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Supplementary Information for

## Sparseness and Smoothness Regularized Imaging for Improving the Resolution of Cryo-EM Single Particle Reconstruction

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## 41 Supplementary Information Text

### 42 Methods

#### 43 Log marginal likelihood

44 The log marginal likelihood function in cryo-EM refinement problem can be derived as follows. For  
45 simplicity, we assume the particles are from a single structure, whereas structural heterogeneity can be  
46 easily incorporated into our generative model by treating the class membership of each particle as a  
47 hidden variable(1). Since the images collected by cryo-EM are 2D projections of a 3D molecular structure,  
48 the Fourier transform of the image has the following relation with the Fourier transform of the 3D  
49 molecular structure. Let the Fourier transform of a 3D molecular structure be  $V$ , we first arrange the 3D  
50 volume  $V$  into a vector with  $L$  elements. Assume an  $N \times N$  image  $i$  is formed by rotating the 3D volume  $V$   
51 with the Euler angle set  $\phi$  and projecting along with the  $z$  axis, and shifting by  $[\Delta_x, \Delta_y]$  from the origin,  
52 using projection-slice theorem, the Fourier transform of the image  $i$  can be expressed as

$$X_{ij} = e^{i\frac{2\pi}{N}(\Delta_x h + \Delta_y k)} \text{CTF}_{ij} \sum_{l=1}^L P_{jl}^\phi V_l, \quad (\text{Equation S1})$$

53 where  $X_{ij}$  is the  $j$ th component of the Fourier transform of the image  $i$  whose corresponding 2D index is  
54  $[h, k]$ ,  $\text{CTF}_{ij}$  is the  $j$ th component of the contrast transfer function for the image  $i$ , and  $P_{jl}^\phi$  is the slice  
55 operator which cuts out the plane in the 3D Fourier transform  $V$  which is rotated from the  $xy$  plane  
56 according to the Euler angle set  $\phi$ . We elaborate on the slice operator  $P^\phi$  by giving its formal definition.  
57 Let the index of a voxel in the 3D Fourier transform  $V$  be  $[h', k', l']$ , and the index of the corresponding  
58 pixel of the Fourier transform of the image  $i$  is  $[h, k]$ , the slice operator  $P^\phi$  transforms the 3D index to 2D  
59 index by the following equation,

$$\begin{pmatrix} h \\ k \\ 0 \end{pmatrix} = R_\phi \begin{pmatrix} h' \\ k' \\ l' \end{pmatrix}, \quad (\text{Equation S2})$$

60 where  $R_\phi$  is a rotation matrix parameterized by Euler angles  $\phi$ . Furthermore, suppose the Fourier  
61 component  $X_{ij}$  is distributed according to Gaussian with the mean defined in Equation S1 and variance  
62  $\sigma^2$ , and the Gaussian noise of each component is independent, the marginal probability of observing  
63 image  $i$  can be obtained by integrating over all possible orientations  $\phi$  and translations  $\Delta = [\Delta_x, \Delta_y]$  as  
64 follows,

$$P(X_i|V) \propto \int_{\phi, \Delta} \exp \left\{ -\frac{1}{\sigma^2} \sum_{j=1}^J \left( X_{ij} - e^{i\frac{2\pi}{N}(\Delta_x h + \Delta_y k)} \text{CTF}_{ij} \sum_{l=1}^L P_{jl}^\phi V_l \right)^2 \right\} d\phi d\Delta. \quad (\text{Equation S3})$$

65 We omit translation factors in the squared difference term in Equation S3 to simplify expressions  
66 elsewhere. The marginal probability of an image can then be leveraged to construct the log marginal  
67 likelihood in Equation 1 in main text.

#### 68 Expectation maximization

69 The expectation maximization algorithm works as follows. Since the difference between log  
70 likelihoods of the marginal probability can be lower bounded by the difference between the sums of log  
71 likelihoods of the joint probability weighted by their corresponding posterior probabilities for latent  
72 variables, *i.e.*,

$$\begin{aligned} & \log P(X_i|V) - \log P(X_i|V_{k-1}) \\ & \geq \sum_{\phi} -P(\phi|X_i, V_{k-1}) \left( \|X_i - \text{CTF}_i P^\phi V\|^2 - \|X_i - \text{CTF}_i P^\phi V_{k-1}\|^2 \right), \end{aligned} \quad (\text{Equation S4})$$

73 thus maximizing the lower bound improves the log likelihood of the marginal probability at least as  
74 much(3). At the expectation step, we calculate the posterior probability of hidden variables conditioned on  
75 a given image and the map. The method to compute the posterior probability derived in RELION(1) can  
76 be applied in the context of our method without any modification.

77 **Weighted approximation**

78 The derivative of  $\sum_j \log(|x_j| + \epsilon)$  at the point  $x = x^i$  is  $\frac{\text{sign}(x_j^i)}{|x_j^i| + \epsilon}$  for each component. Hence, the

79 tangent lines of  $\sum_j \log(|x_j| + \epsilon)$  at this point are  $\sum_j \frac{\text{sign}(x_j^i)}{|x_j^i| + \epsilon} (x_j - x_j^i) + \log(|x_j^i| + \epsilon)$ , where  $\text{sign}(x_j) =$

80  $\text{sign}(x_j^i)$ . Note that  $x_j$  and  $x_j^i$  are of the same sign, the form of tangent lines can be simplified as,

$$g(x|x^i) = \sum_j \frac{|x_j|}{|x_j^i| + \epsilon} + \text{const}_j, \quad (\text{Equation S5})$$

81 where  $\text{const}_j = \frac{|x_j^i|}{|x_j^i| + \epsilon} + \log(|x_j^i| + \epsilon)$ . We hence obtain the weighted approximation for the log norm. Since

