | mp-763452  | 0.0477      | 31.178          | 0.8                     |
| Ba4Li(SbO4)3       | mp-7971    | 0           | 26.219          | 1.78                    |
| Liln(IO3)4         | mp-973966  | 0           | 30.187          | 1.07                    |
| Lisn(PO3)3         | mp-684503  | 0.0361      | 28.535          | 0.5                     |
| LiBiP207           | mp-684109  | 0.0451      | 19.314          | 1.88                    |
| LiSb2P5O16         | mp-684104  | 0.0471      | 30.541          | 1.52                    |
| LiZr2(PO4)3        | mp-681439  | 0.0159      | 29.963          | 2.62                    |
| Li2P5WO15          | mp-763414  | 0.0302      | 31.429          | 0.31                    |
| Li2Sn(PO3)4        | mp-26855   | 0.0416      | 26.749          | 0.52                    |
| Lisn(PO3)3         | mp-26897   | 0.0406      | 27.292          | 0.5                     |
| LiInGeO4           | mp-17854   | 0           | 34.065          | 1.79                    |
| LiP2WO8            | mp-763370  | 0.0488      | 14.895          | 0.8                     |
| Lisn(PO3)3         | mp-758086  | 0.0464      | 31.686          | 0.52                    |
| LiZr2(PO4)3        | mp-773068  | 0.0111      | 29.848          | 2.62                    |
| LiBiP2O7           | mp-672979  | 0.0196      | 28.165          | 1.86                    |
| LiSn2(PO4)3        | mp-562473  | 0.0053      | 22.662          | 1.69                    |
| Li2Pr(NO3)5        | mp-555979  | 0           | 31.056          | 3.54                    |
| LiSbWO6            | mp-19694   | 0.0368      | 34.053          | 1.42                    |
| LiGa(WO4)2         | mp-19695   | 0.0218      | 29.531          | 1.31                    |
| LiSi6Bi9O26        | mp-757434  | 0.0381      | 23.271          | 0.99                    |
| LiB(S2O7)2         | mp-1020060 | 0           | 31.674          | 1.53                    |
| LiSi6Bi9O26        | mp-772721  | 0.0282      | 17.502          | 0.97                    |
| Li2TIPCO7          | mp-754034  | 0.0371      | 30.574          | 0.48                    |
| LiNbSn3(PO4)6      | mp-767091  | 0.0091      | 19.654          | 1.66                    |
| K2LiTa6(PO8)3      | mp-684817  | 0           | 33.994          | 2.5                     |
| LiMgAl3(SO4)6      | mp-677680  | 0.017       | 34.913          | 1.64                    |
| Li3MgPCO7          | mp-768190  | 0.0455      | 34.902          | 2.03                    |
| LiCa9Mg(PO4)7      | mp-686484  | 0           | 18.274          | 3.05                    |
| LiCuS              | mp-753371  | 0.0456      | 33.26           | 0.01                    |
| LiCuS              | mp-766467  | 0.0361      | 33.657          | 0.01                    |
| LiSmS2             | mp-34477   | 0           | 15.743          | 1.08                    |
| LiNdS2             | mp-36791   | 0.0204      | 15.087          | 1.02                    |





### **Transparent methods**

Database and screening criteria

The structure of all compounds in our study were sourced from the Materials Project database(Jain et al., 2013) containing  $\sim$  14,000 Li-containing compounds using the API functionality(Ong et al., 2015) implemented in the *Pymatgen* software package.(Ong et al., 2013) Four criteria were used to pre-screen all these compounds before proceeding to the computation of phonon band centers: (1) The stability of each compound was assessed based on their energy above hull,  $E_{\text{hull}}$  which is the distance between the point representing the compound in the phase diagram and the convex hull. (Ong et al., 2010; Ping Ong et al., 2008) Materials having  $E_{\text{hull}} > 0$ are thermodynamically unstable at 0 K. However, if  $E_{\text{hull}}$  is not too large, the material can still be

stable at finite temperatures due to the entropic contribution to the free energy. We set the threshold of *E*hull to 50 meV/atom based on the observation that e.g. LGPS (space group *P42mc*) has  $E_{\text{hull}} = 30$  meV/atom but can still be synthesized and is stable RT.(Kamaya et al., 2011) (2)

The next criterion was to have a computed band gap of at least 1 eV in order to filter out compounds with too high electronic conductivity to be used as electrolytes. (3) We included only ternary, quaternary and quintenary compounds and excluded the binary compounds due to their rather restricted chemical space. (4) Finally, we excluded all compounds containing 3d transition elements, platinum-group elements and elements in the actinide series, due to concerns about cost, toxicity and internal redox-activity. After this pre-screening,  $\sim 1200$  compounds were retained, and their phonon band center computed.

