1	Supplementary information for						
2	Efficient long-range conduction in cable bacteria through nickel						
3	protein wires						
4	Henricus T. S. Boschker, Perran L.M. Cook, Lubos Polerecky, Raghavendran Thiruvallur						
5	Eachambadi, Helena Lozano, Silvia Hidalgo-Martinez, Dmitry Khalenkow, Valentina						
6	Spampinato, Nathalie Claes, Paromita Kundu, Da Wang, Sara Bals, Karina K. Sand,						
7	Francesca Cavezza, Tom Hauffman, Jesper Tataru Bjerg, Andre G. Skirtach, Kamila Kochan						
8	Merrilyn McKee, Bayden Wood, Diana Bedolla, Alessandra Gianoncelli, Nicole M.J.						
9	Geerlings, Nani Van Gerven, Han Remaut, Jeanine S. Geelhoed, Ruben Millan-Solsona,						
10	Laura Fumagalli, Lars-Peter Nielsen, Alexis Franquet, Jean V. Manca, Gabriel Gomila, Filip						
11	J. R. Meysman						
12							
13	This file includes:						
14	Supplementary Notes						
15	Supplementary Figures 1 to 16						
16	Supplementary Table 1						
17							
18	Other Supplementary information for this manuscript include the following:						
19	Supplementary Movie 1						
20	Excel files with Supplementary Data 1 and 2						
21							

22 Supplementary Notes

23

24 Note 1: Additional results and discussion of the ToF-SIMS analysis

The principal results of the ToF-SIMS analysis are described in the main text. Here we provide additional experimental results and discussion to support the conclusions drawn in the main text.

28 Fragment identification. In total 71 fragment ions were identified in positive mode and 173 in negative mode in the ToF-SIMS runs of the fiber sheaths. Identification was based on 29 the exact mass in ToF-SIMS analysis and was checked for identity and purity with the 30 31 Orbitrap-SIMS runs. For transition metals and S, the isotopic mass distribution was verified. For Ni, the four main isotopes (⁵⁸Ni, ⁶⁰Ni, ⁶¹Ni and ⁶²Ni) were found with the expected 32 isotopic abundance, whereas the minor isotope ⁶⁴Ni showed mass interference from a low 33 amount of Zn. For Cu, both isotopes (⁵³Cu and ⁶⁵Cu) showed the expected abundance. For Fe, 34 the two main Fe isotopes (⁵⁴Fe and ⁵⁶Fe) showed the expected abundance, whereas the minor 35 Fe isotopes (⁵⁷Fe and ⁵⁸Fe) showed mass interference from an unknown source and Ni 36 respectively. Zn counts were low and not reliable as the main isotope $(^{64}$ Zn) showed mass 37 interference from ⁶⁴Ni. Other transition metals found in metalloproteins like Mn, Co, Mo and 38 W were not detected or had very low counts. The main isotopes of Sulfur (³²S and ³⁴S) also 39 showed the expected abundances, whereas the minor isotopes ³³S and ³⁶S showed mass 40 interference from an unknown source. Assignment of the identified fragment ions to 41 biochemical components of intact cable bacteria and the fiber sheath was based on 42 literature^{14–20} and specific depth profiles (see Supplementary Data 1 and 2, Fig. 1, and 43 Supplementary Figure 1). 44

Transition metals. The ToF-SIMS depth profiles of Ni and Fe are discussed in detail in 45 46 the main text. Cu showed a peak at the first data point of the depth profile in both the fiber sheaths and the intact cable bacteria (Fig. 1D, Supplementary Figure 1A-B, Supplementary 47 Figure 11), though with a highly variable count number. This surface peak and its variability 48 between samples was consistent with our LEXRF analysis (Fig. 3), which also showed a 49 highly variable Cu signal. However, Cu counts from STEM-EDX analysis and the majority of 50 the ToF-SIMS depth profiles in the fiber sheaths were low compared Ni counts. This suggests 51 that Cu is not present in high concentrations, but likely derives from an variable impurity that 52 adsorbs to the sample surface at some stage during sample preparation (either cable bacteria 53 collection, fiber sheath extraction procedure) or analysis. 54

Polysaccharide fragments. As discussed in the main text, the basal polysaccharide 55 layer (Fig. 1D and 1E) most likely consists of peptidoglycan, which consists of an acidic 56 amino sugar backbone interconnected by short peptides²¹. We therefore expected to find 57 fragments that contain both oxygen and nitrogen, as derived from the amino sugar backbone 58 59 of peptidoglycan. To verify that we can detect peptidoglycan specific fragments among these O- and N rich fragments, we performed ToF-SIMS analysis of four protein and 60 polysaccharide reference samples: Bovine Serum Albumin (BSA), starch, pectin, and 61 62 peptidoglycan (Supplementary Figure 8). Although the mass spectra of these reference samples showed major differences in both negative and positive mode (Supplementary Figure 63 8), we could not detect any peptidoglycan specific fragments in the fiber sheaths. Moreover, 64 fragments that contained both nitrogen and oxygen showed either the same depth profile as 65 amino acid (protein) derived fragments (i.e. an initial peak at the start of the depth profile), or 66 were specific for nucleic acids that showed low counts and variable depth profiles 67 (Supplementary Data 1 and 2). The latter suggests that some residual DNA/RNA remains 68 69 associated with the fiber sheaths after extraction. The lack of peptidoglycan specific

fragments is most likely due to the abundance of protein in the fiber sheath, which leads to
high signals of amino acid derived fragments that mask fragments derived from
peptidoglycan. Further work is needed to confirm that the polysaccharide layer is indeed
made of peptidoglycan.

