**Title:** A simple model for determining affinity from irreversible thermal shifts

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**Supplemental Information**

## 6 **Extended derivations**

$$
7 \quad \Delta G_L = -RTln(Q) \tag{1}
$$

8 
$$
Q = \frac{[F]}{[F]} + \frac{[FL]}{[F]} = 1 + \frac{[L]}{K_D}
$$

$$
9 \quad \Delta G_L = RTln\left(1 + \frac{[L]}{K_D}\right) \tag{2}
$$

10 where *T* is in Kelvin, and is fixed at the temperature at which the KD was determined (*eg*, 298 K).

*ΔL Gu - ΔGu = Δ<sup>u</sup>* 11 *GL - ΔGL* (3)

$$
12 \qquad k = Ae^{-Ea_1/RT} \tag{4}
$$

$$
13 \t Ea = \Delta G + \Delta G^+ \t (5)
$$

$$
14 \t A1 \cong A2
$$
\t(6)

$$
15 \quad \Delta TM = TM_{,bound} - TM_{,apo} \tag{7}
$$

$$
16 \qquad kU_I = A_I e^{\left(\frac{-E a_1}{RT m_{apo}}\right)} \tag{8}
$$

$$
17 \qquad kU_2 = A_2 e^{\left(\frac{-E a_2}{R(Tm_{\text{apo}} + \Delta T m)}\right)} \tag{9}
$$

18 ∵ *kU1* at *TM,apo* is approximately equal to *kU2* at *TM,bound* (see equations (25-28), **SI Fig. 1** and

$$
19 \qquad \textbf{SI Fig 2)}
$$

20 
$$
I = \frac{kU_1}{kU_2} = \frac{A_1}{A_2} \frac{e^{\left(\frac{-Ea_1}{RTm, apo}\right)}}{e^{\left(\frac{-Ea_2}{R(Tm, apo + \Delta Tm)}\right)}} = e^{\left(\frac{Ea_2}{R(Tm, apo + \Delta Tm)}\right) - \left(\frac{Ea_1}{RTm, apo}\right)}
$$
(10)

21 
$$
\therefore
$$
  $ln(e^{(x)}) = x$  and  $ln(1) = 0$ 

$$
22 \qquad \frac{Ea_2}{R(Tm_{apo} + \Delta Tm)} = \frac{Ea_1}{RTm_{apo}} \tag{11}
$$

$$
23 \t Ea_2 = \frac{Ea_1R(Tm,_{apo}+ATm)}{RTm,_{apo}} = Ea_1\left(1 + \frac{\Delta Tm}{Tm,_{apo}}\right) = Ea_1 + Ea_1\left(\frac{\Delta Tm}{Tm,_{apo}}\right)
$$
(12)

$$
24 \quad \Delta G_v = E a_l - \Delta G v_r^{\dagger} \tag{13}
$$

$$
25 \t\t ALGU = Ea2 - ALGU \dot{\tau}
$$
\t(14)

26 Substitute equation (13) and (14) into equation (3):

$$
27 \t(Ea_2 - \Delta^L G_{U_T}^+) - (Ea_1 - \Delta G_{U_T}^+) = \Delta^U G_L - \Delta G_L \t\t(15)
$$

28 Substitute equation (12) into equation (15) and simplify:

$$
29 \t E a_l \left( \frac{\Delta T m}{T m_{\text{apo}}} \right) + \Delta G_{U_T} t - \Delta^L G_{U_T} t = \Delta^U G_L - \Delta G_L \t (16)
$$

$$
30 \quad \Delta G_{U_T}{}^{\dagger} \cong \Delta^L G_{U_T}{}^{\dagger} \tag{17}
$$

$$
31 \t E a_l \left( \frac{\Delta T m}{T m_{\text{apo}}} \right) = \Delta^{U} G_L - \Delta G_L \t (18)
$$

## 32 Since ligand does not bind to unfolded protein:

$$
43 \t E a_l \left(\frac{\Delta T m}{T m_{\text{app}}} \right) = -\Delta G_L \tag{19}
$$

## 34 Substitute equation (2) into equation (19), followed by rearrangement gives:

$$
35 \qquad \qquad -RTln\left(1 + \frac{[L]}{K_D}\right) = RTln\left(\frac{K_D}{K_D + [L]}\right) = -Ea_l\left(\frac{\Delta Tm}{Tm_{\text{apo}}}\right) \tag{20}
$$

36 Experimentally, the apparent ΔTM (*ΔTM,APP*), will be the sum of the ligand binding to the site of 37 interest (ΔTM), plus any additional extra-site binding (*ΔTMX*):

38 *ΔTM,APP=ΔTM + ΔTMX* (21)

39 Combining equations (20) and (21) we find:

40 
$$
\Delta TM_{APP} = -\left(\frac{Tm_{Appo}RT}{Ea_1}\right)ln\left(\frac{K_{D,app}}{K_{D,app}+[L]}\right) = -\left(\frac{Tm_{Appo}RT}{Ea_1}\right)\left(ln\left(\frac{K_D}{K_D+[L]}\right) + ln\left(\frac{K_{Dx}}{K_{Dx}+[L]}\right)\right) \tag{22}
$$

41 where KD,APP is the apparent affinity, and KD<sub>X</sub> is the affinity of extra-site binding. After 42 rearrangement and simplification we find:

43 
$$
ln\left(\frac{K_D}{K_D+[L]}\right) = -\left(\frac{Ea_1}{RT}\right)\left(\frac{\Delta Tm_{APP}}{Tm_{aapo}}\right) - ln\left(\frac{K_{Dx}}{K_{Dx}+[L]}\right)
$$
(23)

44 where *T* is the temperature at which the KD was determined (*eg*, 298 K).

