

# Proteoglycan degradation mimics static compression by altering the natural gradients in fibrillar organisation in cartilage - Supplementary Information

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## Abstract

We use the equations of mechanical equilibrium to derive an expression for the swelling pressure in hydrated cartilage. We consider a structural model composed of a pre-stressed fibrillar network inside a matrix.

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## 1. The swelling pressure

At equilibrium in free swelling conditions, the Cauchy (true) stress  $\boldsymbol{\sigma}$  in the tissue is zero:

$$\boldsymbol{\sigma} = -\Delta\Pi \mathbf{I} + (1 - \phi_F) \boldsymbol{\sigma}_M + \phi_F \boldsymbol{\sigma}_F = 0 \quad (1)$$

where  $\Delta\Pi$  is the swelling pressure,  $\mathbf{I}$  is the identity tensor,  $\phi_F$  is the volume fraction of the fibrillar network,  $\boldsymbol{\sigma}_M$  is the true matrix stress,  $\boldsymbol{\sigma}_F$  is the true fibre stress. From equation (1), we get

$$\Delta\Pi = \frac{1}{3} \text{trace} \left( (1 - \phi_F) \boldsymbol{\sigma}_M + \phi_F \boldsymbol{\sigma}_F \right) \quad (2)$$

The constitutive model of the matrix is assumed isotropic Neo-Hookean, while the one for the fibrils is anisotropic Neo-Hookean. The strain energy density function  $W_M$  for the Neo-Hookean model is (Holzapfel, 2001)

$$W_M = C_1 (I_1 - \log I_3 - 3) + D_2 (\log I_3)^2 \quad (3)$$

where  $C_1 = \mu/2$  with  $\mu$  being the shear modulus,  $D_2 = K/2$  with  $K$  being the bulk modulus,  $I_1 = \text{trace}(\mathbf{C})$  and  $I_3$  are respectively the first and the third invariant of the *Right Cauchy-Green* strain tensor  $\mathbf{C}$

$$\mathbf{C} = \mathbf{F}^T \mathbf{F} \quad (4)$$

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where  $\mathbf{F}$  is the *deformation gradient*, and  $I_3 = J^2$ , where  $J = \det \mathbf{F}$ .

The strain energy density function  $W_F$  for the anisotropic Neo-Hookean model is (Ateshian et al., 2009; Nagel and Kelly, 2013)

$$W_F = C_4 (I_4 - 1)^\beta \quad (5)$$

where  $C_4$  and  $\beta$  are material parameters and  $I_4$  is the invariant

$$I_4 = \mathbf{a}_0^T \mathbf{C} \mathbf{a}_0 \quad (6)$$

where  $\mathbf{a}_0$  indicates the direction of the fibrils (Holzapfel, 2001), with Cartesian components

$$\mathbf{a}_0 = \begin{bmatrix} \sin \varphi \cos \theta \\ \sin \varphi \sin \theta \\ \cos \varphi \end{bmatrix} \quad (7)$$

with  $\varphi$  being the polar angle and  $\theta$  the azimuthal angle. The value of  $\beta$  follows the restriction  $\beta > 2$ , which allows a smooth transition for the fibril response from compression to tension (Ateshian et al., 2009).

The *Second Piola-Kirchhoff* stresses are given by

$$\mathbf{S}_M = 2 \frac{\partial W_M}{\partial \mathbf{C}} \quad \mathbf{S}_F = \int_{-\pi}^{\pi} d\varphi \int_0^{2\pi} 2 \frac{\partial W_F}{\partial \mathbf{C}} \sin \varphi d\theta \quad (8)$$

where the anisotropic fibrillar stress has been averaged over all the possible directions. The *Second Piola-Kirchhoff* stress is related to the Cauchy stress through

$$\boldsymbol{\sigma} = \frac{1}{J} \mathbf{F} \mathbf{S} \mathbf{F}^T \quad (9)$$

Let us assume an isotropic deformation

$$\mathbf{F} = \lambda \mathbf{I} \quad (10)$$

with  $\lambda$  being the stretch ratio,  $\lambda = 1 + \epsilon$  with  $\epsilon$  being the strain. Then, substituting equation (9) into equation (2) leads to (Nagel and Kelly, 2013)

$$\Delta\Pi = (1 - \phi_F) \frac{2}{\lambda^3} \left[ C_1 (\lambda^2 - 1) + 12 D_2 \log \lambda \right] + \phi_F \frac{8}{3} \pi C_4 \beta \frac{(\lambda^2 - 1)^{\beta-1}}{\lambda} \quad (11)$$

## References

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