#### Detailed methodology of *Quickhess* method

The main idea of this method is that the eigenvalues of the Hessian matrix can be estimated with a single DFT force projection if reasonably good approximations to the eigenvectors are available. In the original implementation (Voss and Vegge, 2008) approximate eigenvectors were constructed from a point charge force field evaluated with Ewald summation and thus with negligible effort compared to a DFT calculation, and point charges were optimized to reproduce the DFT ground-state ionic positions. Here, instead of using Ewald summation, we screen the Coulomb interaction:  $V_{\text{screen}} \sim \exp(-Kr)/r$  with  $K=1/2\text{\AA}^{-1}$  and sum in real space only. Charges are

not fitted but we simply use Bader charges. Neither of the two simplifications changes the trends in approximated phonon band centers. Following Voss *et al.* (Voss and Vegge, 2008), the Hessian matrix of this purely electrostatic force field is diagonalized to obtain the eigenvalues and

eigenvectors (displacements)  $u_i$ . From these eigenvectors, a special displacement  $w$  is constructed simply as a sum of all  $u_i$ . The more accurate Hessian **H** eigenvalues  $h_i$  can be obtained by projecting the Hessian of a more accurate DFT Hamiltonian onto this displacement:

$$
h_i \approx \underline{u}_i^{\mathrm{T}} \underline{H} \underline{w} \tag{S1}
$$

The projection of the Hessian  $H$  onto  $\overline{w}$  is approximated as the difference between the force on a new displaced configuration of the initial structure which is generated by applying the displacement  $w$  to the structure and the force on the initial structure:

$$
H\underline{w} \approx \frac{1}{\varepsilon} \left[ \underline{\nabla} E \left( \underline{r}_o + \varepsilon \underline{w} \right) - \underline{\nabla} E \left( \underline{r}_o \right) \right]
$$
 (S2)

Where E is the (DFT) energy and  $r<sub>o</sub>$  are the ground state coordinates and  $\varepsilon$  is a parameter to scale the displacement for the finite difference as defined in equation (2). If the initial configuration is a well relaxed structure, then the term  $\nabla E(r_{o})$  can be neglected. The parameter  $\varepsilon$  must be chosen carefully such that the induced force is well above numerical noise in the DFT calculations but also not too large potentially invalidating the harmonic approximation.

#### **Synthesis and characterization of Li3ErCl6**

The preparations and sample treatment for all compounds were carried out under argon atmosphere. Erbium chloride (ErCl3, anhydrous, 99.9 %, Alfa Aesar) and lithium chloride (LiCl, ultra-dry, 99.9 %, Alfa Aesar) were mixed in the appropriate stoichiometric molar ratios, while adding approximately 2 wt% excess ErCl<sub>3</sub> relative to the total weighted mass to counter the loss of ErCl3 during the hand grinding procedure. The mixture was then hand ground in an agate mortar to homogenize the powders and afterwards loaded into ball mill vials with a 30:1 mass ratio of the ZrO2 milling media (3 mm in diameter) to the starting powder. The precursors were milled for 297 total cycles, while one cycle consists of 5 minutes of milling, followed by 15 minutes of rest time. Additionally, after every 99th cycle, the cups were opened and the powder homogenized manually

again to reduce the loss of sample on the inner walls of the vials. Crystallization of the as-prepared amorphous compounds were achieved by subsequent annealing for 1 h at 550 °C.

X-ray powder diffraction was carried out using an Empyrean powder diffractometer (PANalytical, Netherlands) with CuKa radiation ( $\lambda_1 = 1.54051 \text{ Å}$ ,  $\lambda_2 = 1.54051 \text{ Å}$ ) in Bragg–Brentano  $\theta-\theta$ geometry, and a PIXcel3D area detector with 255 measuring channels. Powder samples were placed on (911)-oriented silicon zero background holders that were sealed with Kapton foil under an argon atmosphere. Patterns were collected in the range of 10–90° 2*θ* with a step size of 0.026 and an exposure time of 150 s per step for the ball milled sample and 425 s for the subsequently annealed sample, to achieve a sufficiently high intensity and therefore acquire reliable lattice parameter values. Pawley fits(Pawley, 1981) were performed using the TOPAS-Academic V6 software package (Coelho, 2018). A manual background using 62 points was used, whereas the peak profiles were described by a Thompson–Cox–Hastings pseudo-Voigt function.

The ionic conductivities were measured by AC impedance spectroscopy, using pellets (1.4 mm average thickness,  $>80\%$  density for the amorphous compound ( $>75\%$  for the crystalline compound)) with vapor deposited gold layers (200 nm on each side; surface area of 0.53 cm<sup>2</sup>). Electrochemical impedance analysis (EIS) was conducted in the temperature range of -40 °C to 60 °C using a VMP300 impedance analyzer (Biologic) at frequencies from 7 MHz to 100 mHz with an amplitude of 10 mV. All measurements were performed in pouch cells under an argon atmosphere.

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