45 
$$
\therefore \qquad \ln\left(\frac{K_D}{K_D+[L]}\right) \cong \ln\left(\frac{K_D}{[L]}\right) \text{ when } [L] \gg KD
$$

$$
46 \qquad ln\left(\frac{K_D}{[L]}\right) = -\left(\frac{Ea_1}{RT}\right)\left(\frac{\Delta Tm_{,APP}}{Tm_{,apo}}\right) \qquad ln\left(\frac{K_{Dx}}{K_{Dx}+[L]}\right) \tag{24}
$$

47

# 48 **Extended derivations for value of** *kU1* **at** *TM,apo* **and** *kU2* **at** *TM,bound*

49 Beginning at equation (10), without assuming the ratio of the unfolding rate constants of apo and

50 bound is unity,  $\frac{k u_1}{k u_2}$  =1, substitution with equation (7) will yield:

$$
51 \qquad \frac{Ea_2}{RTm_{bound}} = \frac{Ea_1}{RTm_{app}} + ln\left(\frac{ku_l}{ku_2}\right) \tag{25}
$$

### 52 Rearrangement of equation (25) gives:

$$
53 \t Ea_2 = \frac{Ea_1 Tm, bound}{Tm, apo} + RTm, bound \ln\left(\frac{ku}{ku}\right) \t (26)
$$

54 From Sanchez-Ruiz *et al.* (Sanchez-Ruiz et al. 1988) we get the equivalence:

55 
$$
\frac{vEa_1}{RTm_{,apo}^2} = kU_I = A_I e^{-Ea_1/RTm_{,apo}} \tag{27}
$$

56 where  $\nu$  is the heating rate of the irreversible unfolding experiment. If the apo and the bound 57 experiments are collected at the same heating rate, and  $A_1 \cong A_2$ , the ratio of the rate constants at 58 their TM temperatures is:

$$
59 \qquad \frac{k u_l}{k u_2} = \frac{\left(\frac{E a_1}{R T m_a p_o^2}\right)}{\left(\frac{E a_2}{R T m_b o u n d^2}\right)} = \frac{\left(\frac{E a_1}{R T m_a p_o^2}\right)}{\left(\frac{E a_1 T m_b o u n d}{R m_a p_o} + R T m_b o u n d \ln\left(\frac{k u_l}{k u_2}\right)\right)} = \frac{E a_1 T m_b o u n d}{E a_1 T m_a p_o + R T m_a p_o^2 \ln\left(\frac{k u_l}{k u_2}\right)}\tag{28}
$$

At the condition of  $TM$ , apo =  $TM$ , bound,  $\Delta TM = 0$ ,  $\frac{k u_1}{k u_2} = 1$ , therefore  $\ln \left( \frac{k u_1}{k u_2} \right)$ 60 At the condition of  $TM$ , apo =  $TM$ , bound,  $\Delta TM = 0$ ,  $\frac{k_{\text{u}_1}}{k_{\text{u}_2}} = 1$ , therefore  $\ln\left(\frac{k_{\text{u}_1}}{k_{\text{u}_2}}\right) = 0$ . Starting at  $\Delta TM = 0$ , steps can be taken through a range of  $\Delta TM$  values using an estimate of  $\ln \left( \frac{k u_1}{k u_2} \right)$ 61 steps can be taken through a range of  $\Delta TM$  values using an estimate of  $\ln\left(\frac{R_{\text{U}_1}}{R_{\text{U}_2}}\right)$  from the prior  $\Delta TM$ step (*eg*, to approximate  $\frac{k u_l}{k u_2}$  at  $\Delta T M = 1$ , the ln  $\left(\frac{k u_1}{k u_2}\right)$ 62 step (eg, to approximate  $\frac{k u_1}{k u_2}$  at  $\Delta T M = 1$ , the  $\ln \left( \frac{k u_1}{k u_2} \right)$  value from  $\Delta T M = 0$  is used). In this way it can 63 be demonstrated that under typical protein unfolding conditions (*eg*, TM between 50-70 °C), *kU1* 64 at *TM,apo* is approximately equal to *kU2* at *TM,bound* (**SI Fig 1**). Similarly, using a numerical 65 integration approach, a simulation of unfolding can be made which also shows *kU1* at *TM,apo* is 66 approximately equal to *kU2* at *TM,bound* (**SI Fig 2**). Please note that since refolding is negligible in 67 this model, the rates constants of folding are not considered; if you have been reading this 68 extension of the derivation as a loop-out from the main or supplemental text assertion of rate 69 equivalence of *kU1* at *TM,apo* and *kU2* at *TM,bound*, you may now return to the rest of the derivation.

