1	Supplementary Information for
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3	Ionic contrast across a lipid membrane for Debye length extension
4	: towards an ultimate bioelectronic transducer
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23 This file contains following information:

24 Supplementary Notes

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40 Supplementary Note 1. Device fabrication and measurement setup



43 Supplementary Fig. 1. Schematic representation of fabrication of FETs for molecular detection
44

45 Extended gate (EG) fabrication: The top gate electrode of the FET is wired to an indium tin 46 oxide (ITO, thickness: 300 nm and resistivity: 5 Ω/sq) electrode. A 80-nm-thick SiO₂ layer was 47 deposited on the ITO for bonding the reaction chamber to the ITO electrode. The reaction chamber 48 was made of polydimethylsiloxane (PDMS, Sylgard 184 silicon elastomer kit, Dow Corning). The 49 PDMS mixture (base:curing agent = 10:1) was poured into a Petridish and baked for 12 h at 80° C. 50 After curing, an 8-mm diameter and a 5-mm height hole was formed in the PDMS mold using a 51 biopsy punch. Note that the SiO_2 layer plays a bifunctional role here: (1) a surface for forming the supported lipid bilayer (SLB) membrane and (2) serving as a hydrophilic polar silanol (Si-OH) 52 53 group cross-linked surface for the PDMS chamber via covalent bonding after an oxygen plasma 54 treatment¹.

Importance of disposable EG: The EG reaction chamber is disposable. This measurement configuration for the EG and the remotely packaged FET allows for reliable data acquisition, eliminating any possible contamination or damage to the active current channel upon direct exposure to target analytes in ionic solution. Moreover, direct comparison between measurements is possible because an identical FET transduces signals from the disposable EG. This is important during quantification of analyte concentration.



62

63 Supplementary Fig. 2. Photograph images of the EG with the reaction chamber placed in a
64 Faraday shielding box

Back gate bias voltage conditions: In our measurement configuration, the potential variation ($\Delta \psi$) of the sensing probe (EG) is capacitively coupled to top gate voltage (V_{TG}), modulating drain current (I_D) in the FET. The proportionality between the variations of V_{TG} and I_D is defined by transconductance (g_m),

$$g_{\rm m,max} = g_{\rm m}|_{V_{\rm TG} = V_{\rm th}} = \frac{\mathrm{d}I_{\rm D}}{\mathrm{d}V_{\rm TG}} \tag{1}$$

where dI_D and dV_{TG} are the small variation of the drain current and the top gate voltage. As seen in the Supplementary Eq. 1, the g_m represents the amplification of the output response in the transducing process from the molecular bindings. Namely, the dV_{TG} caused by the molecular binding is transduced into the obtained current variation dI_D . The g_m is typically maximized at around the threshold voltage V_{th} . However, this bias voltage condition to maximize the g_m is affected by the ionic strength of the analyte solution.

In our dual gate configuration, we can address this issue systematically by tuning back gate bias voltage V_{BG} . As shown in Supplementary Fig. 3, the threshold voltage (= $g_{m,max}$) can be shifted by the V_{BG} . Thus, we characterized the FET in the measurement condition to determine the back gate bias voltage for the maximized g_m . With this optimizing process, the applied g_m is ensured to be maximized throughout the measurements. For instance, the g_m was maximized with the V_{BG} of 13.5 V in the DIW environment.



Supplementary Fig. 3. I_D - V_{TG} curves at various back gate voltages V_{BG} . ($V_S = 1$ V)

86 Supplementary Note 2. Lipid preparation and fluorescence (FL) measurement

88	FL imaging and mobility test: For the FL imaging, we monitored the SLB using an exposure
89	time of 100 ms and gain 5.8× for every FL micrograph in the SLB coverage experiments. The FL
90	and bright-field (BF) images of particulate materials were monitored using a FL microscopy
91	system (LV100ND, Nikon) and then analyzed using image processing software (ImageJ, National
92	Institutes of Health, USA). For the FRAP test, a 20-µm-circle region of the SLB patch was
93	bleached for 1 min, and time-lapse FL intensities were measured using confocal microscopes (C2
94	C-ER, Nikon, Japan).
95	

97 Supplementary Note 3. Buffer solution exchange

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After the SLB (DOPC:B-PE = 95:5) preparation in DIW, half of the solution in the EG reaction chamber (50 μ L out of 100 μ L) was exchanged with 50 μ L of 1× PBS. This process was repeated until the V_{TG} value was saturated (Supplementary Fig. 4). Note that the signal reached a steady state after 10 repeated exchanges of the outer buffer (OB), which is consistent with the expected concentrations; 0.992× PBS at the 7th buffer solution exchange and 0.999× PBS at the 10th buffer solution exchange (Supplementary Fig. 4).



