# **Supplementary Information**

**Lattice distortion induced internal electric field in TiO<sup>2</sup> photoelectrode for efficient charge separation and transfer**

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# **Supplementary Figures**



**Supplementary Figure 1 |** Electron distribution and density of states. (a, b) top view and side view of the electron density distribution of the pristine  $TiO<sub>2</sub>$  system. The blue and red balls represent Ti and O atoms, respectively. Yellow surface represents the electron distribution. (c, d) Density of states (DOS) (DOS on specified orbitals) of  $TiO<sub>2</sub>$  and  $Li<sub>0.5</sub>TiO<sub>2</sub>$  at the DFT level. The Fermi level is at zero energy and represented by dashed lines.



**Supplementary Figure 2 | SEM images of pure TiO<sub>2</sub>. (a) Top-view and (b) cross** section-view SEM images of pure TiO<sub>2</sub> tubes supported on Ti foil.



**Supplementary Figure 3** | Discharge curve of Li-TiO<sub>2</sub> nanotube array film based battery under the current of 0.03 mA (around 0.08 C).



**Supplementary Figure 4** | X-ray photoelectron spectroscopy spectra of samples. (a) Ti 2p and (b) O 1s for the pristine  $TiO<sub>2</sub>$  (black line) and Li-TiO<sub>2</sub> (red line). In addition to two dominant peaks at 465.1 eV (Ti 2p1/2) and 459.3 eV (Ti 2p3/2) associated with  $Ti<sup>4+</sup>$  in Li-TiO<sub>2</sub>, two additional peaks at 462.1 eV (Ti 2p1/2 level) and 456.6 eV (Ti 2p3/2 level) can be assigned to  $Ti^{3+}$ . In O 1s spectrum of Li-TiO<sub>2</sub>, the strong peak at 532.4 eV can be assigned to oxygen vacancies, which is typically observed for reduced  $TiO<sub>2</sub>$ .<sup>1</sup>



**Supplementary Figure 5** | The Williamson-Hall (W-H) relationship for the pristine and 5 h-lithiation treated  $TiO<sub>2</sub>$  samples. It is derived from the full-width Half Maximum (FWHM,  $\beta$ ) and the diffraction peaks (2 $\theta$ ) of TiO<sub>2</sub>. A least squares (R<sup>2</sup>) method is used to fitting the original data (dots) with a linear relationship (solid line)



**Supplementary Figure 6** | HR-TEM image and FFT analysis of TiO<sub>2</sub> nanotubes. The selected area with accurate focus in (a) is processed with fast Fourier transfer (FFT) as indicated in (b). According to the angle and distance between the diffraction points, the facets in the FFT image can be indexed. The lattice distance is calculated based on the distance in FFT images (c~f), where d1, d2, d3, and d4 equal to  $0.471$  nm (2/4.25), 0.350 nm (2/5.72), 0.350 nm (2/5.72), and 0.188 nm (2/10.64). In the meanwhile,  $d_{004}$ ,  $d_{101}$ , and  $d_{200}$  was 0.238 nm, 0.352 nm, and 0.189 nm in anatase TiO<sub>2</sub> (PDF#21-1272). Thus, it is concluded that the facets in  $(1)$ ,  $(2)$  and  $(4)$  belong to  $\{002\}$ ,  $\{101\}$  and  ${200}$  as shown in (b).



**Supplementary Figure 7** | HR-TEM image and FFT of TiO<sub>2</sub>. According to these TEM images (a, b, c, g, h, i) and the corresponded FFT images (d, e, f, j, k, l), the lattice distance of  $d_{101}$ ,  $d_{200}$ , and  $d_{004}$  is 0.348±0.0007 nm, 0.187±0.0006 nm, and 0.249±0.0199 nm, respectively.



**Supplementary Figure 8** | HR-TEM image and FFT analysis of Li-TiO<sub>2</sub>. According to these TEM images (a, b, e, f) and the corresponded FFT images (c, d, g, h),, the lattice distance of  $d_{101}$ ,  $d_{200}$ , and  $d_{004}$  is 0.348±0.0003 nm, 0.189±0.0012 nm, and  $0.231 \pm 0.0016$  nm, respectively.