82 the log norm is concave, its tangent line is its upper bound, namely,  $g(x|x^i) \geq \sum_j \log(|x_j| + \epsilon)$ , and

83  $g(x^i|x^i) = \sum_j \log(|x_j^i| + \epsilon)$ . If  $g(x^{i+1}|x^i) < g(x^i|x^i)$ , we have  $\sum_j \log(|x_j^{i+1}| + \epsilon) < g(x^{i+1}|x^i) < g(x^i|x^i) =$

84  $\sum_j \log(|x_j| + \epsilon)$ . Combining the previous relation with Equation S4, we have  $\log P(X_i|V) - \log P(X_i|V_{k-1}) -$

85  $\sum_{j=1}^L (\alpha \log(|x_j| + \epsilon) + \beta \log(\|\nabla x_j\|_2 + \epsilon')) \geq \sum_{\phi} -P(\phi|X_i, V_{k-1}) (\|X_i - \text{CTF}_i P^{\phi} V\|^2 - \|X_i - \text{CTF}_i P^{\phi} V_{k-1}\|^2) -$

86  $g(x|x^i) - g(\nabla x|\nabla x^i)$ . Thus, we can then prove that improving the lower bound will cause the left-hand

87 side of the inequality to improve at least as much by induction. The right-hand side of the preceding

88 inequality is the Equation 4 in main text ignoring constant terms and implicit gradient restraint.

89 **Nesterov smoothed TV norm**

90 To derive the gradient of Nesterov smoothed TV norm, we begin by stating the discrete form of  
91 the gradient in TV norm. For a voxel  $x[i, j, k]$  of a 3D map  $x$ , the gradient of the 3D map obtained by  
92 discrete differentiation operator  $D$  at this voxel is of the form

$$Dx[i, j, k] = \begin{bmatrix} D_1 x[i, j, k] \\ D_2 x[i, j, k] \\ D_3 x[i, j, k] \end{bmatrix} = \begin{bmatrix} x[i, j, k] - x[i-1, j, k] \\ x[i, j, k] - x[i, j-1, k] \\ x[i, j, k] - x[i, j, k-1] \end{bmatrix}. \quad (\text{Equation S6})$$

93 By denoting  $D_i$  as the matrix representation of the discrete differentiation operator along the  $i$ th

94 dimension, we can express the gradient of the map  $x$  along the  $i$ th dimension as  $D_i x$ . By abuse of

95 notation, let  $D = [D_1, D_2, D_3]^T$  be a matrix composed by concatenating  $D_i$  by rows, the TV norm can be

96 defined as,

$$\|x\|_{\text{TV}} = \max_{u \in Q_d} \langle u, Dx \rangle, \quad (\text{Equation S7})$$

97 where  $u = [u_1, u_2, u_3]^T \in Q_d$  is the vector of dual variables of the gradients of  $x$  along three directions, and

98  $Q_d$  is the dual space in which each vector satisfies the inequality  $u_1 [i, j, k]^2 + u_2 [i, j, k]^2 + u_3 [i, j, k]^2 \leq 1$ .

99 Using Nesterov smoothing with smoothing parameter  $\mu$ , the smoothed TV norm  $f_{\mu}(x)$  is the TV norm with

100 a quadratic restraint on the dual variables, which can be written as,

$$f_{\mu}(x) = \max_{u \in Q_d} \langle u, Dx \rangle - \frac{\mu}{2} \|u\|_2^2, \quad (\text{Equation S8})$$

101 according to Candes *et al.*(4). The gradient of the smoothed TV norm  $f_{\mu}(x)$  can be expressed as

$$\nabla f_{\mu}(x) = D^T u_{\mu}(x), \quad (\text{Equation S9})$$

102 according to Candes *et al.*(4), where  $D$  is the discrete differentiation operator defined in Equation S6,

103  $u_{\mu}(x)$  is a vector of the form  $[u_1, u_2, u_3]^T$  and for each dimension  $a \in [1, 2, 3]$  and voxel  $[i, j, k]$ ,

$$u_a[i, j, k] = \begin{cases} \mu^{-1} (D_a x)[i, j, k], & \text{if } \|\nabla x[i, j, k]\| < \mu \\ (\|\nabla x[i, j, k]\|)^{-1} (D_a x)[i, j, k], & \text{otherwise.} \end{cases} \quad (\text{Equation S10})$$

104 For the weighted smoothed TV norm in Equation 4 in main text, its gradient  $u_a[i, j, k]$  can be obtained by

105 multiplying Equation S10 with the weight  $\frac{1}{\|\nabla x^i_{(i,j,k)}\| + \epsilon'}$ . Equation S10 shows that the gradient of smoothed

106 TV norm can be easily obtained by first calculating the norm of discrete gradient of the volume  $x$  at each

107 voxel  $[i, j, k]$ , and then setting the norm of gradient at this voxel  $\|\nabla x[i, j, k]\|$  with value smaller than the

108 smoothing parameter  $\mu$  to  $\mu$ , thus keeping the denominator of the gradient of TV norm in a valid range

109 and avoiding the non-differentiability of the non-smoothed TV norm at zero. With the form of discrete  
 110 differentiation operator  $D$  in Equation S6, we can write the gradient of the smoothed TV norm at a voxel  
 111  $[i, j, k]$  as follows,

$$\nabla f_{\mu}(x)_{ijk} = \sum_{a=1}^3 u_a [i, j, k] - u_a [(i, j, k) + \Delta_a], \quad (\text{Equation S11})$$

112 where  $\Delta_a$  is a  $1 \times 3$  vector with one on the  $a$ th entry and zeros elsewhere. Substituting Equation S10 into  
 113 Equation S11 leads to the complete form of the gradient in Equation 5 in main text. We then observe that  
 114 the gradient of TV norm at the voxel  $[i, j, k]$  depends on gradients of 3D map around this voxel.