P-containing fragments. The fiber sheaths also showed high counts of various P-containing fragment ions in both positive and negative mode (Supplementary Data 1 and 2)with variable depth profiles between runs. The likely source of these P-containing ions is thepoly-phosphate granules that are commonly found in cable bacteria²². Likely, the poly-phosphate granules are incompletely removed during the fiber sheath extraction procedure,and they comprise the particles that are seen in the interior of fiber sheaths during STEM 3D-tomography (Fig. 1A).

81 S and Ni containing fragments in negative mode. Negative mode ToF-SIMS depth profiles showed a subsurface peak of various S-derived anions $({}^{32}S^-, {}^{34}S^-, SH^- and S_2^-)$ at 82 approximately the same position as the Ni peak in positive mode (Fig. 1D and 1E, 83 84 Supplementary Figure 1, Supplementary Data 2). In addition, a substantial number of Nicontaining fragments were found in negative mode that also showed this distinct subsurface 85 peak (Fig. 1D, Supplementary Figure 1, Supplementary Data 2). The most prominent Ni 86 containing fragments were a series of Ni_xS_y⁻ cluster ions (x = 1 to 6, y = x +/- 1, both ⁵⁸Ni and 87 ⁶⁰Ni isotopologues were detected, Supplementary Figure 9). Additionally, we found $NiC_xN_x^-$ 88 (x = 1 or 2) and NiCSN⁻, and even a small but detectable Ni⁻ signal (Supplementary Data 2). 89 All these Ni-containing fragments showed a similar sharp subsurface peak in the depth profile 90 as found for Ni in positive mode (Fig. 1E). The observed NiS clusters are most likely formed 91 in the ion plume of the ToF-SIMS rather than being natively present in the fiber sheath^{23–25}. 92 As a control, we analyzed an artificial mixture of protein (BSA) and freshly precipitated 93 mineral NiS, and we indeed observed the same Ni_xS_v cluster ions (Supplementary Figure 94

95 9D). Still, due to an effect called self-focusing²³, the detection of NiS clusters suggests that Ni
and S must be present in close proximity in the fiber sheath (lateral (XY) within <0.5 nm,
97 depth (Z) within <2 nm), otherwise they would not be formed in the ToF-SIMS ion plume.
98 This observation is hence in agreement with the presence of a S-ligated Ni cofactor in the
99 fiber sheaths.

We also found a number of organic S fragments in negative mode ToF-SIMS spectra 100 101 (Fig. 1E, Supplementary Figure 1, Supplementary Data 2). Most of these fragments showed a similar depth profile as S⁻ with a sharp subsurface peak on top of a general S signal, and so 102 103 again, we cannot exclude that they were formed during ToF-SIMS analysis. However, two 104 fragments ($C_2S_2^-$ and $C_2S_2H^-$) showed specific depth profiles, with the subsurface peak much more pronounced than for S⁻ (Fig. 1E). This suggests that these two fragments originate 105 directly from the fiber sheath. These fragments could indicate that the fiber proteins are rich 106 in disulfide bridges (C-S-S-C) or they could be derived from the Ni/S group. Fiber sheaths 107 display an exceptional chemical resistance, as they remain their integrity and conductivity 108 109 after SDS and EDTA treatments, and this resistance could be aided by protein disulfide bridges. 110

Sputtering depth calibration. Sputtering depth was calibrated by in-situ AFM height 111 112 measurements at three different times (Supplementary Figure 2): before sputtering, just after the Ni or S peak appeared, and finally just after the carbohydrate peak appeared. Sputtering 113 114 depth was approximately linear with sputtering time, although the initial protein layer seemed to sputter somewhat faster (Supplementary Figure 2). A representative example of the 115 resulting AFM height maps is shown in Supplementary Figure 10, where sputtering was 116 117 stopped at 57 sec (i.e. just after the Ni peak had emerged). There was considerable lateral variation in sputtering depths for a given sputtering time as seen in transects across the central 118 cell area (Supplementary Figure 10D) and at the cell junctions (Supplementary Figure 10C). 119

As the junctions are higher than the cell areas (e.g. Supplementary Figure 10A), the shading in sputtering seen at the cell junctions probably due the angled Ar-cluster sputter beam hitting the side of the cell junction facing the beam at a higher angle whereas the opposite side is hit much less. This lateral variation in sputtering rate together with the left-over cytoplasm content and the presence of the cart-wheel structure at the junctions leads to smearing of the ToF-SIMS depth profiles. This explains why only the first fiber sheath layer is clearly seen in the ToF-SIMS depth profiles.

127

128 Note 2: Additional results and discussion of the SDM analysis

The key results of the Scanning Dielectric Microscopy (SDM) analysis are described in the 129 main text. Here we describe additional results that strengthen the conclusions obtained. The 130 finite element model we used is the same as presented earlier 5-7 except that here, fibers have a 131 core-shell structure with a conductive core and an insulating outer shell. The whole simulation 132 domain (which encompasses the ellipsoid fiber and the surrounding space) is cylindrical with 133 height 34 µm and radius 17 µm. Insulating boundary conditions are assumed on the lateral 134 and top borders of the simulation domain. A cross-section of the electric potential distribution 135 obtained from the model calculations is shown in the right insert in Fig. 5C. At the frequency 136 of the calculations (2 kHz) and for the parameters in Fig. 5, the electric potential is real, and 137 hence, the phase of the electric potential is constant in space. The length of the fiber in the 138 model is $L = 1 \mu m$. It has been shown previously that for L > 100 nm the force acting on the tip 139 is independent from the fiber length⁶. 140