70

#### 71 **Extended discussion on simulations**

 In our laboratory experiments, we observed a ligand-dependent ΔTM shift beyond what can be explained by initial occupancy. One possibility is that the effect is due to unfolding through the apo unfolding pathway, which favors net ligand dissociation; thus, the ligand-dependent ΔTM is due to kinetic competition between ligand rebinding to apo protein and unfolding through the 76 apo path. To test this hypothesis, the simulation was adjusted to allows us to predict what  $\Delta TM$ 77 results for the system when we have high initial occupancy of the bound form ( $\geq 90\%$ ), and then allow the system to evolve over time as the temperature increases. When no interconversion of 79 apo and bound is allowed, we see a  $\Delta TM$  that is the average of  $TM$ , apo and  $TM$ , bound values weighted by the initial proportions of apo and bound protein. This model represents the extreme of a very fast rate of rebinding or a very slow rate of dissociation, where no ligand-bound protein is ever lost to the apo unfolding pathway; we see that this system is saturable as the limit of 100-percent *TM,bound*, after which further ligand addition has no effect on ΔTM.

 In contrast, a model that allows instantaneous transfer between the bound form to apo, but no transfer of apo to the bound state (unidirectional replacement of every molecule of apo that unfolds with a molecule of bound), represents the limit of very fast dissociation with very slow rebinding; this is the most aggressive example of a system where no ligand rebinding is allowed to protect the protein from unfolding beyond the initial equilibrium distribution of bound protein. As with the prior model, we see a system that is strongly influence by the initial occupancy but 90 even this extreme has only a slightly reduced apparent  $\Delta TM$  compared to the  $\Delta TM$  of the fast rebinding model (**SI Fig. 3**).

 As can be seen in **SI Fig 3**, slow rebinding predicts a ΔTM that is very close to the first model as the initial occupancy approaches 100-percent bound. Therefore, both these models lack the non-saturable effect seen in experiments for increasing ligand concentration, suggesting this

 phenomenon is not due to rate competition between ligand binding and the apo unfolding pathway for irreversible unfolding.

 Lastly, these models begins with the equilibrium state, which then evolves with time and temperature. Based on these models, the initial equilibrium occupancy of the system seems to have a dominant influence on the final ΔTM. Inspection shows this is because ΔTM is equally dependent on the rate constants and the ligand occupancy (*eg*, *kU2*×*[FL]*); since the fold-change between the rate constants *kU1* and *kU2* at the same temperatures is small (*ca* 4-fold) (**SI Fig. 2**) compared to the fold-change in the concentration of apo and bound at 95-percent occupancy (*ca* 19-fold), most of the information for the resultant ΔTM is contained in the initial occupancy term.

## **References**

 Sanchez-Ruiz, J. M., J. L. Lopez-Lacomba, M. Cortijo, and P. L. Mateo. 1988. 'Differential scanning calorimetry of the irreversible thermal denaturation of thermolysin', *Biochemistry*, 27: 1648-52.



 **SI Fig 1.** The value of the apo unfolding rate constant, *kU1*, at *TM,apo* divided by the value of the bound unfolding rate constant, *kU2*, at *TM,bound*. For a two-state unfolding process, with unfolding 111 rates for apo ( $kU_1$ ) and bound ( $kU_2$ ) that are Arrhenius functions, where  $A_1 \cong A_2$ , the value of  $kU_1$  at *TMapo* is approximately the same value as *kU2* at *TM,bound*. Data were generated using equation (28), with *TMapo* values of 50, 60 and 70 °C (vantablack, ultramarine and porphyry).



 **SI Fig 2.** Simulated data showing unfolding and rate of unfolding. These data show that for a two- state unfolding process, with unfolding rates for apo (*kU1*) and bound (*kU2*) that are Arrhenius 117 functions, where  $A_1 \cong A_2$ , the value of  $kU_1$  at  $TM_{\text{app}}$  is approximately the same value as  $kU_2$  at *TM,bound*. Apo (black) and bound (green) unfolding profiles (left) or rate constant traces (right) are shown; red and cyan dashed lines have been added to guide the eye. Unfolding data were modeled using the methods described in equations (20-22) of the main text, with a data integration steps of 121 0.1 s, a ligand concentration of 100  $\mu$ M, and KD values of 0.005, 0.05, 0.5 and 5  $\mu$ M (top to bottom).



 **SI Fig 3.** Simulations showing the TM expected for 30 (left), 60 (middle) or 90-percent (right) initial occupancy when there is very fast rebinding (red line) or very slow rebinding (cyan line) relative to the apo unfolding rate. 100-percent apo (black) or bound (green) unfolding traces are 127 shown as reference. Data are for a heating rate of  $4 \degree C/\text{min}$ . An excel file programming these simulations is available to downloaded.