## **Supporting Information**

## **Phonon Anharmonicity in Bulk** *Td***-MoTe2**

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**Figure S**1 a) Temperature-dependent Raman spectra for Peak 1 normalized by the intensity of Peak 2 as labelled in Figure 1c. Solid black lines represent single Lorentzian fits with a linear background. b) FWHM of Peak 1 extracted from the fits at each temperature. c) Amplitudes of Peak 1, extracted from the fits in (a). 1σ error bars are included.

## **Raman Mode Assignment**

Raman experiments in a backscattering geometry are sensitive to modes with *A<sup>g</sup>* and *B<sup>g</sup>* symmetry for the *1T*' phase and with  $A_I$  and  $B_I$  symmetries in the  $T_d$  phase.<sup>1</sup> In polarization-dependent Raman

measurements of *1T'*-MoTe2, which will be reported elsewhere, we have determined that Peaks 1- 4 are of  $A_g$  symmetry. We now argue that these  $A_g$  modes evolve into  $A_1$  modes in  $T_d$ -MoTe<sub>2</sub> based upon the similarity of their Raman tensors. The tensors for  $A<sub>q</sub>$  and  $A<sub>1</sub>$  symmetry Raman modes are $<sup>1</sup>$ </sup>

$$
\overrightarrow{R}_{A_g} = \begin{pmatrix} a & d & 0 \\ d & b & 0 \\ 0 & 0 & c \end{pmatrix} \qquad \overrightarrow{R}_{A_1} = \begin{pmatrix} a & 0 & 0 \\ 0 & b & 0 \\ 0 & 0 & c \end{pmatrix}.
$$
 (S1)

The Raman cross-section  $S$  can be calculated from<sup>2</sup>

$$
S = |\hat{e}_s \cdot \vec{R} \cdot \hat{e}_t|^2 \tag{S2}
$$

where  $\hat{e}_i$  and  $\hat{e}_s$  are unit vectors for the polarization of the incident and scattered photons. In our experiment the excitation and collection paths are collinear and perpendicular to the basal plane of MoTe<sub>2</sub> (a backscattering geometry). The laser is polarized at an arbitrary angle  $\theta$  with respect to the  $\alpha$  crystal axis and  $\phi$  is the polarization angle of the scattered light with respect to the same axis.  $\phi = \theta$  and  $\phi = \theta + \frac{\pi}{2}$  correspond to parallel and perpendicular collection configuration respectively. Therefore,

$$
\hat{e}_t = (\sin \theta \quad 0 \quad \cos \theta), \tag{S3}
$$

$$
\widehat{e_s} = (\sin \phi \quad 0 \quad \cos \phi). \tag{S4}
$$

We now evaluate Eq. S2 for  $A<sub>g</sub>$  symmetry modes using Eq. S3 and Eq. S4

$$
S_{A_g} = \begin{vmatrix} (\sin \phi & 0 & \cos \phi \end{vmatrix} \begin{pmatrix} a & d & 0 \\ d & b & 0 \\ 0 & 0 & c \end{pmatrix} \begin{pmatrix} \sin \theta \\ 0 \\ \cos \theta \end{pmatrix} \Big|^2 \tag{S5a}
$$

$$
= \begin{vmatrix} (\sin \phi & 0 & \cos \phi) \begin{pmatrix} a \sin \theta \\ d \sin \theta \\ c \cos \theta \end{pmatrix} \end{vmatrix}^2
$$
 (S5b)

$$
= |a\sin\theta\sin\phi + c\cos\theta\cos\phi|^2. \tag{S5c}
$$

We can perform a similar calculation for the  $A_1$  symmetry mode

$$
S_{A_1} = \begin{vmatrix} (\sin \phi & 0 & \cos \phi) \begin{pmatrix} a & 0 & 0 \\ 0 & b & 0 \\ 0 & 0 & c \end{pmatrix} \begin{pmatrix} \sin \theta \\ 0 \\ \cos \theta \end{pmatrix} \end{vmatrix}^2
$$
 (S6a)

$$
= \begin{vmatrix} (\sin \phi & 0 & \cos \phi) \begin{pmatrix} a \sin \theta \\ 0 \\ c \cos \theta \end{pmatrix} \end{vmatrix}^2 \tag{S6b}
$$

$$
= |a\sin\theta\sin\phi + c\cos\theta\cos\phi|^2. \tag{S6c}
$$

Equations S5c and S6c are identical which, when combined with the continuous evolution of each mode with temperature, suggests that  $A<sub>g</sub>$  modes will evolve into  $A<sub>1</sub>$  modes at low temperature.

A similar analysis can be performed for the  $B<sub>g</sub>$  and  $B<sub>1</sub>$  symmetry Raman tensors

$$
\overrightarrow{R}_{B_g} = \begin{pmatrix} 0 & 0 & f \\ 0 & 0 & g \\ f & g & 0 \end{pmatrix} \qquad \overrightarrow{R}_{B_1} = \begin{pmatrix} 0 & 0 & f \\ 0 & 0 & 0 \\ f & 0 & 0 \end{pmatrix}.
$$
 (S7)

Following the same approach, we find

$$
S_{B_g} = \begin{vmatrix} (\sin \phi & 0 & \cos \phi \end{vmatrix} \begin{pmatrix} 0 & 0 & f \\ 0 & 0 & g \\ f & g & 0 \end{pmatrix} \begin{pmatrix} \sin \theta \\ 0 \\ \cos \theta \end{pmatrix}^2 = |f|^2 (\sin(\theta + \phi))^2, \tag{S8}
$$

$$
S_{B_1} = \begin{vmatrix} (\sin \phi & 0 & \cos \phi \end{vmatrix} \begin{pmatrix} 0 & 0 & f \\ 0 & 0 & 0 \\ f & 0 & 0 \end{pmatrix} \begin{pmatrix} \sin \theta \\ 0 \\ \cos \theta \end{pmatrix}^2 = |f|^2 (\sin(\theta + \phi))^2
$$
 (S9)

Equations S8 and S9 are identical which, when combined with the continuous evolution of this mode with temperature, suggests that  $B<sub>g</sub>$  modes will evolve into  $B<sub>1</sub>$  modes at low temperature.

## **Supplementary References**

<sup>1</sup> M.I. Aroyo, J.M. Perez-Mato, D. Orobengoa, E. Tasci, G. De La Flor, and A. Kirov, Bulg. Chem. Commun. **43**, 183 (2011).

<sup>2</sup> H.B. Ribeiro, M.A. Pimenta, C.J.S. De Matos, R.L. Moreira, A.S. Rodin, J.D. Zapata, E.A.T. De Souza, and A.H. Castro Neto, ACS Nano **9**, 4270 (2015).