Supplementary Figure 12A shows experimentally derived capacitance gradient crosssection profiles corresponding to SDM images (similar to the one shown in the left insert in
Fig. 5B of the main text) acquired at different tip-substrate distances (continuous lines,

ranging from z=60 nm to z=240 nm). The profile at z=66 nm is shown in Fig. 5C of the main 144 145 text. The dashed lines represent the result of the theoretical calculations with the model described above for the same fiber parameters as those use in Fig. 5 of the main text (h=42 146 nm, w=87 nm, d=12 nm, $\varepsilon_s = \varepsilon_c = 3$, $\sigma_s = 0$ S/cm, $\sigma_c = 20$ S/cm). The tip-substrate distances have 147 been determined from a capacitance gradient approach curve measured on a bare part of the 148 substrate (Supplementary Figure 12B, symbols) following procedures previously reported⁵. 149 150 The calculated capacitance gradient profiles nicely reproduce the experimental ones with no adjustable parameter. At all distances, the theoretical (and experimental) electric force phase 151 152 contrast is zero (not shown). Similar conclusions are reached if we analyze the capacitance gradient approach curves measured on the substrate and on the fiber (pink and black thick 153 lines in Supplementary Figure 12B). A least square fitting of the model curves to the 154 experimental curve for $\varepsilon_s = \varepsilon_c = 3$ and $\sigma_s = 0$ S/cm, $\sigma_c = 20$ S/cm gave d=12±2 nm (red continuous 155 156 line in Supplementary Figure 12B), in agreement with the value obtained from the capacitance gradient cross-section profile analysis. The sensitivity of the capacitance gradient 157 profiles on the thickness of the insulating shell is shown by the dashed lines in Supplementary 158 Figure 12B for the approach curves and in Supplementary Figure 12C for the profiles at z=66 159 nm. In these figures we also compare the predictions corresponding to an homogeneous 160 conductive model, corresponding to d=0 nm (or to $\sigma_s = \sigma_c = 20$ S/cm, dark grey line), and to an 161 homogeneous insulating model, corresponding to d=h/2 (or $\sigma_c=0$ S/m, light grey line). In all 162 cases we assumed a protein composition of all layers, $\varepsilon_s = \varepsilon_c = 3$. The pure conductive model 163 overestimates the force acting on the tip, while the pure insulating model underestimates it. 164

We have considered other possible sets of electric parameters for the fiber to explore alternative interpretations of the SDM results. Figs. S13A and S13B show, respectively, the contrast values of the amplitude (in zF/nm) and phase of the electrical force for the tip at the center of the fiber at a distance z=66 nm from the substrate as a function of the conductivity

of the core in the range from $\sigma_c = 10^{-9}$ S/m (insulator) to 10^3 S/m (conductor). The dielectric 169 constants have been fixed to $\varepsilon_s = \varepsilon_c = 3$ (proteins). The thickness of the insulating layer has been 170 varied from d=0 nm to d=20 nm. The calculations were done for the frequency of the 171 experiments (2 kHz). The amplitude as a function of the core conductivity displays two 172 173 plateaus for low and high conductivities, separated by a transition region, which tends to show a third plateau not fully displayed. The phase shows a minimum for every transition region 174 mentioned above. Such behavior is characteristic of materials with (equivalent) permittivities 175 showing real (dielectric) and imaginary (conductive) parts. An analytical expression for the 176 equivalent homogeneous permittivity of the core-shell cylinder in a non-uniform electric field 177 cannot be derived, but the behavior observed is qualitatively similar to the one corresponding 178 to a core-shell spherical particle in a uniform electric field²⁶. For a given measured 179 capacitance gradient contrast (grey band in Supplementary Figure 13A), one can obtain 180 couples of values for the shell thicknesses and core conductivity that match^{5,27}. For instance, 181 we observe that the solution found above for the shell thickness (d=12 nm) is valid for $\sigma_c > 10^{-2}$ 182 S/m (and, in particular, for $\sigma_c = 2 \cdot 10^3$ S/m as employed here). Other couples of values are for 183 instance d=5 nm and σ_c =5·10⁻⁶ S/cm or d=0 nm, σ_c =7·10⁻⁷ S/cm, among others. For these 184 185 values the same capacitance gradient profile (matching the experimental one) is obtained, as shown in the Supplementary Figure 13C. However, the different couples of values predict 186 187 different electrical phase contrast profiles (see Supplementary Figure 13D). Only the couple of values corresponding to the solution reported in the main text (d=12 nm and $\sigma_c > 10^{-2}$ S/m) 188 predicts a null phase contrast, as seen in the experimental results. 189

We have also considered a homogeneous dielectric model, with no conductivity, and analyzed the (equivalent) homogeneous dielectric constant that the fiber should have to explain the experimental results. A least square fitting of the calculated capacitance gradient approach curves for different dielectric constants to the experimental curve measured on the

fiber gives $\varepsilon_s = \varepsilon_c = 11 \pm 3$ (see Supplementary Figure 14A). For this equivalent dielectric 194 constant value the constant height capacitance gradient profiles at the different tip substrate 195 196 distances adequately reproduce the experimental ones (Supplementary Figure 14B). The phase contrast for this model is zero (not shown). However, a problem is that the equivalent 197 dielectric constant value found ($\varepsilon_s = \varepsilon_c = 11 \pm 3$) is much larger than values obtained, with the 198 same technique, on other (dry) bio-samples made of lipids ($\varepsilon = 2$)²⁸, proteins ($\varepsilon = 3-5$)^{27,28,6,7} and 199 even nucleic acids $(\varepsilon \sim 8.5)^{5,27}$ (see Supplementary Figure 14C). Our composition data all 200 201 suggest that the fibers are made primarily of protein. During ToF-SIMS analysis of fiber sheaths, some fragments derived from residual RNA or DNA were detected, but these were 202 not associated with the initial fiber protein layer (see Additional results and discussion of the 203 ToF-SIMS analysis) and it therefore seems unlikely that an isolated fiber as analyzed here by 204 SDM would contain major amounts of DNA. For this reason, the pure dielectric model is 205 discarded in favor of the conductive core-shell model. 206