### 115 Local kernel regression

116 Nonparametric regression is often used to estimate the value of a point given the values of its  
 117 neighborhoods. Denote  $Y_i$  as the value at a certain point  $x_i \in \mathbb{R}^N$ , and let  $y$  be the value at the point  $x \in$   
 118  $\mathbb{R}^N$  which is to be predicted, as in Takeda *et al.*(2), we can define  $y$  as the maximizer of

$$- \sum_{i=1}^L K(x_i, x) \|Y_i - y\|^2, \quad (\text{Equation S12})$$

119 where  $K(\cdot)$  represents a chosen kernel function. In the context of cryo-EM refinement, for an orientation  
 120  $\phi$ , let  $n_j(\phi) = [h_j, k_j, l_j] \in \mathbb{R}^3$  be the back-projected voxel in the 3D volume  $V$  which is rotated from the  
 121 Fourier coefficient  $j$  of an image  $X_i$  (see “Back-projection as lock kernel regression”), we can define the  
 122 value of the target voxel  $n = [h, k, l] \in \mathbb{Z}^3$  as the maximizer of the following regression problem,

$$- \sum_{i=1}^N \sum_{j=1}^J \sum_{\phi} P(\phi | X_i, V_{k-1}) K(n_j(\phi), n) \|X_{ij} - \text{CTF}_{ij} V_n\|^2, \quad (\text{Equation S13})$$

123 where  $K(\cdot)$  is a kernel function measuring the closeness of the back-projected voxel and the target voxel.  
 124 A summation on losses of all voxels leads to Equation 7 in main text.

### 125 Parameter settings

126 Theoretically optimal scales of the parameters in OPUS-SSRI can be obtained by using the  
 127 closed form solution of our new target function. For the LASSO type problem, the closed form solution  
 128 can be derived from its dual form(5). The first step towards the dual form of our new target function is  
 129 converting our new target function to a matrix form. Assume  $x$  is the 3D volume which is rearranged into a  
 130 vector, that is, a voxel with index  $[i, j, k]$  is mapped to the  $h$ th component  $x_h$  of  $x$ , and let  $A$  be the  
 131 corresponding 3D Fourier transform matrix, we can express the Fourier coefficients of the 3D volume  $V$   
 132 as the result of matrix vector multiplication, namely,  $V = Ax$ . With Equation 7 in main text in hand, we can  
 133 write the matrix form of our 3D reconstruction problem as

$$\min_x \frac{1}{2} \|y - DAx\|^2 + \alpha \sum_{j=1}^L \frac{|x_j|}{|x_j^i| + \epsilon} + \beta \sum_{j=1}^L \frac{\|\nabla x_j\|_2}{\|\nabla x_j^i\|_2 + \epsilon'} + \gamma \|x - x^{k-1}\|_2^2, \quad (\text{Equation S14})$$

134 where  $y$  is a vector representation of the 3D Fourier transform data with  $y(n) =$

135  $\frac{\sum_{i=1}^N \sum_{j=1}^J \sum_{\phi} P(\phi | X_i, V_{k-1}) K(n_j(\phi), n) \text{CTF}_{ij} X_{ij}}{\sqrt{N(n)}}$  for  $n = [h, k, l] \in [0, L]$ ,  $D$  is a  $L \times L$  diagonal matrix with diagonal

136 element  $D(n, n) = \sqrt{N(n)}$ , and  $Ax$  is the Fourier transform of the 3D map  $x$ . We can derive the dual form  
 137 of Equation S14 by simplifying our restraint. According to Tibshirani *et al.*(5), substituting the restraints in  
 138 Equation S14 by a generalized LASSO restraint,  $\lambda \|Gx\|_1$ , the dual of Equation S14 with new restraint is of  
 139 the form,

$$\min_u (A^H D y - G^T u)^T (A^H D^2 A)^+ (A^H D y - G^T u) \quad (\text{Equation S15})$$

140 subject to  $\|u\|_{\infty} \leq \lambda$ ,  $G^T u \in \text{row}(DA)$ , where  $(A^H D^2 A)^+$  is the Moore-Penrose inverse of  $A^H D^2 A$ ,  $\|u\|_{\infty} =$   
 141  $\max_{u_i \in u} |u_i|$ ,  $\lambda$  is the parameter for  $l_1$  restraint and  $u$  is the dual variable of the 3D volume  $x$ . Given  $u$ , the

142 closed form solution of  $x$  can be written as,

$$x = (A^H D^2 A)^+ (A^H D y - G^T u). \quad (\text{Equation S16})$$

143 In Equation S16,  $A^H D y$  represents the inverse Fourier transform of the data, which is the unregularized  
 144 solution.  $G^T u$  is the dual variable of restraint, which regularizes  $A^H D y$ . To achieve sparseness in the  
 145 solution  $x$ ,  $u$  needs to zero out certain components of  $A^H D y$ . Since the dual variable  $u$  is bounded by  $\lambda$ ,  
 146 the restraint parameter  $\lambda$  should be of the same scale as the average of the magnitudes of  $A^H D y$ . Though  
 147 our restraint is of a more complex form than  $\|Gx\|_1$ , the dual of our restraint is in the space of a  
 148 combination of two domains similar to  $\|u\|_\infty \leq \lambda$ . Detailed derivations about the dual space of the  
 149 combination of two norm can be found in Rockafellar *et al.*(6). Therefore, the simplified discussion drives  
 150 us to set the parameters of our restraint to be of the scale as the square root of the average of the  
 151 squares of  $A^H D y$ . We denote the square root of the average of the squares of  $A^H D y$  as  $\|A^H D y\|_2$   
 152 henceforth.