105

Supplementary Fig. 4. Potential variation of the top gate electrode during the OB exchange
process with the ionic contrast across the SLB

109 Supplementary Note 4. Control experiment with identical ionic conditions across the SLB

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111 To validate the role of the DIW layer between the SiO₂ surface and the SLB, a negative control 112 experiment was carried out using symmetrical ionic conditions across the SLB. Namely, the SUV 113 was prepared in $1 \times PBS$ (0.1 mg/mL) and exposed to the hydrophilic SiO₂ surface of the EG. After 114 SLB formation (marker 2 in Supplementary Fig. 5), we waited for at least 5 min for signal 115 stabilization. No signal change was observed after the 800 pM-avidin injection (marker 3 in 116 Supplementary Fig. 5), in contrast to the recognizable ΔV_{TG} changes under 800 pM-avidin 117 bindings under asymmetric ionic conditions across the SLB (main manuscript Figs. 3a and h). This 118 is attribute to the ultra-thin Debye length (approximately 0.74 nm in $1 \times PBS$), which is thinner 119 than IB layer (1–2 nm) between SLB and ITO surface.

120



122 Supplementary Fig. 5. Real-time ΔV_{TG} measurement of the response to biotin–avidin reactions

- 123 under identical ionic conditions across the SLB
- 124



- **Supplementary Fig. 6.** Representation of the liquid cell for X-ray reflectivity measurement and
- 127 the SLB on a SiO_2 wafer placed in the cell

129 Supplementary Note 5. Minimum molecular mass for the full coverage of the SLB

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131 We estimated the required number of lipid molecules to cover the entire surface of the SiO₂ 132 and the inside wall of the reaction chamber A_{EG} ,

133
$$A_{\rm EG} = \pi r^2 + 2\pi r h = 125.27 \,({\rm mm}^2)$$
 (2)

where *r* and *h* are the diameter and height of the EG chamber, respectively. Considering the head area of DOPC molecules² ($A_{SLB} = 72.4 \text{ Å}^2$), the required number of DOPC molecules was calculated to be

137
$$\frac{A_{\rm EG}}{A_{SLB}/2} = \frac{125.27 \ (mm^2)}{36.2 \ ({\rm \AA})} = 0.57 \times 10^{-9} \ ({\rm mol}). \tag{3}$$

138 With an SLB composition of DOPC/B-PE = 95/5, the averaged mass of the SLB is

139
$$\left(\underbrace{\frac{786.1 \times 0.95}{DOPC} + \underbrace{1105.5 \times 0.05}_{B-PE}}\right)(g \ mol^{-1}) \times (0.57 \times 10^{-9})(mol) \approx 0.46 \ \mu g.$$
 (4)

This means that $0.46 \ \mu g$ of lipids were required for full coverage of the EG including the PDMS wall and the SiO₂ surface. Considering that the mass of the lipid in solution was $0.1 \ \mu g$ for every drop, the observed saturation after five lipid injections (markers 2 to 6 in main manuscript Fig. 2b) seems reasonable amount for the full coverage of the EG surface. Additional casting of lipid molecules (marker 7 in the main manuscript Fig. 2b) confirms that 5 times of lipid introduction was sufficient for the complete passivation of the EG surface.

147 <u>Supplementary Note 6. Langmuir–Nernst isotherm model considering nonspecific bindings</u> 148

The SLB allows covering the EG surface, and thereby prohibit the nonspecific binding of analytes to the EG surface. To investigate and quantify the role of SLB as the 'suppressor' for the nonspecific binding, we conducted an experiment with the EG functionalized by sulfo-NHS-biotin and compared this with the isotherm analysis of the SLB-FET (main manuscript Fig. 3d). The Langmuir–Nernst isotherm model is described as follows:

155
$$\Delta V_{\rm TG}([A]) = \frac{q_A}{C_{TG}} [B]_{\rm max} \frac{[A]}{[A] + K_{eq}} + k_N[A]$$
(5)