**Supplementary Figure 9 | TEM images of Li-TiO<sub>2</sub> tubes. (a) TEM image of TiO<sub>2</sub>** tube photoelectrode cut by focused ion beam (FIB). (b) HR-TEM image of the wall of two adjacent tubes. (c) and (d) The HRTEM images of a single  $TiO<sub>2</sub>$  tube at two adjacent zone with a total length of *ca.* 300 nm.



**Supplementary Figure 10** | *I-V* characterization of the electrode. (a) cyclic voltammetry (CV) scan of the devices fabricated from (black line)  $TiO<sub>2</sub>$  and (red line) Li-TiO<sub>2</sub>. (b) zoom-in image of (a) to indicate the intersection voltage  $(V_i)$  at I = 0. (c) schematic of the device used for CV measurement. (d) theoretical equivalent circuit of (c). (e) ideal *I-V* response during CV scan with (red line) and without (blue line) a constant bias  $(E_0)$  on the resistor  $(R)$  according to  $(d)$ . (f) The plots of log *I* as a function of V. The intersection voltage  $(V_i)$  in (e) is highlighted as sharp spikes in (f).



**Supplementary Figure 11** | The incident photon-to-current efficiency (IPCE) of TiO<sub>2</sub> and  $Li-TiO<sub>2</sub>$  photoelectrodes. The  $Li-TiO<sub>2</sub>$  photoanode shows a threshold at the wavelength of around 420 nm, with is in consistent with the bandgap of  $TiO<sub>2</sub>$ . The increased light absorption beyond 420 nm caused by the lithiation-induced reduction effect makes negligible contribution to the photocurrent response.



**Supplementary Figure 12** | Electrochemical impedance spectroscopy (EIS) of (black line)  $TiO<sub>2</sub>$  and (red line) Li-TiO<sub>2</sub> and the equivalent circuit used for fitting.



**Supplementary Figure 13** | The photocurrent stability tests in 1 M NaOH solution ( $pH = 13.6$ ) and 0.5 M Na<sub>2</sub>SO<sub>4</sub> ( $pH = 7$ ) at 1.0 V *vs* SHE. The photocurrent keeps a good stability in NaOH solution for a whole test but experiences a rapid decrease at the first 5 minutes and then gradual decrease in  $Na<sub>2</sub>SO<sub>4</sub>$  solution.



**Supplementary Figure 14** | Mott-Schottky (M-S) curves of samples. (a, b) M-S curves of (black line)  $TiO<sub>2</sub>$ , (red line) Li-TiO<sub>2</sub> and (orange line) reduced  $TiO<sub>2</sub>$  (r-TiO<sub>2</sub>) TiO<sup>2</sup> nanotube photoanodes.



**Supplementary Figure 15** | The photocurrent of (black line) pristine  $TiO<sub>2</sub>$ , (orange line) reduced  $TiO<sub>2</sub>$ , and (red line) Li-TiO<sub>2</sub>. For the r-TiO<sub>2</sub> electrode, electrochemical reduction at -1.6 V vs Ag/AgCl was carried out deliberately. The amount of charge used to reduce  $TiO<sub>2</sub>$  is the same (0.08 C) as that passed through Li-TiO<sub>2</sub> during the discharging. In this way, the amount of reduced Ti species is anticipated to be close in  $r$ -TiO<sub>2</sub> and Li-TiO<sub>2</sub> electrodes, as indicated by the Mott-Schottky curves in Fig. S10.