153 The scale of implicit gradient ascent restraint is easy to set since it is quadratic. Note that each  
 154 quadratic data loss term in our target function is scaled by  $N(n)$  in Equation 8 in main text. Using the  
 155 heuristic that the penalty term should match the loss term, we can set the restraint parameter  $\gamma$  to be on  
 156 the scale of the average of  $N(n)$ . We denote the average of  $N(n)$  as  $\bar{N}(\bar{n})$  henceforth.

157 Other important parameters to be set are  $\epsilon$  and  $\epsilon'$ , which are in the denominators of our sparsity  
 158 and smoothness restraint, respectively. If they are too small compared to the density values of the 3D  
 159 volume, the weights in our weighted norms will be very flexible and strongly depend on the magnitude of  
 160 the value of each voxel in the 3D volume. Such kinds of restraints might not be able to effectively remove  
 161 background noises and cause two independent refinements to diverge. If they are too large compared to  
 162 the density values of the 3D volume, the restraints degrade to the original  $l_1$  and TV norms and leads to  
 163 more biased solutions. Optimal values of  $\epsilon$  and  $\epsilon'$  should assign large weights to background noises and  
 164 small weights to true molecular densities. Hence, we can set  $\epsilon$  to the level of density values  
 165 corresponding to molecular content in the 3D volume.  $\epsilon'$  can be set close to  $\epsilon$ . This level can be easily  
 166 obtained from the intermediate volumes generated by the refinements using RELION 3.0. This is also  
 167 similar to choose a threshold for creating a mask when computing masked FSC.

168 In conclusion, we should set the restraint parameters  $\alpha$  and  $\beta$  to be on the same scale as  
 169  $\|A^H D y\|_2$ . Since the corresponding restraints are inversely weighted by quantities with two other  
 170 parameters  $\epsilon$  and  $\epsilon'$ , we multiply the scale  $\|A^H D y\|_2$  by  $\epsilon$  or  $\epsilon'$  to counter acting the effects of  $\epsilon$  and  $\epsilon'$ . For  
 171 zero elements, their restraint parameters are then normalized to be on the scale of  $\|A^H D y\|_2$ .

## 172 Experiment process

173  $\beta$ -galactosidase (EMPIAR-10017) : Since there was no ready-to-use particle stack for model  
 174 building, our test began with extracting particles from micrographs using the coordinates manually picked  
 175 by Richard Henderson(7). We carried out 3D refinement in RELION 3.0 with a 50 Å low-pass filtered initial  
 176 map while enforcing D2 symmetry. Then, we performed further 3D refinement using OPUS-SSRI with a  
 177 grid search to determine the possible ranges of parameters  $\alpha$ ,  $\beta$ ,  $\gamma$  and  $\epsilon$ . We started by setting  $\epsilon$  to 0.1,  
 178 which is higher than the level of density values of the EM map from RELION 3.0. We considered setting  $\epsilon'$   
 179 to be  $\epsilon/3$  since the magnitude of gradient is often smaller than the density value of molecule. The initial  
 180 guesses for  $\beta$  and  $\gamma$  were  $\beta = 2/3\|A^H D y\|_2\epsilon$ ,  $\gamma = 0.05\bar{N}(\bar{n})$ . We scanned through  $\alpha \in [0.4, 0.6]\|A^H D y\|_2\epsilon$   
 181 with a step size of  $0.1\|A^H D y\|_2\epsilon$ . The final resolutions based on gold standard FSC=0.143 for different  
 182  $\alpha$  is shown in **Fig.S4a**. The best resolution was obtained at  $\alpha = 0.5\|A^H D y\|_2\epsilon$ . We then set  $\alpha =$   
 183  $0.5\|A^H D y\|_2\epsilon$  and scanned through  $\beta \in [1/3, 2.2/3]\|A^H D y\|_2\epsilon$  with a step size of  $0.2/3\|A^H D y\|_2\epsilon$ . The  
 184 best resolution was obtained at  $\beta = 1.6/3\|A^H D y\|_2\epsilon$  or  $1.8/3\|A^H D y\|_2\epsilon$  (**Fig.S4b**). To calculate the model  
 185 vs. map FSCs, we fitted the atomic coordinates of an *E. coli*  $\beta$ -galactosidase structure (PDB code  
 186 3I3E)(8) into the post-processed density maps reconstructed by different methods.

187 80S ribosome (EMPIAR-10002): We extracted particles from the micrographs using the semi-  
 188 automated selection process in RELION 3.0(9). The particles were pruned by one round of 2D  
 189 classification where only the particles classified to major classes were kept. We then constructed an *ab*  
 190 *initio* map in RELION 3.0 through 3D classification. 3D refinements continued from the 70 Å low pass  
 191 filtered initial map. For OPUS-SSRI, we used the same  $\epsilon$ ,  $\epsilon'$  and  $\gamma$  as  $\beta$ -galactosidase and scanned  
 192 through  $\alpha \in [0.3, 0.5]\|A^H D y\|_2\epsilon$  with a step size of  $0.1\|A^H D y\|_2\epsilon$ . The resolutions with different  $\alpha$  are

193 shown in **Fig.S4c**. The best result is obtained at  $\alpha = 0.4\|A^H D y\|_2 \epsilon$ . We then scanned through  $\beta \in$   
194  $[1.2, 2.4]$  with a step size of  $0.4/3\|A^H D y\|_2 \epsilon$  for  $\alpha = 0.4\|A^H D y\|_2 \epsilon$  and  $\alpha = 0.5\|A^H D y\|_2 \epsilon$ . The results are  
195 shown in **Fig.S4d**. The best results were obtained at  $\alpha = 0.4\|A^H D y\|_2 \epsilon$  and  $\beta = 1.6/3\|A^H D y\|_2 \epsilon$  or  
196  $2/3\|A^H D y\|_2 \epsilon$ . To assess the cryo-EM maps determined using different methods, we fitted 80S crystal  
197 structure (PDB code 3U5B(10)) using a simple rigid-body fit into post-processed maps. The 40S and 60S  
198 subunits were fitted separately.