154 where K_{eq} , and k_N stand for equilibrium constant and nonspecific binding constant, respectively³. 156 The strong suppression of nonspecific binding in our SLB-FET is described in the main 157 manuscript in Fig. 3d. The measured result (circles in orange in main manuscript Fig. 3d) from our 158 SLB-FET is highly consistent with the conventional Langmuir isotherm model (solid orange line 159 in main manuscript Fig. 3d, $k_N[A] = 0$). Biotin-avidin binding signals measured with the sulfo-160 NHS-biotin functionalized EG (circles in grey in main manuscript Fig. 3d) exhibited strong 161 agreement with the Langmuir-Nernst isotherm (solid grey line in main manuscript Fig. 3d). From 162 the fitting, k_N and K_{eq} were calculated to be 3.685 \pm 0.754 mV/ μ M·m and 54.2 fM, respectively, which was within the reasonable range⁴. Note that our SLB-FET platform allowed for negligence 163 164 of the second term, whereas protein bindings without SLB should take into account the possibility 165 of nonspecific binding.

166 Supplementary Note 7. Sensing mechanism and analytical models

167

Molecular detection using FETs is based on conductivity changes in response to variance in the gate potential. In our case, using an n-type channel for FET and applying a positive gate voltage leads to accumulation of carriers. When analytes bind to receptors on the SLB, the charge density variation in the SiO₂ layer changes the top gate voltage, and eventually the conductance in the FET. The time-dependent top gate voltage induced by adsorbed molecules is described by Supplementary Eq. 6,

174
$$\frac{d\Delta V_{\rm TG}(t)}{dt} = \frac{q_A}{C_{\rm TG}}[AB]$$
(6)

175 where [*AB*] is the density of an adsorbed analyte.

176 In general, the binding kinetics of the receptor when capturing an analyte in solution can be177 described as follows:

178

$$[A] \stackrel{k_M}{\leftrightarrow} [A]_{\mathsf{S}} + [B] \stackrel{k_1, \ k_{-1}}{\longleftrightarrow} [AB] \tag{7}$$

where k_M is a diffusion-limiting rate constant, k_I and k_{-1} are association and dissociation rate constants, respectively, and $[A]_s$ is the surface concentration of an analyte⁵. We assumed the fast mixing model ($[A]_s = [A]$) to describe our SLB-FET system, because the k_I in the biotin–avidin reaction (10^5 to 10^7 M⁻¹s⁻¹) is two or more orders of magnitude smaller than the k_M (~ 10^9 M⁻¹s⁻¹) for diffusion limited reactions⁶. In other words, we suppose that the analyte concentration at the surface remains almost the same as in the bulk in the following analysis. With this, the reaction can be described in a simpler fashion by the first order Langmuir equation (Supplementary Eq. 8a).

186
$$\frac{d[AB]}{dt} = k_1[A]([B]_{\max} - [AB]) + k_{-1}[AB]$$
(8a)



188
$$[AB]_{t} = \frac{k_{1}[B]_{\max}[A]}{k_{1}[A] + k_{-1}} \left(1 - e^{-(k_{1}[A] + k_{-1})t}\right) = [B]_{\max} \frac{[A]}{[A] + K_{eq}} \left(1 - e^{-(k_{1}[A] + k_{-1})t}\right)$$
(8b)

By combining Eqs. 6 and 8b, the time-dependent voltage variation is described by Supplementary Eq. 9,

191
$$\Delta V_{\text{TG}_{adsorption}}(t) = \frac{q_A}{C_{\text{TG}}} [B]_{\max} \left(1 - e^{-(k_1[A] + k_{-1})t} \right)$$
(9)

In our measurements, the target analyte in solution is dropped in the reaction chamber; therefore, this inevitably leads to a sudden increase in the ionic concentration in the outer buffer. Thus, the top gate potential should reflect relaxation behavior via thermal diffusion. We model its transient state towards equilibrium by a resistor–capacitor (RC) circuit model as in Supplementary Eq. 10.

197
$$\Delta V_{\text{TG}_{\text{RC}}}(t) = \begin{cases} V_p (1 - e^{-\frac{t}{\tau_1}}) & t < T \\ V_p e^{-\frac{t}{\tau_2}} & t \ge T \end{cases}$$
(10)

198 where the peak voltage V_p is the difference between $q_A[B]_{\text{max}}/C_{\text{TG}}$ and the maximum value of 199 $\Delta V_{\text{TG}}(t)$, τ_1 and τ_2 are RC time constants, and *T* is an effective duration for the sudden increase in 200 the ionic concentration in the outer buffer.

Finally, the measured signal should be the superposition of these responses described above (Eqs. 9 and 10).