**Supplementary Figure 16** | Performance based on TiO<sub>2</sub>|FTO photoelectrodes. (a, b) Optical images of pristine  $TiO<sub>2</sub>$  nanorod electrode, lithiated  $TiO<sub>2</sub>$  nanorod photoanode  $(Li-TiO_2|FTO)$  and annealed Li-TiO<sub>2</sub>|FTO (a-Li-TiO<sub>2</sub>|FTO). The pristine TiO<sub>2</sub>|FTO and a-Li-TiO<sub>2</sub>|FTO have nearly white color, and the Li-TiO<sub>2</sub>|FTO has blue color. (c) Mott-Schottky curves of the  $TiO_2|FTO$ , Li-TiO<sub>2</sub>|FTO, and a-Li-TiO<sub>2</sub>|FTO photoanodes and (d) photocurrent response curves of three photoelectrodes under 1 sun simulated light. Based on the slope of M-S curves, the charge carrier concentration of Li-TiO<sub>2</sub> was determined to be much higher than that of TiO<sub>2</sub> due to the reductive effect of lithiation. The carrier concentration of a-Li-TiO<sub>2</sub> is comparable to that of TiO<sub>2</sub>. The Li-ions accommodated in the a-Li-TiO<sub>2</sub> can cause lattice distortion. Therefore, comparison of the photocurrent density of a-Li-TiO<sub>2</sub>|FTO,  $Li-TiO<sub>2</sub>|FTO$  and  $TiO<sub>2</sub>|FTO$  in (d) clearly indicates the substantial role of lattice distortion in enhancing photocurrent.



**Supplementary Figure 17** | Characterization of Li doped TiO<sub>2</sub> reference sample. (a)  $XPS$  spectrum of Li 1s in Li doped  $TiO<sub>2</sub>$  reference sample and (b) photocurrent-bias curves of undoped and Li doped TiO<sub>2</sub> photoanodes.

## **Supplementary Tables**



**Supplementary Table 1 | The charge distribution around each atom in**  $Li_{0.5}TiO_2$ lattice cell.

\*It is the average valence electron number around each atom in the lattice cell schemed in Fig. 2b. According to the results, there are 0.145 electrons around Li atoms on average, suggesting that each Li atoms will donate 0.845 electrons.

	Slope	intercept	
Pristine $TiO2$	0.00217	0.00110	0.667
5 h lithiated $TiO2$	0.00498	0.00014	0.664
$(Li-TiO2)$			

**Supplementary Table 2** | The W-H relationship for pristine  $TiO<sub>2</sub>$  and the Li-TiO<sub>2</sub>

## **Supplementary Notes**

**Supplementary Note 1.** In Supplementary Figure 5, according to the Williamson-Hall relationship, the broadening of peaks can be expressed as Supplementary Equation 1:<sup>2</sup>

$$
\beta \times \cos \theta = \frac{k \times \lambda}{D} + 4\varepsilon \times \sin \theta \tag{1}
$$

Where k is a constant,  $\lambda$  is the irradiation wavelength, D is the particle diameter and  $\epsilon$ is the lattice strain. Accordingly, the  $\beta \times \cos \theta \sim 4 \times \sin \theta$  should give a linear relationship with the slope of strain  $(\varepsilon)$ . Based on the diffraction peaks of (101), (004), (200), (211), (105) and (204) from the synchrotron XRD, the W-H relationship for pristine  $TiO<sub>2</sub>$ and the one after 5 h lithiation can be fitted within Table S1.

**Supplementary Note 2.** In Supplementary Figure 10, an indium particle is placed on the top of photoelectrode to enhance the contact between the tip and semiconductor film, and to avoid the short-circuiting. Soft indium particle can easily be pasted on the electrode top so it can work as an intermediate layer to enhance the contact between the tip and nanotube film. Moreover, the solid-state nature of indium prevents its penetration through the tubes to reach the bottom Ti foil, causing short-circuit.

The device in (c) is similar to a physical plane capacitor with a metal|dielectric|metal  $(Ti|TiO<sub>2</sub>|In)$  structure. So a parallel capacitor with a resistor in (d) is used to simulate the equivalent circuit. In this case, we can predict the *I-V* relationship at the scan rate of ±*v* during the CV scan as Supplementary Equation 2:

$$
I=(V-E_0)/R \pm \nu C \tag{2}
$$

In Supplementary Equation 3,  $(V-E_0)/R$  is the current passed through the resistor, which is independent of the scan rate. The slope of *I-V* curves is determined by the resistance  $(R)$  of the film. The  $\pm \nu C$  is the current passed through the capacitor, which is determined by the scan rate  $(v)$  and capacitance  $(C)$ , but independent of the applied voltage (*V*).  $E_0$  is voltage created by dipolar in the bulk of TiO<sub>2</sub>.