199 Hemagglutinin (EMPIAR-10097): The 3D refinements were performed by using a 40 Å low-  
200 passed filtered initial map and enforcing the C3 symmetry. For OPUS-SSRI, the optimal parameters are  
201  $\epsilon = 0.015$  and  $\epsilon' = \epsilon/3$ ,  $\alpha = 0.4\|A^H D y\|_2 \epsilon$ ,  $\beta = 2.6/3\|A^H D y\|_2 \epsilon$  and  $\gamma = 0.2\overline{N(n)}$ . To compare the post-  
202 processed density maps obtained from different refinement methods, we used a crystal structure of HA  
203 trimer (PDB code 3WHE)(11) as a reference following the original publication(12).

204 TRPM4 (EMPIAR-10126): The 3D refinement rounds were performed with a 50 Å low-passed  
205 filtered initial map and C4 symmetry. The best parameters of our method were  $\epsilon = 0.01$ ,  $\epsilon' = \epsilon$ ,  $\alpha =$   
206  $0.8\|A^H D y\|_2 \epsilon$ ,  $\beta = 4\|A^H D y\|_2 \epsilon$  and  $\gamma = 0.1\overline{N(n)}$ . The cryo-EM maps were compared with the atomic  
207 model in PDB code 6BQR(13).

208 Hrd1/Hrd3 (EMPIAR-10099): Due to the heterogeneity of this dataset, the 3D classification was  
209 used to classify the particles, and generate the corresponding initial maps of different complexes for 3D  
210 refinements. The particles classified as Hrd1/Hrd3 dimer were selected, and then subject to 3D  
211 refinements. We performed 3D refinements using different methods and the same 20 Å low-passed  
212 filtered initial map. The best parameters of our method are  $\epsilon = 0.01$ ,  $\epsilon' = \frac{\epsilon}{2}$ ,  $\alpha = 0.4\|A^H D y\|_2 \epsilon$ ,  $\beta =$   
213  $0.6\|A^H D y\|_2 \epsilon$  and  $\gamma = 0.1\overline{N(n)}$ . The final results were compared with the atomic models of Hrd1 dimer  
214 (PDB code 5V6P) and Hrd3 monomer (PDB code 5V7V)(14). 5V6P and 5V7V were fitted into density  
215 map separately.

216 TRPV5 (EMPIAR-10254): The 3D refinement rounds were performed by using a 40 Å low-pass  
217 filtered 3D initial map and enforcing C4 symmetry. The best parameters of our method were  $\epsilon = 0.01$ ,  $\epsilon' =$   
218  $2\epsilon$ ,  $\alpha = 0.5\|A^H D y\|_2 \epsilon$ ,  $\beta = 2\|A^H D y\|_2 \epsilon$  and  $\gamma = 0.1\overline{N(n)}$ . The cryo-EM maps were compared to the atomic  
219 model in PDB code 6O1P(15).

220 TMEM16A (EMPIAR-10123): The 3D refinements were performed by using a 40 Å low-passed  
221 filtered initial map and enforcing the C2 symmetry. The best parameters of our method are  $\epsilon = 0.01$ ,  $\epsilon' =$   
222  $\epsilon$ ,  $\alpha = 0.5\|A^H D y\|_2 \epsilon$ ,  $\beta = 2\|A^H D y\|_2 \epsilon$  and  $\gamma = 0.1\overline{N(n)}$ . The final results were compared to the atomic  
223 model in PDB code 6BGI(16) by calculating model vs. map FSC.  
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**Table S1. Comparison of the final reconstructions refined by RELION 3.0, THUNDER or OPUS-SSRI**

Proteins	Model vs. Map FSC=0.143					
	RELION	THUNDER		OPUS-SSRI		
	Resolution (Å)	Resolution (Å)	$\Delta\text{Å}$ over RELION <sup>1</sup>	Resolution (Å)	$\Delta\text{Å}$ over RELION <sup>1</sup>	$\Delta\text{Å}$ over THUNDER <sup>2</sup>
<i>β</i> -galactosidase (EMPIAR-10017)	4.05	4.08	-0.03	3.91	0.14	0.17
80S ribosome (EMPIAR-10002)	4.04	3.70	0.34	3.89	0.15	-0.19
Hemagglutinin (EMPIAR-10097)	4.06	3.89	0.17	3.72	0.34	0.17
TRPM4 (EMPIAR-10126)	3.34	/	/	2.93	0.41	/
Hrd1/Hrd3 (EMPIAR-10099)	4.70	4.88	-0.18	4.25	0.45	0.63
TRPV5 (EMPIAR-10254)	3.05	2.98	0.07	2.76	0.29	0.22
TMEM16A (EMPIAR-10123)	3.87	/	/	3.14	0.73	/
<b>Average improvement</b>			<b>0.07</b>		<b>0.30</b>	<b>0.20</b>

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<sup>1</sup>: the value in negative indicates the resultant resolution is worse than that from RELION, while the value in positive indicates the resultant resolution is better than that from RELION.