203
$$\Delta V_{\rm TG}(t) = \begin{cases} \frac{q_A}{C_{\rm TG}} [B]_{\rm max} \left(1 - e^{-(k_1[A] + k_{-1})t}\right) + V_p (1 - e^{-\frac{t}{\tau_1}}) & t < T\\ \frac{q_A}{C_{\rm TG}} [B]_{\rm max} \left(1 - e^{-(k_1[A] + k_{-1})t}\right) + V_p e^{-\frac{t}{\tau_2}} & t \ge T \end{cases}$$
(11)

Supplementary Fig. 7 shows the response curve for 100 pM-avidin binding over the SLB (DOPC:B-PE = 95:5), fitted to our suggested theory model. From these fittings, important reaction parameters are obtained. Those parameters are summarized in Supplementary Table 1.





Supplementary Fig. 7. Real-time top gate voltage variation in response to 100 *pM*-avidin fitted
by an RC circuit model and the first order Langmuir equation.

Supplementary Table 1. Fitting parameters in Supplementary Eq. 10

	V_p	$ au_1$	$V_{ m eq}$	k_{I}	<i>k</i> -1	$ au_2$	Т
	(mV)	(sec)	(V)	$(\times 10^7 M^{-1} s^{-1})$	(×10 ⁻² s ⁻¹)	(sec)	(sec)
100 pM	1.3200	29.682	0.0022	1.6835	1.0469	231.790	432.984
1 nM	0.7802	12.971	0.0053	1.6642	1.0350	106.525	219.648
5 nM	-	-	0.0079	1.6606	1.0311	-	-
10 nM	-	-	0.0088	1.6227	1.0010	-	-
100 nM	-	-	0.0089	1.7163	1.0410	-	-

213 Supplementary Note 8. X-ray reflectivity measurement for electron density fitting



214

215 Supplementary Fig. 8. Structure for a slab model. a, the SLB and b, the SLB with a layer of





217

Supplementary Fig. 9. X-ray intensities and electron density profiles of SLB (blue) and SLB
with avidin (yellow). a, X-ray reflectivity intensities and their fittings. b, Electron density curves
along the z-direction.

As shown in Supplementary Fig. 9a, both measurements of 'SLB' and 'SLB with avidin' were well fitted by the model described above. Slight deviation at high q regime (> 0.5 Å^{-1}) is due to the low peak intensity. The electron density plots in Supplementary Fig. 9b clearly exhibit the typical electron density curve of SLB, showing consistency with the literature in terms of the size of the conjugated avidin² (~50 Å).

227 Supplementary Note 9. Chemical potential calculation using the electron density

Poisson's equation (Supplementary Eq. 12) was applied to calculate electric potential variation due to the redistribution of the charge density in the lipid membrane^{7, 8, 9}.

231
$$\psi(z) - \psi(0) = -\frac{1}{\varepsilon_0} \int_0^z dz' \int_0^{z'} \rho(z'') dz''$$
(12)

232 The chemical potential at the probing surface was calculated by integrating the charge density 233 $(\Delta \rho = \rho_{SLB} - \rho_{SLB+avidin})$, and multiplied by the area of SLB film ($\pi \times R^2 \approx 5.027 \times 10^{-5} \text{ m}^2$) assuming 234 in-plane homogeneity.

235
$$\Delta \psi(z) - \Delta \psi(0)$$

236
$$= -\frac{1}{\varepsilon_0} \int_0^z \int_0^{z'} \Delta \rho(z'') dz'' dz'$$

237
$$= -\frac{e^{-} \times \pi r^{2}}{\varepsilon_{0}} \int_{0}^{z} \int_{0}^{z'} \Delta \rho(z'') dz'' dz'$$

238
$$= -\frac{1.602 \times 10^{-19} \times 5.027 \times 10^{-5}}{8.854 \times 10^{-12}} \int_0^z \Delta E(z') - \Delta E(0) dz' = 0.697 \text{ (V)} (13)$$

As seen in the main manuscript Fig. 4d, the obtained signal mainly originates from the charge density change in proximity of the probing surface because the chemical potential is inversely proportional to the distance squared.