According, the intersected voltage  $(V_i)$  can be expressed as:

$$
V_i = E_0 \pm \nu C/R \tag{3}
$$

For pure TiO<sub>2</sub>,  $E_0$  is null, therefore,  $V_i$  should be  $\pm vC/R$ , which is a mirror symmetry at *V*=0. This result is constant with the result in (b) and Figure 5b ( $V_i = \pm 0.03$  V). For Li-TiO<sub>2</sub>, the symmetry axis shows a positive shift ( $V_i = -0.07, 0.37$  V), suggesting that there is a positive  $E_0$  generated in the device. Considering the configuration of the device and measurement, it is concluded that there should be an electric field directed towards the surface from the bottom as indicated in (c).

**Supplementary Note 3.** In Supplementary Figure 14, according to the *M-S*

Supplementary Equations 4 and  $5$ :

$$
C^{-2} = \frac{2}{q \varepsilon N_d} \left( E - E_{fb} - \frac{k_B T}{q} \right)
$$
  

$$
N_d = \frac{2}{q \varepsilon} \left( dC^{-2} / dE \right)^{-1}
$$
 (5)

Where *q* is the elemental charge,  $\varepsilon$  is the dielectric constant,  $k_B$  is the Boltzman constant,  $E_{\text{fb}}$  is the flat band potential, *T* the temperature in kelvin,  $N_{\text{d}}$  is the carrier concentration.

The carrier concentration can be evaluated from the slope of M-S curves  $\left( \frac{d^{2}C^{-2}}{dE} \right)$ .

Upon lithiation, the  $N_d$  is significantly improved due to the reductive effect. For the electrochemically reduced  $r-TiO_2$ , the amount of charge passing the electrode is the same as that passing  $Li-TiO<sub>2</sub>$ . Thus, r-TiO<sub>2</sub> and  $Li-TiO<sub>2</sub>$  were reduced at a similar extent and have comparable concentrations of charge carriers as indicated by the slope of M-S curves in Supplementary Figure 14 (b).

### **Supplementary Methods**

**Synthesis.** Rutile TiO<sub>2</sub> photoelectrodes grown on FTO glass substrates are prepared by a hydrothermal method.<sup>[4](#page-24-0)</sup> Typically, 1.0 mL of titanium butoxide (TBT, Sigma-Aldrich) was dissolved in a mixed solution containing 20 mL of hydrochloric acid (Sigma-Aldrich) and 30 mL of distilled water. A piece of clean FTO glass substrate was placed into the solution vertically in an autoclave with the Polytetrafluoroethylene (PTFE) line. The hydrothermal reaction was carried out at 180  $^{\circ}$ C for 2 h. After cooling down to room temperature, the TiO<sub>2</sub>|FTO electrode was cleaned and further annealed at 450  $^{\circ}$ C (the ramping rate of 5  $^{\circ}$ C min<sup>-1</sup>) for 2 h. TiO<sub>2</sub> in the electrode has the nanorod morphology. The lithiation of TiO<sub>2</sub>|FTO photoelectrode was conducted following the same procedures as for the  $TiO<sub>2</sub>$  nanotube photoelectrode. The resultant electrode is denoted as  $Li$ -TiO<sub>2</sub>|FTO photoelectrode. In order to quench oxygen vacancies caused by the lithiation, the  $Li-TiO<sub>2</sub>|FTO$  photoelectrode was then annealed in air at 450 °C (the ramping rate of 1 °C min<sup>-1</sup>) for another 2 h and the sample obtained is denoted as a-Li-TiO<sub>2</sub>|FTO photoelectrode.

The Li doped TiO<sub>2</sub>|FTO photoelectrodes were prepared by spin coating.<sup>5-7</sup> 0.2 M TBT ethanol solution was used as precursor and LiAc was dissolved in it to make the Li/Ti molar ratio of 0.45 (the same ratio in Ti foil supported Li-TiO<sub>2</sub> photoelectrode). The film was fabricated at a spin rate of 1000 rpm for 1 min with 0.1 mL of precursor solution dispersed on an FTO glass. After spin coating, the film was heated at 120  $^{\circ}$ C for 5 minutes on a hot plate. The procedure was repeated 10 times to increase the film thickness. In the end, the film was annealed at 450 <sup>o</sup>C for 3 h in a muffle oven to achieve the anatase phase.