We have also analyzed the homogeneity of the fiber electrical properties along its 207 longitudinal direction in view of the slight variations of the electric contrast observed in the 208 SDM images along the fiber (see left insert in Fig. 5C of the main text). To this end we 209 considered the core-shell geometrical model that accounts for the observed (tiny) height 210 variations of the fiber height $(\pm 5 \text{ nm})^{29}$ shown in Supplementary Figure 15A. Examples of 211 calculated longitudinal and transversal electric potential distributions are shown in figs. S15B 212 and S15C, respectively. The tip and electrical parameters of the fiber are the same as those 213 214 used in the calculations of in Fig. 5 of the main text, except for the shell thickness, which here is $d=10\pm 2$ nm. The slightly smaller value obtained, as compared to $d=12\pm 2$ nm for the 215 cylindrical model, is due to the tip convolution effects included in this geometrical model (see 216 Supplementary Figure 15D), as discussed elsewhere²⁹. Supplementary Figure 15E shows a 217 calculated constant height SDM image at z=66 nm corresponding to the region enclosed by 218

the dashed rectangle in the insert (which corresponds to the image in the left insert in Fig. 5C
of the main text). The calculated SDM image is nearly identical to the experimental one, a
fact that is further evidenced by comparing the capacitance gradient cross-section profiles
calculated on the hills and valleys of the image with the experimental ones (figs. S15F and
S15G, respectively). This result shows that, to a good approximation, the electrical properties
of the fiber are homogeneous along its length.

225 Finally, we have also analyzed the electric properties of the fibers when they are still embedded within the fiber sheath by SDM. Supplementary Figure 16A shows an AFM 226 227 topographic image of a fiber sheath. Figs. S16B and S16C show, respectively, topographic 228 and constant height SDM images acquired in the area enclosed by the dashed rectangle in Supplementary Figure 16A. The corresponding height and capacitance gradient profiles along 229 the dashed lines in figs. S16B and S16C are shown in figs. S16D and S16E, respectively 230 (black lines). The geometry of the tip has been calibrated with a capacitance gradient 231 approach curve acquired on the substrate (not shown) giving R= 26 ± 2 nm, $\theta=22\pm 3^{\circ}$, 232 C'_{offset}=109±3 zF/nm. The sample geometry has been reconstructed by using a topographic 233 reconstruction procedure²⁹, as above (see Supplementary Figure 16D). For simplicity, we 234 have considered an equivalent homogeneous dielectric model characterized by an (equivalent) 235 dielectric constant ε_{sheath} . An example of a calculated electric potential distribution is shown in 236 Supplementary Figure 16F for a tip-substrate distance z=155 nm and $\varepsilon_{sheath}=7$. Supplementary 237 Figure 16E shows calculated constant height capacitance gradient profiles along the dashed 238 line in Supplementary Figure 16C for different values of ε_{sheath} (thin lines). The experimental 239 values on the hills (corresponding most likely to fiber positions) agree with the calculated 240 241 ones for $\varepsilon_{\text{sheath}} \sim 7-11$ (except the first hill that gives a somewhat larger value). These values are in reasonable agreement with those found for the isolated fiber when the equivalent 242 dielectric model is considered $\varepsilon \sim 11 \pm 3$. This result implies that the electrical properties 243

obtained for the fibers when within the fiber sheath are consistent with those obtained on 244 245 isolated fibers. Similar results are expected for a fiber sheath model that included its internal structure, composition and electrical properties (e.g. conductive core-shell model for the 246 fibers), as the one sketched in Fig. 5A of the main text. However, obtaining quantitative 247 predictions from SDM measurements for such complex model involving buried structures lies 248 outside the current capabilities of SDM³⁰ since too many unknowns are present in the model 249 (e.g. fiber position within the fiber sheath, number of fibers present within the polysaccharide 250 layer, thickness of the polysaccharide layer, etc.). 251

252

Note 3: Calculations for estimating the contribution of the fibers to the fibers sheath, the distance between Ni atoms in the conductive fibers and electron tunneling times.

Electron hopping rates between redox sites in metalloproteins decreases logarithmic 255 with the distance between the redox sites³¹. Given the relatively low Ni content of the fiber 256 sheaths of 0.017 and 0.034 Atm% as determined by STEM-EDX for two sample batches 257 (Supplementary Table 1), we tried to estimate what the inter atomic distance of the Ni atoms 258 would be. At the cell area thick fiber sheaths are ca. 4 µm wide and 117 nm high (this study) 259 and contain on average 60 parallel fibers² that each contain a conductive core of 26 nm 260 diameter (this study). From this we calculate that the conductive fiber cores explain only 7% 261 of the cross-section area of the fiber sheath. If we do the same calculations for the total 50 nm 262 263 fiber with that includes the non-conductive outer shell, we arrive at a 25% contribution from the fibers to the fiber sheath. Assuming that the conductive fiber core is made of protein with 264 a density of 1.4 g·cm⁻³ (density from³²) and that Ni is concentrated in the conductive core (this 265 study), we estimate that the Ni concentration in the conductive fiber core is between 0.24 and 266 0.54 Atm% or 260 and 580 µMol·cm⁻³, which are substantial concentrations. Further 267

assuming that Ni is homogenously distributed in the core of the fibers, this Ni concentration in the fiber core leads to an interatomic Ni distance of 1.4 - 1.9 nm, which is within the range as found for metalloproteins involved in electron transfer³¹.

