1 Supplementary information

3	Thermal structure of the Venusian atmosphere from the sub-
4	cloud region to the mesosphere as observed by radio
5	occultation
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17 1. Detail of the comparison of the temperature distributions obtained by the radio

18 occultation measurements, VIRA and VIRA-2

19 Fig. S1 shows temperature differences among the radio occultation measurements,

20 VIRA and VIRA-2. Fig. S2 shows the temperature distributions obtained from the

21 Akatsuki and Venus Express radio occultation measurements. Although the temperature

- 22 obtained on the Akatsuki mission is 2–5 K lower than that obtained on the Venus
- 23 Express mission at all altitudes (see also Fig. S1), the shape and curvature of these
- 24 temperature profiles are similar. The static stabilities obtained from these radio
- 25 occultation measurements are then almost consistent as shown in Fig. S3.
- 26



28 Figure S1 | Differences in temperature among radio occultation measurements,







36 Figure S2 | Comparison of mean temperatures at altitudes of 40–85 km for



38 Akatsuki and (red) Venus Express missions. The number within parentheses in each

39 panel is the number of radio occultation measurements made at an altitude of 50 km. An



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44 latitudes of 0°-30° obtained from radio occultation measurements made on (black)

45 Akatsuki and (red) Venus Express missions. The number within parentheses in each

46 panel is the number of radio occultation data measurements made at an altitude of 50

47 km. An error bar represents the standard deviation of the temperature.



49 **2. Measurement error due to horizontal drift**

50 Venus Express had a polar orbit. Horizontal drift of the ray path tangent point thus 51 generated measurement error in the temperature at low latitudes. The error is estimated 52 here following the work of Kursinski et al. (1997). 53 The bending angle error, $\delta \alpha(z, z_0, y)$, due to horizontal drift of the ray path tangent 54 point is expressed to the first order as $\delta \alpha(z, z_0, y) = \Delta y(z, z_0, y) \frac{\partial \alpha}{\partial y}(z, y) ,$ 55 (A1) where y represents latitude and z is the height of the error, z_0 is the altitude of the 56 57 lower integration limit of the Abel transform and the altitude of the refractivity retrieval, 58 and Δy is the horizontal distance that the ray path tangent point drifts between z and z_0 . For simplicity, we assume that α is proportional to the vertical density gradient 59 such that $\alpha \sim c\rho/H$, where c is a scale factor, ρ is the mass density and H = RT/g60 61 is the density scale height, with R, T and g respectively being the gas constant, 62 temperature and gravity acceleration. Making this assumption, it follows that $\frac{1}{\alpha}\frac{\partial\alpha}{\partial y}(z,y) \sim \frac{1}{\rho}\frac{\partial\rho}{\partial y}(z,y) - \frac{1}{T}\frac{\partial T}{\partial y}(z,y) \; .$ 63 (A2)

64 In the case of Venus, the cyclostrophic balance can be applied as

65
$$\frac{u_c^2 \tan\varphi}{a} = -\frac{1}{\rho} \frac{\partial p}{\partial y} , \qquad (A3)$$

66 where a is the radius of Venus, u_c is the cyclostrophic wind, p is pressure and φ is the

67 latitude. Taking the vertical gradient of (A3) gives

68
$$\frac{\partial}{\partial z} \left(\frac{\rho u_c^2}{a} \right) = -\frac{1}{\tan\varphi} \frac{\partial}{\partial y} \left(\frac{\partial p}{\partial z} \right) = \frac{g}{\tan\varphi} \frac{\partial \rho}{\partial y} .$$
(A4)

69 The first term on the right-hand side of (A2) can therefore be rewritten as

70
$$\frac{1}{\rho}\frac{\partial\rho}{\partial y}(z,y) = \frac{\tan\varphi}{ag}\left(2u_c\frac{\partial u_c}{\partial z} + \frac{u_c^2}{\rho}\frac{\partial\rho}{\partial z}\right) = \frac{\tan\varphi}{ag}\left(2u_c\frac{\partial u_c}{\partial z} - \frac{u_c^2}{H}\right).$$
(A5)