242

228

Supplementary Table 2. Electric parameters of FET-SLB circuit components

	$C_{ m EG}$	$C_{ m TG}$	$C_{ m BG}$	$C_{ m SLB}$
Dielectric constant (<i>E</i> _r)	3.9	3.9	3.9	1.9
Area (m ²)	5.03E-05	1.50E-10	1.50E-10	5.03E-05
thickness (m)	8.00E-08	1.50E-08	7.50E-07	5.00E-09
Capacitance (F)	2.17E-08	3.45E-13	6.91E-15	4.36E-06
Voltage (V)	2.18E-07	1.37E-02	6.83E-01	1.08E-09

243 Supplementary Note 10. Comparison of lipid membranes on a hydrophilic polymer layer vs

244 <u>a SiO₂ layer</u>

Hydrophilic SAM on metal surface



50 nm SiO2 film on metal surface



245

246 Supplementary Fig. 10. Direct comparison of lipid membrane over the hydrophilic polymer

vs the SiO₂ layer. a, Non-uniform FL distribution of supported lipid membrane on the hydrophilic

248 polymer (2-Mercaptoethanol) on the EG. b, corresponding schematic illustration. c, Partial 249 immobility of lipid compositions during a FRAP test. Albeit become brighter, the circular FL 250 bleached area does completely recovered in 10 min. **d**, Time-dependent ΔV_{TG} after the outer buffer 251 exchange (DIW \rightarrow 1× PBS) showing the leakage flow of ions through the pin-hole of the lipid 252 membrane. e, f, Uniform SLB formation over the 50-nm-SiO₂ layer (e) with a schematic 253 illustration (f). g, FRAP test showing later fluidity of the composing lipids. Note that the FL 254 bleached (black) region disappeared within 10 min. **h**, Time-lapse ΔV_{TG} variation after the OB 255 exchange (DIW \rightarrow 1× PBS). The ΔV_{TG} was saturated after the ionic contrast between IB (DIW) 256 and OB (1× PBS) was obtained.

257

This is the major reason why we prepared the SLB on the SiO₂ film. Although mass production of biosensors using polymeric SAM seems far-fetched, our suggested SLB-FET with the SiO₂ film on the EG guarantees the robust formation of the SLB in terms of coverage (Supplementary Fig. 10e) and lipid mobility (Supplementary Fig. 10g). Most importantly, stable OB exchange manifests the importance of the defect-free SLB. This allows reliable molecular detection with high reproducibility in real time (Supplementary Fig. 10h).

264 Supplementary Note 11. Molecular dynamics (MD) simulation for investigating the



266

267 Supplementary Fig. 11. MD simulation on electron density and conformational change in the 268 SLB membrane (DOPC:B-PE=95:5) upon biotin-avidin binding. a. Final configurations of the SLB before (left) and after (right) avidin binding. The DOPC lipids, B-PE lipids, avidin, and water 269 270 molecules are represented by the green, red, purple, and cyan colours, respectively. The SiO₂ 271 surface are represented by red (O atoms) and yellow (Si atoms) spheres. b. Electron density profiles 272 of the upper leaflet in the membrane before and after biotin-avidin bindings. c. Potential difference 273 $(\Delta \psi)$ upon biotin-avidin binding calculated by Poisson's equation. **d**. Top view of avidin-bound 274 membrane and corresponding area per lipid calculated by Voronoi analysis. e. Chart for calculated 275 areas per lipid for constituent lipids.

276 Supplementary Fig. 11 shows the obtained simulation results for the final configurations of binary 277 lipid membrane (DOPC:B-PE=95:5) upon avidin binding events under a symmetrical ionic 278 condition across the lipid bilayer (DIW for both of IB and OB). We then calculated the electron 279 density profile of lipids to investigate whether the avidin-biotin binding event can modulate the 280 profiles of the lipid membrane. The electron density profile achieved from the centre of the lipid 281 membrane shows that the maximum peak of electron density at the upper leaflet of the lipid bilayer 282 shifts toward a higher value of z after avidin binding, indicating slight thickening of the lipid 283 bilayer (Supplementary Fig. 11b). For quantitative comparison with experimental results, the 284 potential difference, $\Delta \psi$, via the avidin binding was calculated using Poisson's equation 285 (Supplementary Fig. 11c). As a result, the avidin bound membrane induced chemical potential 286 change of $\Delta \psi = 499$ mV at the surface of the SiO₂ layer. This agrees very well with the 287 experimental results shown in main Fig. 4d. We further measured the structural characteristic of 288 the lipid bilayer, the area per lipid, A_p , to quantify the conformational change of the membrane 289 upon avidin binding. Top view of the avidin bound DOPC/B-PE membrane and corresponding 290 Voronoi cells represented for DOPC and B-PE are shown in Supplementary Fig. 11d. This shows 291 in-plane (x-y plane) compression of lipid membrane when two biotins bound for two binding sites 292 of an avidin (indicated as 'Avidin-B-PE', Supplementary Fig. 11d). The color map analyzed in the 293 Voronoi cell (Supplementary Fig. 11d) and corresponding data in area per lipid (Supplementary Fig. 11e) clearly indicate that avidin bound B-PEs become the seed spots causing the membrane 294 295 packing effect (a decreased area per lipid). Successive decrease of the area per lipid around the 296 avidin bound B-PE accompanies the thickening effect over the lipid membrane, which results in a 297 potential decrease of the SiO_2 surface. Considering that the packing density is one of the essential 298 criteria in defining the modulation of the lipid membrane during the avidin binding, the MD