We then used these estimates of interatomic Ni distances to calculate electron 271 tunneling times needed to support the high average normalized currents in fiber sheaths of 28 272 273 nA (calculated from the standard SDS+EDTA extraction data in Fig. 4D). Estimated tunneling times are 50 - 90 nsec, which is possible but on average however approximately a 274 factor 10 faster than expected for metalloproteins with similar hopping distances³¹. It seems 275 therefore unlikely that long distance electron redox hopping through the Ni/S-group can on its 276 own support the high conductivity of the fibers in cable bacteria (see Discussion in the main 277 text). 278

279

282



283

Supplementary Figure 1. Principal Component Analysis (PCA) of ToF-SIMS depth profiles 284 for a selection of ion fragments recorded from fiber sheaths of cable bacteria. ToF-SIMS data 285 were collected in both positive (A, B) and negative modes (C, D), and the results of three 286 replicate depth profiling runs are shown. Counts for fragments were normalized and centered. 287 Panels (A, C) show the scores plot (labels indicate sputtering time in seconds; colors denote 288 separate runs). Panels (B, D) show the loading plots of the fragments. Only the data until the 289 peak carbohydrate-derived fragments at 125-164 seconds of sputtering time were retained 290 (representing the first half of the double-folded fiber sheath). 291



Supplementary Figure 2. Calibration of sputtering depth as a function of sputter time during 293 ToF-SIMS analysis. The in-situ AFM within the ToF-SIMS instrument was used to record 294 height images at three different times (start (sputter time = 0 sec), just after the Ni or S peak 295 (ca. 50 sec) and just after the carbohydrate peak (ca. 160 and 190 sec)). Sputtering depth 296 297 (black diamonds, average +/- SD) showed a linear relation with sputtering time (dashed line = regression line through origin), though the initial protein layer seemed to sputter somewhat 298 299 faster. The total thickness of the double-folded fiber sheath amounts to 117 ± 10 nm 300 (determined within the middle of cells at t = 0 sec; black dotted line) and is well in agreement with previous independent AFM imaging ². The average position of the Ni/S peak (15 ± 3 nm) 301 and the carbohydrate peak (59 ± 6 nm) are indicated (red and green dotted lines respectively, 302 303 N = 6 from 3 positive and 3 negative depth profiles in Supplementary Figure 1).



Supplementary Figure 3. (A) Green-laser Raman spectra of intact cable bacteria from 306 additional sediments also show the characteristic low-frequency bands at 371 and 492 cm⁻¹. 307 Spectra are from ca. 1 µm wide freshwater cable bacteria (Aarhus pond sediment, intact living 308 filaments in glass micro-chambers, average background-corrected spectrum) and ca. 1 µm 309 wide brackish cable bacteria (Yarra River sediment, intact filaments air-dried on a CaF₂ 310 cover, average raw spectrum). (B) Interpretation of the Raman signal distribution in the scan 311 across an individual cable bacterium in Fig. 2B. General components of a bacterial cell such 312 as proteins and CH are more-or-less evenly distributed in the cross section (green filled 313 circle), whereas the Ni/S group is only found in the periplasmic space (red open circle). In the 314 depth integrated Raman signals, this leads to a unimodal distribution for Raman bands from 315 the general components and a bimodal distribution for the two low frequency Ni/S bands with 316 a peak at both edges of the filament. 317



320

Supplementary Figure 4. Effect of ¹³C labelling on the green-laser Raman spectra from 321 intact cable bacteria. A) The low-frequency bands either do not shift (371 cm⁻¹ band) or are 322 only slightly affected (491 cm⁻¹ band shifts to 489 cm⁻¹) by the labelling. B) The characteristic 323 cytochrome bands clearly shift to lower values suggesting that cable bacteria filaments were 324 highly labelled with ¹³C. The average of N = 32 spectra is shown for both unlabeled control 325 and ¹³C treatments and shifts in wavenumbers are indicated. C) PCA analysis of the 350 to 326 550 cm⁻¹ region of the spectra containing the low-frequency bands shows a difference 327 between the labelled versus unlabeled treatments. D) Variable (wavenumber) loadings on the 328 first PCA axis show that only the second lower band at 492 cm⁻¹ shifted slightly to lower 329 330 values in the labelled spectra.



Supplementary Figure 5. NanoSIMS images of A) negative ions and B) positive ions of metals from single fiber sheaths deposited on a gold-coated polycarbonate filter. In A) only the first 100 planes of analysis were selected as this most clearly showed the fibers structure in the ³²S image. In B) the first 50 planes were selected as this showed the fibers most clearly in the ⁵⁸Ni image. Independent replica's (N = 2) yielded similar results.



Supplementary Figure 6. Atomic Force Microscopy (AFM) images of fiber sheaths as
extracted with the standard protocol (A and B, SDS+1 mM EDTA for 10 min) and the high
EDTA treatment to remove Ni (C and D, SDS+50 mM EDTA for 10 min). Shown are the
height (A,C) and peak force error (B,D) data. The parallel fiber structures are retained after
high-concentration EDTA extraction. Independent replica's (N = 2) yielded similar results.







350 sheath showing the parallel fibers.





Supplementary Figure 8. ToF-SIMS positive (A) and negative (B) mode mass spectrum of
additional protein and polysaccharide reference samples analyzed to aid in the interpretation
of the fiber sheath mass spectra.







Supplementary Figure 9. ToF-SIMS negative mode Ni and S cluster ions from fiber sheaths 359 360 and a NiS/BSA mixture. A) ToF-SIMS negative mode mass spectrum of a mixture of NiS mineral in BSA. Positions of a selection of important Ni and S cluster ions are indicated. B) 361 Distribution of Ni₃S₃⁻ isotopomers as a seen in the spectrum of fiber sheaths and C) compared 362 to the expected distribution showing the good fit (similar fits were found for the other Ni_xS_y-363 clusters described, data not shown). D) Distribution of Ni_xS_y-cluster ions in fiber sheaths and 364 365 NiS/BSA mixture (counts were normalized to Ni₃S₃⁻). Cluster ions with 3 or less Ni atoms were sufficiently mass separated from fragments containing 2 oxygens instead of sulfur such 366 as Ni₃S₂O₂⁻. 367