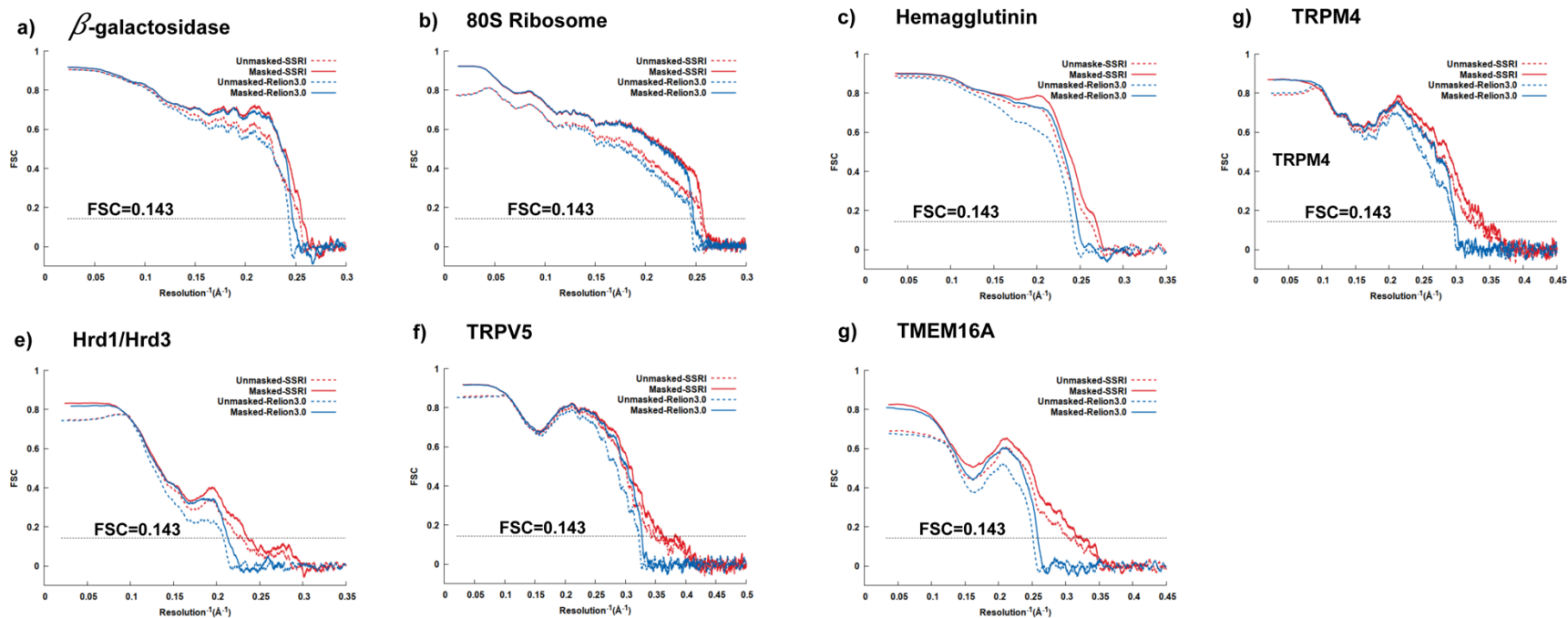
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<sup>2</sup>: the value in negative indicates the resultant resolution is worse than that from THUNDER, while the value in positive indicates the resultant resolution is better than that from THUNDER.

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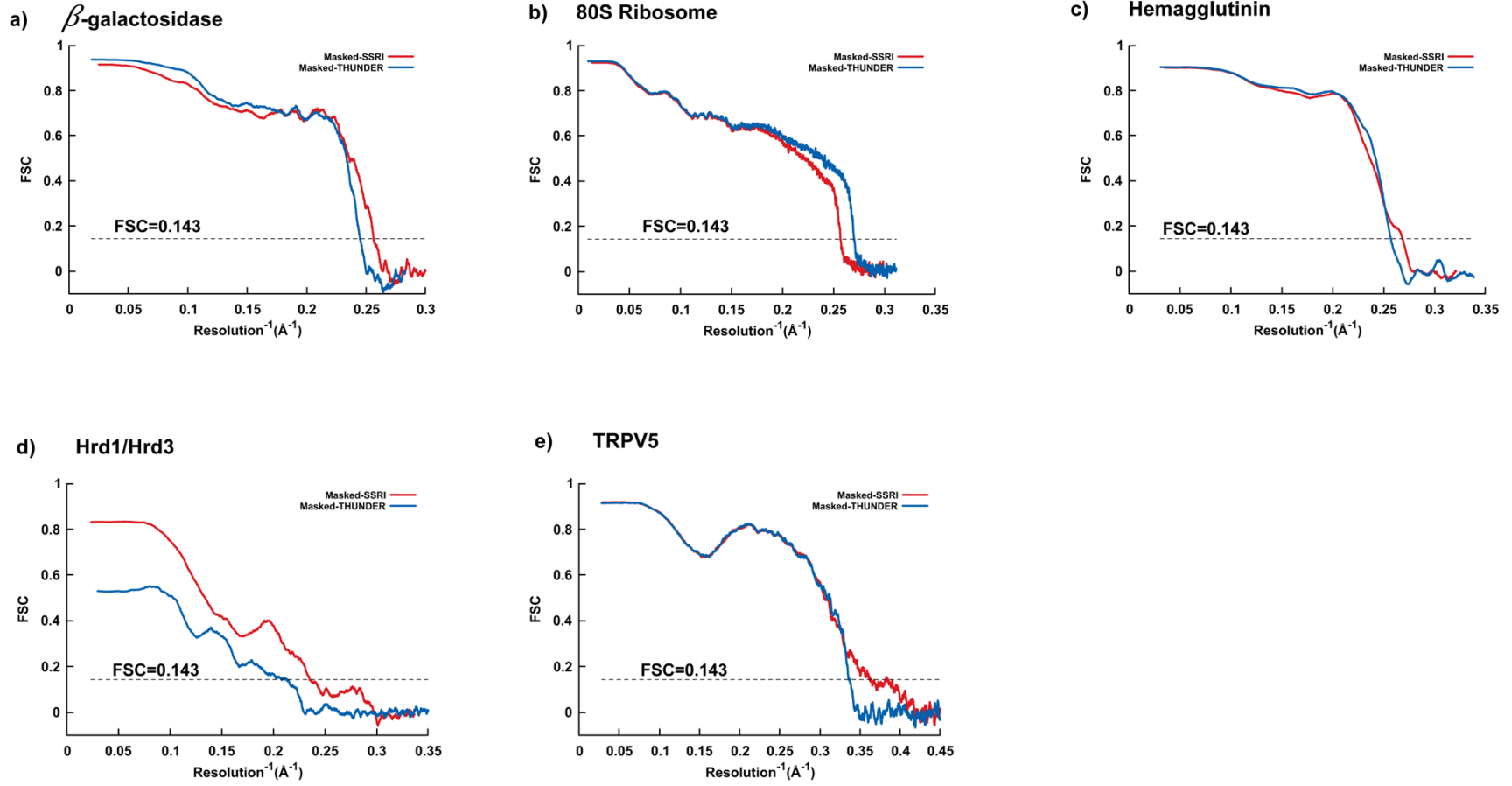
/: indicates that the comparison was unavailable in two cases where THUNDER failed to execute due to computer incompatibility.