**The incident photon-to-current efficiency (IPCE) test** The IPCE was obtained using an Oriel Cornerstone 260 1/4 m monochromators with a 300 W Oriel Xenon lamp as the simulated light source. The light power was calibrated with a standard Si diode. The power density at a specific wavelength was measured by a Newport 1918-c power meter. The IPCE can be calculated according to following Supplementary Equation 6. 3

$$
I PCE = \frac{j (mA cm^2 - 2) \times 1239.8 (V.nm)}{\lambda (nm) \times I (mW cm^2 - 2)} \times 100\%
$$
 (6)

*j* is the photocurrent density (mA cm<sup>-2</sup>) measured from the electrochemical workstation, *λ* refers to the incident light wavelength (nm), and *I* is the light density measured at a specific wavelength  $(mW cm<sup>-2</sup>)$ .

**DFT calculation of the formation energy** The formation energies of  $Li<sub>n</sub>TiO<sub>2</sub>$  were defined as Supplementary Equation 7:

$$
\Delta E = (E_{\text{Li}_n \text{TiO}_2} - E_{\text{TiO}_2} - nE_{\text{Li}}) / (n+1). \tag{7}
$$

where  $E_{\text{Li}_n \text{TiO}_2}$ ,  $E_{\text{TiO}_2}$ , and  $E_{\text{Li}}$  are the total energies of  $\text{Li}_0$ <sub>5</sub>TiO<sub>2</sub>, TiO<sub>2</sub> and a Li atom in the structure, respectively. Based on the calculation,  $E_{\text{Li}_2\text{Ti}_4\text{O}_8}$ ,  $E_{\text{Ti}_4\text{O}_8}$ , and  $E_{\text{Li}}$  were -98.078eV, -91.378eV, and -1.344 eV.

$$
\Delta E = (E_{\text{Li}_2\text{Ti}_4\text{O}_8} - E_{\text{Ti}_4\text{O}_8} - 2E_{\text{Li}}) / (2+1) = 1.34 \text{eV} \quad (8)
$$

Based on Supplementary Equation 8, *ΔE* is 1.34 eV.

### **Supplementary References**

- 1. Yang Y.; Yin L.; Gong Y.; Niu P.; Wang J.; Gu L.; Chen X.; Liu G.; Wang L.; Cheng H., An Unusaual Strong Visible-Light Absorption Band in Red Antase TiO<sub>2</sub> Photocatalyst Induced by Atomic Hydrogen-Occupied Oxygen Vacancies, *Adv. Mater.*, **2018,** 30(6), 1704479.
- 2. Williamson, G.; Hall, W., X-ray line broadening from filed aluminium and wolfram. *Acta Metal.*  **1953,** *1* (1), 22-31.
- 3. Wang, Z.; Wang, L., Progress in designing effective photoelectrodes for solar water splitting. *Chinese J. Catal.* **2018,** *39* (3), 369-378.
- 4. Yang, Y.; Wang, S.; Jiao, Y.; Wang, Z.; Xiao, M.; Du, A.; Li, Y.; Wang, J.; Wang, L., An unusual red carbon nitride to boost the photoelectrochemical performance of wide bandgap photoanodes. *Adv. Funct. Mater.* **2018,** *28* (47), 1805698.
- 5. Wang, C.; Chen, Z.; Jin, H.; Cao, C.; Li, J.; Mi, Z., Enhancing visible-light photoelectrochemical water splitting through transition-metal doped TiO<sub>2</sub> nanorod arrays. *J. Mater. Chem. A* **2014,** 2 (42), 17820-17827.
- <span id="page-24-0"></span>6. Feng, X.; Shankar, K.; Paulose, M.; Grimes, C. A., Tantalum‐Doped Titanium Dioxide Nanowire Arrays for Dye‐Sensitized Solar Cells with High Open‐Circuit Voltage. *Angew. Chem. Int. Ed.*  **2009,** *48* (43), 8095-8098.
- 7. Pan, L.; Zou, J.-J.; Zhang, X.; Wang, L., Photoisomerization of norbornadiene to quadricyclane using transition metal doped TiO<sub>2</sub>. Ind. Eng. Chem. Res. 2010, 49 (18), 8526-8531.