- 299 simulation results with DOPC/B-PE lipid mixtures shows great correspondence with the electron
- 300 density profiles obtained from XRR.





303 Supplementary Fig. 12. MD simulation on electron density and conformational change in the 304 SLB (DPPC:B-PE=95:5) upon biotin-avidin binding. a. Final configurations of the SLB before 305 (left) and after (right) avidin binding. The DPPC lipids, B-PE lipids, avidin, and water molecules 306 are represented by the green, red, purple, and cyan colours, respectively. The SiO₂ surface are 307 represented by red (O atoms) and yellow (Si atoms) spheres. b. Electron density profiles of the upper leaflet in the membrane before and after biotin-avidin bindings. c. Top view of the avidin-308 309 bound membrane and corresponding area per lipid calculated by Voronoi analysis. d. Chart for 310 calculated areas per lipid for constituent lipids.

312 Supplementary Fig. 12a shows the final configurations of the DPPC lipid mixed with B-PE 313 (DPPC:B-PE=95:5) before and after the avidin binding. In contrast to the membrane with the 314 mobile DOPC lipid, the stiffer and more crystalized DPPC with B-PE shows a negligible effect on 315 the modulation of the electron density profile (Supplementary Fig. 12b). Top view of the avidin 316 bound DPPC/B-PE membrane, and corresponding Voronoi cells represented the area per lipid, A_p, 317 for DPPC and B-PE are shown in Supplementary Fig. 12c. The area per lipid calculated from the 318 Voronoi cells represents that negligible conformational change of the lipid membrane was 319 observed in the binary DPPC/B-PE membrane after avidin binding (Supplementary Fig. 12d). In 320 summary, insignificant physicochemical modulation was observed in the DPPC/B-PE under 321 avidin binding.

323 **Supplementary References** 324 325 1. Y. S. Ryu *et al.*, Kinetics of lipid raft formation at lipid monolayer-bilayer junction probed 326 by surface plasmon resonance. *Biosens. Bioelectron.* 142, 111568 (2019). 327 2. N. Kucerka, S. Tristram-Nagle, J. F. Nagle, Structure of fully hydrated fluid phase lipid 328 bilayers with monounsaturated chains. J. Membrane Biol. 208, 193-202 (2005). 329 L. Maletinska et al., Human glioblastoma cell lines: Levels of low-density lipoprotein 3. 330 receptor and low-density lipoprotein receptor-related protein. *Cancer Res.* 60, 2300-2303 331 (2000).332 4. R. A. Kohanski, M. D. Lane, Monovalent avidin affinity columns. *Methods Enzymol.* 184, 333 194-200 (1990). 334 5. X. Duan et al., Quantification of the affinities and kinetics of protein interactions using 335 silicon nanowire biosensors. Nat. Nanotechnol. 7, 401-407 (2012). 336 6. R. F. Delgadillo et al., Detailed characterization of the solution kinetics and 337 thermodynamics of biotin, biocytin and HABA binding to avidin and streptavidin. PLoS 338 *One* **14**, e0204194–e0204194 (2019). 339 7. I. D. Mayergoyz, Solution of the nonlinear Poisson equation of semiconductor-device 340 theory. J. Appl. Phys. 59, 195-199 (1986). 341 8. H. J. Shen et al., Molecular dynamics simulations of ether- and ester-linked phospholipid 342 Bilayers: A comparative study of water models. J. Phys. Chem. B 122, 9399-9408 (2018). 343 9. K. E. Forsten, R. E. Kozack, D. A. Lauffenburger, S. Subramaniam, Numerical solution of 344 the nonlinear Poisson-Boltzmann equation for a membrane-electrolyte system. J. Phys. 345 Chem. 98, 5580-5586 (1994).

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