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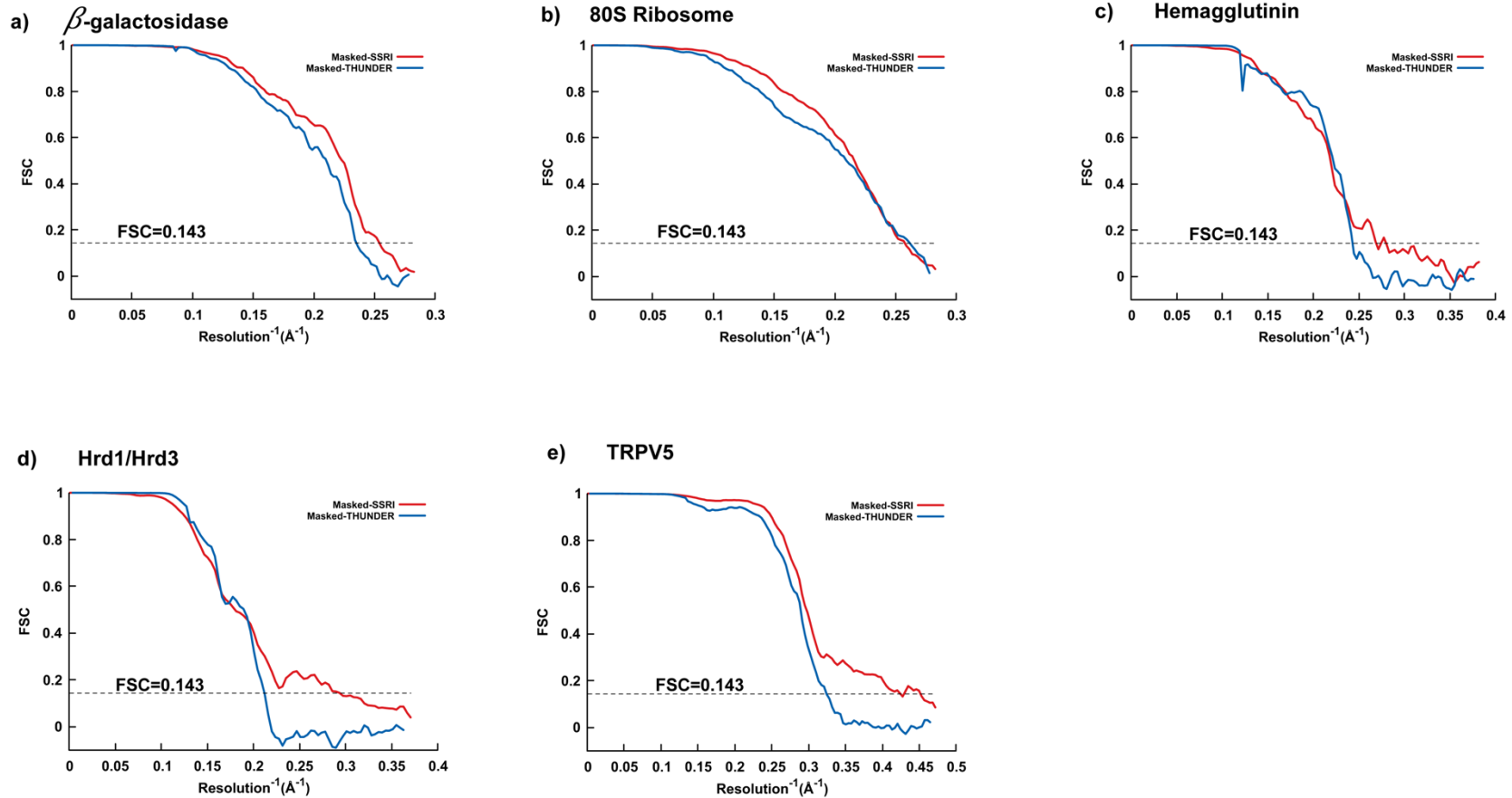


**Fig. S1.** Model vs. map FSC curves between the post processed density maps and the corresponding rigid-body fitted atomic models for seven proteins. **(a)**  $\beta$ -galactosidase with atomic model in PDB code 3I3E; **(b)** 80S ribosome with atomic model in PDB code 3U5B, for which the 40S and 60S subunits were fitted separately; **(c)** Influenza hemagglutinin with atomic model in PDB code 3WHE; **(d)** TRPM4 with atomic model in PDB code 6BQR; **(e)** Hrd1/Hrd3 complex, for which the Hrd1 dimer (PDB code 5V6P) and Hrd3 monomer (PDB code 5V7V) were fitted separately; **(f)** TRPV5 with atomic model in PDB code 6O1P; **(g)** TMEM16A with atomic model in PDB code 6BGI. Dashed black line represents FSC=0.143 in all panels.