Supplementary Figure 10. Example of *in-situ* AFM calibration of ToF-SIMS sputtering 369 depth for fiber sheaths. AFM height map. A) mapping before ToF-SIMS analysis was started. 370 B) mapping after 57 sec of sputtering time (i.e. shortly after the peak in Ni⁺ counts). C) The 371 differential image between (A) and (B) shows the amount of material removed after 57 sec 372 373 sputtering. Notice shading of the junctions with higher than average sputtering rates on the side and lower than average sputtering rates on the opposite side. D) Height profiles as 374 indicated by the blue lines in panel (C) showing the lateral variation in sputtering depth. 375 Actual sputtering depths were calculated from height profiles as the average difference in 376 height of the fiber sheath (relative to the wafer surface) before and after sputtering (in this 377 example 22.8 ± 3.2 nm). 378





381 Supplementary Figure 11. ToF-SIMS depth profiles of intact cable bacteria demonstrate that the Ni/S group is located in the periplasmic space. Representative ToF-SIMS depth profiles 382 from intact cable bacteria are shown in A) positive mode and B) negative mode. A selection 383 of fragments is shown representing the major biomolecule classes, and counts of individual 384 fragments were scaled to fit all data in a plot. Duplicate runs in both positive and negative 385 modes showed similar profiles. The counts from Ni₃S₃⁻ are the sum of all ⁵⁸Ni and ⁶⁰Ni 386 isotopologues. The fatty acid fragment $(C_{16}H_{29}O_2)$ clearly shows two peaks with the first 387 peak at the surface coming most likely from the outer membrane and the broader secondary 388 389 peak from the cell membrane. Phosphate (PO_3) also shows a first peak at the surface probably coming from phospholipids and the secondary rise from phospholipids in the cell membrane 390 and poly-phosphate and nucleic acids in the cytoplasm. Arrows indicate sputtering depths as 391 392 determined with the in-situ AFM in the ToF-SIMS. Average sputtering rate was 0.52 ± 0.05 nm/sec, which places the Ni and S containing peak at approximately 30-40 nm depth in 393 agreement with the expected position of the fibers in intact cable bacteria². 394



Supplementary Figure 12. Additional information on the SDM quantitative analysis shown 397 in Fig. 5 of the main text. A) Capacitance gradient cross-section profiles measured 398 (continuous lines) and calculated (dashed lines) at different tip-substrate distances. The 399 profiles for z=66 nm are those displayed in Fig. 5 of the main text. The parameters used in the 400 calculations are the same as those of Fig. 5 of the main text: H=12.5 μ m, W= 3 μ m, L=3 μ m, 401 $l=1 \mu m, \theta=22, R=54 nm, z=60 nm, h=42 nm, w=87 nm, d=12 nm, l=1 \mu m, \epsilon_s = \epsilon_c = 3, \sigma_s = 0$ 402 S/m and $\sigma_c=2.10^3$ S/m. B) Capacitance gradient approach curves measured on a bare part of 403 the substrate (pink continuous line) and on the fiber (black continuous line). The continuous 404 405 dark blue and red lines correspond to the curves that best fit the experimental results from where the tip geometry (R=54 \pm 1 nm, θ =22° \pm 0.5°, C'_{off}=112.5 \pm 1.5 zF/nm) and shell thickness 406 $(d=12\pm 2 \text{ nm})$ have been extracted. The dashed lines represent calculated capacitance gradient 407 approach curves on the fiber for other values of the shell thickness. The continuous dark and 408 409 grey lines correspond to the homogeneous conductive and dielectric models, respectively. The green and orange symbols correspond to the values measured on the bare substrate and fiber, 410 respectively, on the SDM images at different heights. These values have been used to set the 411 tip-substrate distances of the SDM images. C) Effect of the shell thickness on the calculated 412 capacitance gradient profiles for z=66 nm. The meaning of the lines is the same as in B). 413



415

Supplementary Figure 13. Additional analysis of the core-shell conductive model used to 416 interpret the SDM measurements shown in Fig. 5 of the main text. A) Amplitude (in 417 capacitance gradient) and B) phase of the electric force contrast calculated for the cylindrical 418 419 fiber model as a function of the conductivity of the core for different values of the shell thickness. The grey band corresponds to the experimental values extracted from the profiles 420 shown in Fig. 5 of the main text. Calculated capacitance gradient C) and phase D) profiles for 421 422 the couples of shell thickness-core conductivity that match the experimental measured capacitance gradient contrast in A). The thick lines represent the experimental results (same 423 as in Fig. 5 of the main text). The parameters used in the calculations are the same as those of 424 Fig. 5 of the main text and of Supplementary Figure 12, when not otherwise stated. 425



Supplementary Figure 14. Analysis of the SDM measurements shown in Fig. 5 of the main 427 text with a homogeneous cylinder dielectric model ($\varepsilon_s = \varepsilon_c = \varepsilon_{eff}$, $\sigma_s = \sigma_c = 0$ S/m). A) Calculated 428 capacitance gradient approach curves on the fiber for different values of the equivalent 429 dielectric constant ε_{eff} compared to the experimental measured curves on the substrate and 430 fiber (same as in Supplementary Fig 6B, thick lines). The best fit is obtained for $\varepsilon_{eff}=11\pm3$. B) 431 Calculated capacitance gradient profiles for $\varepsilon_{eff}=11$ for different tip-substrate distances, and 432 comparison with the measured profiles (thick lines, same as in Supplementary Figure 6A). C) 433 Calculated capacitance gradient profiles for z=66 nm and different equivalent dielectric 434 constants ε_{eff} , and comparison with the measured profile (thick line, same as in Fig. 5 of the 435 436 main text). The parameters used in the calculations are the same as those of Fig. 5 of the main text and of Supplementary Figure 6, when not otherwise stated. 437



