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Supplementary Materials for

The cold Leidenfrost regime

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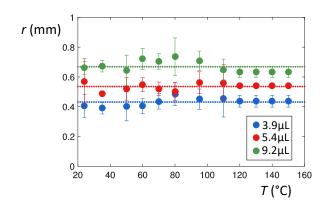


Fig. S1. Contact radius *r* of a water drop placed on a hot superhydrophobic solid, as defined in Fig. 2A. The substrate is Glaco-treated brass brought at temperature *T*. Three drop volumes are investigated: $\Omega = 3.9 \,\mu\text{L}$ (blue), $\Omega = 5.4 \,\mu\text{L}$ (red) and $\Omega = 9.2 \,\mu\text{L}$ (green). The contact radius is found to be roughly independent of the substrate temperature.

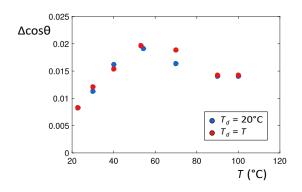


Fig. S2. Contact angle hysteresis $\Delta \cos\theta$ on Glaco-coated substrates as a function of T for drops having initially either a temperature $T_d = 20^{\circ}$ C (blue data) or the same temperature as the substrate ($T_d = T$, red data). Regimes of adhesion appear to be independent of the drop initial temperature.

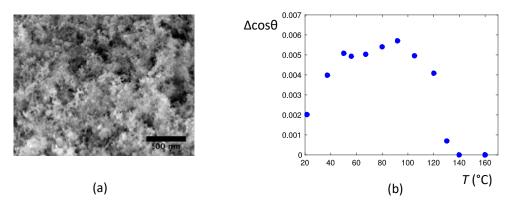


Fig. S3. Water adhesion on heated brass coated by a commercial colloidal repellent material (Ultra-Ever Dry, UltraTech International). (A) SEM picture of the surface of the samples (credits: Anaïs Gauthier). (B) Contact angle hysteresis $\Delta \cos\theta$ for a millimetre-sized water drop as a function of the substrate temperature *T*. We observe the same sequence of regimes as in figure 2c.

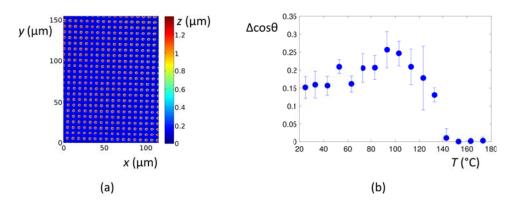


Fig. S4. Water adhesion on heated micrometric posts. (A) Profilometry of the substrate: a silicon wafer etched by Deep Reactive Ion Etching (DRIE) has micro-posts with diameter 2.6 μ m, spacing 6.25 μ m and height 1.2 μ m. (B) Contact angle hysteresis $\Delta \cos\theta$ as a function of temperature *T*. We observe the same regimes as in figures 2c, SI-2 and SI-3, even if adhesion is much larger here, owing to a stronger pinning on micropillars.

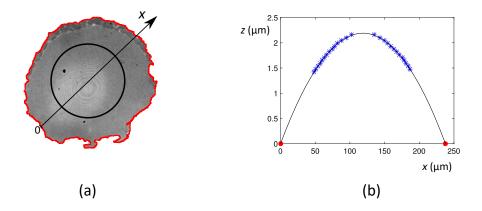


Fig. S5. Morphology of a vapor patch. (A) Fringes at the centre of the main vapor patch observed in figures 4b and 4c at $T = 75^{\circ}$ C. (B) We deduce from these fringes the location of the liquid/vapor interface (blue data), which appears to be nicely fitted by a circle (solid line) far from the contact line, in spite of pinning at the contact line. The intersection of this profile with the (red) contact line observed in (a) provides a vapor contact angle $\theta_v = 2^{\circ}$. This very low value implies a weak contact angle hysteresis $\Delta \cos \theta_v$, yet large enough to generate visible distortion of the contact line. A low θ_v is typical of a superhydrophobic material, and it explains that even a small quantity of vapor (here about 30 pL) can invade macroscopic distances (here about a quarter of a millimeter) at the solid/water interface.

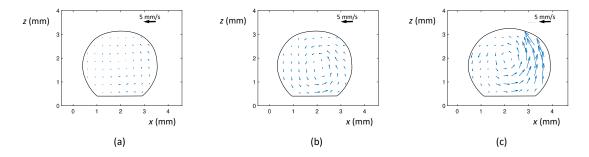


Fig. S6. Internal flow in water drops ($R \approx 1.5$ mm) placed on a hot superhydrophobic solid (Glaco-coated wafer). Observations are performed by PIV (particle image velocimetry) at different solid temperature. In all cases, an inner motion is observed and found to consist in a rolling motion, a symmetry breaking we attribute to confinement. The scale of the velocity vectors is the same in all pictures and given by the bold arrow showing 5 mm/s. (A) At $T=50^{\circ}$ C, the average and maximum measured velocities of water inside the drop are around 0.5 mm/s and 1 mm/s, respectively. (B) At $T=85^{\circ}$ C, the average and maximum velocities are around 1.1 mm/s and 2.5 mm/s. (C) At $T=107^{\circ}$ C, the average and maximum velocities are around 2.4 mm/s and 7.1 mm/s.