71 The second term on the right-hand side of (A2) is related to the vertical gradient of the

72 cyclostrophic wind such that

73
$$\frac{\partial}{\partial z} \left(\frac{u_c^2 \tan \varphi}{a} \right) = -\frac{R}{H} \frac{\partial T}{\partial y} = -\frac{g}{T} \frac{\partial T}{\partial y} .$$
(A6)

74 Accordingly, (A2) is finally described as

75
$$\frac{1}{\alpha}\frac{\partial\alpha}{\partial y}(z,y) = \frac{tan\varphi}{ag}\left(4u_c\frac{\partial u_c}{\partial z} - \frac{u_c^2}{H}\right).$$
(A7)

76 The fractional error in the bending angle due to horizontal drift of the ray path tangent

77 point is then written as

78
$$\frac{\delta \alpha}{\alpha} = \frac{tan\varphi}{ag} \left(4u_c \frac{\partial u_c}{\partial z} - \frac{u_c^2}{H} \right) \Delta y \quad (A8)$$

The radius (a) is 6050 km, gravitational acceleration (g) is 8.87 m s⁻², the density scale

80 height (*H*) is fixed at 5 km and the latitude (φ) is fixed at 20°. We consider the two

81	altitude ranges of 40–50 km and 60–70 km. In the case of the Venus Express radio
82	occultation measurements made at low latitude, Δy is approximately 10 ³ km at
83	altitudes of 40–50 km and 10^2 km at altitudes of 60–70 km, which corresponds to
84	latitudinal movement of ~10° and ~1°, respectively. Inserting these values into (A8) to
85	calculate the temperature measurement error (δT) yields $\delta \alpha / \alpha \sim 6.7 \times 10^{-4}$ and $\delta T \sim$
86	0.05 K at altitudes of 40–50 km and $\delta \alpha / \alpha \sim 2.7 \times 10^{-4}$ and $\delta T \sim 0.03$ K at altitudes
87	of 60–70 km. These δT values are smaller than the measurement error of 0.1 K, which
88	is associated with the stability of the oscillator onboard the spacecraft.
89	
90	References
91	Kursinski, E. R., Hajj, G. A., Schofield, J. T., Linfield, R. P. & Hardy, K. R. Observing
92	Earth's atmosphere with radio occultation measurements using the Global Positioning
93	System. J. Geophys Res. 102, 23429–23466 (1997).
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95	Figure legends
96	Figure S1 Differences in temperature among radio occultation measurements,

97	VIRA and VIRA-2. At latitudes of 0°–30°, the temperature differences between Venus
98	Express and Akatsuki (black), between Venus Express and VIRA (red), and between
99	Venus Express and VIRA-2 (blue) are shown. At other latitudes, the temperature
100	differences between the present study and VIRA (red) and between VIRA-2 (blue) and
101	Magellan radio occultation measurements (green) are shown.
102	
103	Figure S2 Comparison of mean temperatures at altitudes of 40–85 km for
104	latitudes of 0°–30° obtained from radio occultation measurements made on (black)
105	Akatsuki and (red) Venus Express missions. The number within parentheses in each
106	panel is the number of radio occultation measurements made at an altitude of 50 km. An
106 107	panel is the number of radio occultation measurements made at an altitude of 50 km. An error bar represents the standard deviation of the temperature.
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106 107 108 109	panel is the number of radio occultation measurements made at an altitude of 50 km. An error bar represents the standard deviation of the temperature. Figure S3 Comparison of mean static stabilities at altitudes of 40–85 km for
106 107 108 109 110	 panel is the number of radio occultation measurements made at an altitude of 50 km. An error bar represents the standard deviation of the temperature. Figure S3 Comparison of mean static stabilities at altitudes of 40–85 km for latitudes of 0°–30° obtained from radio occultation measurements made on (black)
106 107 108 109 110 111	 panel is the number of radio occultation measurements made at an altitude of 50 km. An error bar represents the standard deviation of the temperature. Figure S3 Comparison of mean static stabilities at altitudes of 40–85 km for latitudes of 0°–30° obtained from radio occultation measurements made on (black) Akatsuki and (red) Venus Express missions. The number within parentheses in each

113 km. An error bar represents the standard deviation of the temperature.