Supplementary Figure 15. Analysis of the SDM measurements shown in Fig. 5 of the main 438 text with a geometrical core-shell model reconstructed from the measured topography. A) 439 Detail of the geometry and mesh of the model generated from the measured topography of the 440 441 fiber (corresponding to the insert in Fig. 5B of the main text). The core-shell structure is defined by assigning a sigmoidal behavior to the dielectric constant and conductivity with 442 plateau representing the shell and core values, respectively. The thickness of the shell 443 444 corresponds to the center of the sigmoid. B) Example of a calculated electric potential distribution along the fiber for z=66 nm, $\varepsilon_s = \varepsilon_c = 3$, $\sigma_s = 0$ S/m, $\sigma_c = 2 \cdot 10^3$ S/m and d=10 nm. The 445 tip parameters are the same as those in Fig. 5 of the main text. C) Example of a cross-section 446 electric potential distribution corresponding to B). D) Comparison of the cross-section of the 447 448 measured topography (red line), the topographically reconstructed fiber model (black continuous line) and of the cylinder model (black dashed line). E) Calculated constant height 449 capacitance gradient SDM image by using the model in A) for the parameters in B). The 450 calculated image corresponds to the area enclosed by the dashed line in the insert. F) and G) 451 Comparison of the calculated capacitance gradient profiles (red lines) with the experimental 452 ones (black lines) on the hills (A,B,C) and valleys (a,b) indicated in E), respectively. 453



Supplementary Figure 16. SDM measurements and analysis on a fiber sheath. A) AFM 455 topographic image of the fiber sheath analyzed. B) and C) AFM topographic and SDM 456 constant height (z=155 nm) images measured on the region enclosed by the dashed rectangle 457 in A). D) and E) Cross-section topographic and capacitance gradient profiles along the dashed 458 lines in B) and C) (thick black lines), respectively. The hills in the image correspond to 459 different fibers. F) Geometrical model reconstructed from the measured topography used in 460 the calculations, with an example of an electric potential distribution overlaid on it. The 461 electrical properties of the fiber sheath have been characterized by an equivalent 462 homogeneous dielectric constant, $\varepsilon_{\text{sheath}}$. The tip geometry has been calibrated from a 463 464 capacitance gradient approach curve on the bare substrate giving R=26±2 nm, θ =22±3°, C'_{offset}=109±3 zF/nm. The theoretically predicted capacitance gradient profiles for this model 465 for different values of $\varepsilon_{\text{sheath}}$ are shown in E) (thin lines). Most of the hills (fibers) correspond 466 to $\varepsilon_{\text{sheath}}=7-11$ in good agreement with the equivalent dielectric constant value found on 467 isolated fibers ($\varepsilon_{eff}=7-11$). For the first hill a higher value is obtained, probably indicating that 468 two fibers are overlaid. 469

- **Supplementary Table 1.** STEM-EDX element compositions for intact cable bacteria and
- 472 fiber sheaths. Data from two separate imaging sessions are shown (dates in first row).

	04-05-17			08-08-17				
	intact cable				intact cable			
	bacteria		fiber sheath		bacteria		fiber sheath	
	N = 3		N = 3		N = 8		N = 7	
Elements								
(atm%)	AVG	SD	AVG	SD	AVG	SD	AVG	SD
Carbon	76.617	2.382	73.013	0.508	82.233	1.231	80.099	1.081
Oxygen	14.370	1.185	15.880	0.322	10.854	1.394	11.319	1.205
Nitrogen	6.133	0.250	9.283	0.224	5.411	0.547	7.333	0.477
Phosphorus	1.120	0.607	0.663	0.100	0.833	0.167	0.736	0.150
Sulphur*	0.510	0.035	0.350	0.225	0.483	0.078	0.329	0.023
Zinc	0.210	0.061	0.103	0.015	0.145	0.035	0.149	0.040
Iron	0.047	0.006	0.023	0.006	0.033	0.007	0.013	0.005
Nickel	0.009	0.001	0.037	0.006	0.009	0.004	0.016	0.005
Copper	0.009	0.004	0.010	0.009	0.006	0.005	0.007	0.005

- 475 *: Sulphur showed interference from a Molybdenum impurity derived most likely from the
- 476 TEM grid and is therefore uncertain.

Supplementary references

- 479
- 1. Burdorf, L. D. W. et al. Long-distance electron transport occurs globally in marine 480 sediments. Biogeosciences 14, 683-701 (2017). 481
- 2. Cornelissen, R. et al. The cell envelope structure of cable bacteria. Front. Microbiol. 9, 482 (2018). 483
- 3. Bjerg, J. T. et al. Long-distance electron transport in individual, living cable bacteria. 484 Proc. Natl. Acad. Sci. 115, 5786-5791 (2018). 485
- 4. Eilers, P. H. C. & Boelens, H. F. M. Baseline correction with asymmetric least squares 486 smoothing. Leiden University Medical Centre Report. (2005). 487
- 5. Fumagalli, L., Esteban-Ferrer, D., Cuervo, A., Carrascosa, J. L. & Gomila, G. Label-free 488 489 identification of single dielectric nanoparticles and viruses with ultraweak polarization forces. Nat. Mater. 11, 808-816 (2012). 490
- 491 6. Lozano, H. et al. Dielectric constant of flagellin proteins measured by scanning dielectric microscopy. Nanoscale 10, 19188-19194 (2018). 492
- 7. Lozano, H., Millán-Solsona, R., Fabregas, R. & Gomila, G. Sizing single nanoscale 493 objects from polarization forces. Sci. Rep. 9, 1-12 (2019). 494
- 8. Wang, D. et al. Interplay between spherical confinement and particle shape on the self-495 assembly of rounded cubes. Nat. Commun. 9, 2228 (2018). 496
- 9. van Aarle, W. et al. The ASTRA toolbox: A platform for advanced algorithm 497
- development in electron tomography. Ultramicroscopy 157, 35-47 (2015). 498
- 10. Gianoncelli, A., Kourousias, G., Merolle, L., Altissimo, M. & Bianco, A. Current status 499
- of the TwinMic beamline at Elettra: A soft X-ray transmission and emission microscopy 500
- station. J. Synchrotron Radiat. 23, 1526–1537 (2016). 501