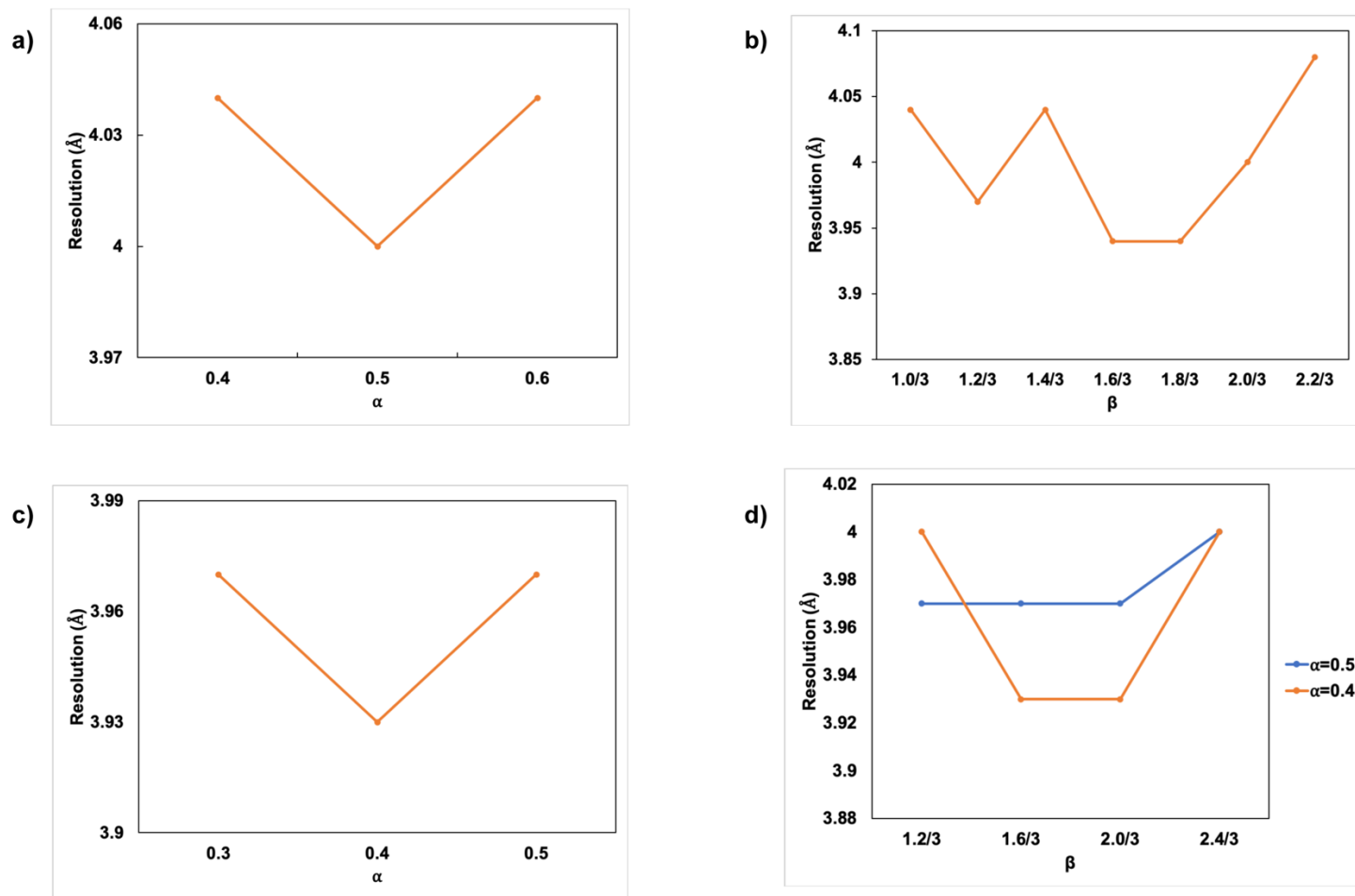




**Fig. S2.** Gold standard masked FSC curves for the final 3D reconstructions refined by OPUS-SSRI or THUNDER for **(a)**  $\beta$ -galactosidase, **(b)** 80S ribosome, **(c)** influenza hemagglutinin, **(d)** Hrd1/Hrd3 and **(e)** TRPV5. In all panels, dashed black line represents FSC=0.143.



**Fig. S3.** Model vs. map FSC curves between the post processed density maps and the corresponding rigid-body fitted atomic models for five proteins. **(a)**  $\beta$ -galactosidase with atomic model in PDB code 3I3E; **(b)** 80S ribosome with atomic model in PDB code 3U5B, for which the 40S and 60S subunits were fitted separately; **(c)** Influenza hemagglutinin with atomic model in PDB code 3WHE; **(d)** Hrd1/Hrd3 complex, for which the Hrd1 dimer (PDB code 5V6P) and Hrd3 monomer (PDB code 5V7V) were fitted separately; **(e)** TRPV5 with atomic model in PDB code 6O1P. Dashed black line represents FSC=0.143 in all panels.



**Fig. S4.** Grid search for optimal parameters in OPUS-SSRI. **(a)** Resolutions of EM maps refined with different  $\alpha$  for  $\beta$ -galactosidase. **(b)** Resolutions of EM maps refined with different  $\beta$  for  $\beta$ -galactosidase. **(c)** Resolutions of EM maps refined with different  $\alpha$  for 80S ribosome. **(d)** Resolutions of EM maps refined with different  $\beta$  and  $\alpha$  for 80S ribosome.

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