502	11. Solé, V. A., Papillon, E., Cotte, M., Walter, Ph. & Susini, J. A multiplatform code for the
503	analysis of energy-dispersive X-ray fluorescence spectra. Spectrochim. Acta Part B At.
504	<i>Spectrosc.</i> 62 , 63–68 (2007).
505	12. Polerecky, L. et al. Look@NanoSIMS - a tool for the analysis of nanoSIMS data in
506	environmental microbiology. Environ. Microbiol. 14, 1009-1023 (2012).
507	13. Meysman, F. J. R. et al. A highly conductive fibre network enables centimetre-scale
508	electron transport in multicellular cable bacteria. Nat. Commun. 10, 1–8 (2019).
509	14. Baugh, L. et al. Probing the orientation of surface-immobilized protein G B1 using ToF-
510	SIMS, sum frequency generation, and NEXAFS spectroscopy. Langmuir 26, 16434-
511	16441 (2010).
512	15. Goacher, R. E., Jeremic, D. & Master, E. R. Expanding the library of secondary ions that
513	distinguish lignin and polysaccharides in time-of-flight secondary ion mass spectrometry
514	analysis of wood. Anal. Chem. 83, 804-812 (2011).
515	16. Wei, W. et al. Characterization of syntrophic Geobacter communities using ToF-SIMS.
516	<i>Biointerphases</i> 12 , 05G601 (2017).
517	17. Lebec, V., Boujday, S., Poleunis, C., Pradier, CM. & Delcorte, A. Time-of-flight
518	secondary ion mass spectrometry investigation of the orientation of adsorbed antibodies
519	on SAMs correlated to biorecognition tests. J. Phys. Chem. C 118, 2085–2092 (2014).
520	18. Ding, Y. et al. In Situ Molecular Imaging of the Biofilm and Its Matrix. Anal. Chem. 88,
521	11244–11252 (2016).
522	19. Lee, CY., Harbers, G. M., Grainger, D. W., Gamble, L. J. & Castner, D. G.
523	Fluorescence, XPS, and TOF-SIMS Surface Chemical State Image Analysis of DNA
524	Microarrays. J. Am. Chem. Soc. 129, 9429–9438 (2007).

- 525 20. May, C. J., Canavan, H. E. & Castner, D. G. Quantitative X-ray Photoelectron
- 526 Spectroscopy and Time-of-Flight Secondary Ion Mass Spectrometry Characterization of
 527 the Components in DNA. *Anal. Chem.* 76, 1114–1122 (2004).
- 528 21. Lovering, A. L., Safadi, S. S. & Strynadka, N. C. J. Structural perspective of
- 529 peptidoglycan biosynthesis and assembly. *Annu. Rev. Biochem.* **81**, 451–478 (2012).
- 530 22. Geerlings, N. M. J., Zetsche, E.-M., Hidalgo Martinez, S., Middelburg, J. J. & Meysman,
- 531 F. J. R. Mineral formation induced by cable bacteria performing long-distance electron
- transport in marine sediments. *Biogeosciences Discuss*. 1–35 (2018) doi:10.5194/bg-
- **533** 2018-444.
- 534 23. Franquet, A. *et al.* Self focusing SIMS: Probing thin film composition in very confined
 535 volumes. *Appl. Surf. Sci.* 365, 143–152 (2016).
- 536 24. Feld, H., Rading, D., Leute, A. & Benninghoven, A. Comparative investigations of the
 537 secondary ion emission of metal complexes under MeV and keV ion bombardment. *Org.*
- 538 *Mass Spectrom.* **28**, 841–852 (1993).
- 539 25. El Nakat, J. H., Dance, I. G., Fisher, K. J., Rice, D. & Willett, G. D. Laser-ablation
- 540 FTICR (Fourier-transform ICR) mass spectrometry of metal sulfides: gaseous anionic
- 541 nickel-sulfur [NixSy] clusters. J. Am. Chem. Soc. 113, 5141–5148 (1991).
- 542 26. Jones, T. B. *Electromechanics of particles*. (Cambridge University Press, 1995).
- 543 27. Cuervo, A. *et al.* Direct measurement of the dielectric polarization properties of DNA.
- 544 *Proc. Natl. Acad. Sci.* **111**, E3624–E3630 (2014).
- 545 28. Dols-Perez, A., Gramse, G., Calò, A., Gomila, G. & Fumagalli, L. Nanoscale electric
- 546 polarizability of ultrathin biolayers on insulating substrates by electrostatic force
- 547 microscopy. *Nanoscale* **7**, 18327–18336 (2015).
- 548 29. Checa, M. et al. Mapping the dielectric constant of a single bacterial cell at the nanoscale
- 549 with scanning dielectric force volume microscopy. *Nanoscale* **11**, 20809–20819 (2019).

- 550 30. Fabregas, R. & Gomila, G. Dielectric nanotomography based on electrostatic force
- 551 microscopy: A numerical analysis. J. Appl. Phys. 127, 024301 (2020).
- 552 31. Gray, H. B. & Winkler, J. R. Electron flow through metalloproteins. *Biochim. Biophys.*
- 553 *Acta BBA Bioenerg.* **1797**, 1563–1572 (2010).
- 554 32. Fischer, H., Polikarpov, I. & Craievich, A. F. Average protein density is a molecular-
- 555 weight-dependent function. *Protein Sci.* **13**, 2825–2828 (2004).