Supplementary Information

Manipulation of Hot Carrier Cooling Dynamics in Two-dimensional Dion-Jacobson Hybrid Perovskites via Rashba Band Splitting

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Supplementary Note 1: Minimizing the Auger heating effect

Hot-carrier re-excitation via interband Auger recombination of carriers at band edges (also called Auger heating) can induce slow hot-carrier cooling in the tens to hundreds of picosecond time range, particularly at high pump fluence.¹ In our work, we kept the pump fluence low, in the range of 0.8–4.0 μ J/cm², in order to minimize the Auger heating effect.

Supplementary Note 2: Carrier density calculations

The carrier density is estimated via $n_0 = \frac{Pa}{\hbar \omega}$, where *P* is the pump fluence, *α* is the absorption coefficient of the film, and ℏ*ω* is the photon energy. The estimated carrier densities for (3AMP)PbI4 and (4AMP)PbI4 are as follows:

Supplementary Note 3: Hot-carrier temperatures (*T***c)**

*T*c values were extracted by fitting the high-energy tails of the TA spectra using the Maxwell-Boltzmann function. When the difference between the carrier energy and the Fermi level is larger than *k*BT (*i.e.*, *E*−*E*f > *k*BT), the Fermi-Dirac distribution can be approximately described by the Maxwell-Boltzmann distribution: $\frac{1}{\sqrt{1-\frac{1}{n}}}$ $\frac{E-E_f}{1+exp\frac{E-F_f}{1+T}}$ $k_B T$ $\approx exp \frac{E-E_f}{k_B T}$, which is a generally accepted practice to

extract the T_c value. More details can be found in refs. $1, 2, 3$.

Supplementary Note 4: Electron-phonon coupling calculations

The electron-phonon coupling is described by the dimensionless Fröhlich parameter *α*, defined as: $\alpha = \frac{1}{2}$ $4\pi\varepsilon_0$ $\frac{e^2}{\hbar} \left(\frac{1}{\varepsilon_{\infty}} - \frac{1}{\varepsilon_{\rm s}} \right) \sqrt{\frac{m}{2\hbar \omega}}$, where ε_0 is the dielectric permittivity of vacuum; ε_{∞} and $\varepsilon_{\rm s}$ are the optical and static dielectric constants, respectively; *e* is the charge of the carrier; 2πℏ is Planck's constant; *m* is the bare electron effective mass; and ω is the characteristic angular frequency of the longitudinal optical (LO) phonon mode. *ε*s was calculated using the linear response method as implemented in the Quantum Espresso package.⁴ The dielectric function at far-infrared frequencies was obtained using the following equation: $\varepsilon(\omega) = 1 + \sum_{n=1, Nph} \frac{F_n^2}{\omega^2}$ $n=1, Nph \frac{F_n^2}{\omega_n^2 - (\omega + i\eta)^2} + \frac{F_{el}^2}{\omega_{el}^2 - (\omega + i\eta)^2}$ $\frac{r_{el}}{\omega_{el}^2 - (\omega + i \eta_{el})^2}$, where *n* is the index of the N_{ph} phonon modes; ω_n is the *n*th phonon frequency; ω_{el} is the electronic band gap; F_n and F_{el} are the oscillator strengths; and η and η_{el} are the broadening factors. The details of the calculations of the far-infrared dielectric function can be found in refs ^{5, 6}.

Supplementary Note 5. Dielectric constant of organic cation layers

The dielectric constant (ε) of PEA is 3.3.⁷ Although the experimental dielectric constants of 3AMP and 4AMP are not available, their values can be expected to be higher than that of PEA as the additional functional group (−CH2−NH3 +) further increases the dielectric constant of piperidine (*ε* $= 5.9$).

Supplementary Figure 1. Optical characterizations of 2D DJ perovskites. (**a**) X-ray diffraction (XRD) patterns and (**b**) absorption spectra for (3AMP)PbI4 and (4AMP)PbI4 films (the Tauc plots of the absorption spectra are shown in the insets).

Supplementary Figure 2. Atomic force microscopy (AFM) images of (3AMP)PbI4 and (4AMP)PbI4 films.

Supplementary Figure 3. Molecular structures and electrostatic potential surfaces of organic spacers phenmethylammonium (PMA) and phenethylammonium (PEA) in the 2D Ruddlesden-Popper perovskites, and 3-(aminomethyl)piperidinium (3AMP) and 4-(aminomethyl)piperidinium (4AMP) in the 2D Dion-Jacobson perovskites.

Supplementary Figure 4. Crystal structures for 2D RP perovskites, (PMA)2PbI4 and (PEA)2PbI4, together with the chemical structures of organic cations and the distances between inorganic layers.

Supplementary Figure 5. Electronic band structures for 2D DJ perovskites (3AMP)2PbI4 and (4AMP)2PbI4, calculated at the GGA/PBE level of theory without and with spin-orbit coupling (SOC).

Supplementary Figure 6. Electronic band structures for 2D RP perovskites (PMA)2PbI4 and (PEA)2PbI4, calculated at the GGA/PBE+SOC level of theory without and with spin-orbit coupling (SOC).

Supplementary Table 2. Excitation wavelengths (energies) used for the TA, CTA, and timeresolved PL measurements on (3AMP)PbI4, (4AMP)PbI4, (PMA)2PbI4, and (PEA)2PbI4 films.

Film	GSB Position	High-Energy Excitation	Low-Energy Excitation	
(3AMP)PbI ₄	554 nm (2.24 eV)	330 nm (3.76 eV)	500 nm (2.48 eV)	
(4AMP)PbI ₄	517 nm (2.40 eV)	316 nm (3.92 eV)	470 nm (2.64 eV)	
$(PMA)_2PbI_4$	520 nm (2.39 eV)	318 nm (3.90 eV)	472 nm (2.63 eV)	
(PEA) ₂ PbI ₄	519 nm (2.39 eV)	318 nm (3.90 eV)	472 nm (2.63 eV)	

Supplementary Figure 7. Time-resolved PL decays for (**a**) the emission at 565 nm of the (PMA)2PbI4 film and (**b**) the emission at 560 nm of the (PEA)2PbI4 film at excitation wavelengths of 318 and 472 nm.

Supplementary Figure 8. Normalized transient absorption spectra at the excitation wavelength of (**a**) 330 nm for the (3AMP)PbI4 film and (**b**) 316 nm for the (4AMP)PbI4 film with pump fluence of 0.8 μ J/cm².

Supplementary Figure 9. Normalized transient absorption spectra at the excitation wavelength of 318 nm for (**a, b, c**) the (PMA)2PbI4 film and (**d, e, f**) the (PEA)2PbI4 film, with pump fluences of 0.8, 2.0, and 4.0 μ J/cm², respectively.

Supplementary Table 3. Fitting parameters of hot-carrier cooling temperature (T_c) decays in $(3AMP)PbI₄$ and $(4AMP)PbI₄$. The time constants for T_c were obtained via biexponential fittings of the initial T_c decay using the equation: $\tau = a_1 \cdot \tau_1 + a_2 \cdot \tau_2$.

Supplementary Figure 10. Transient absorption spectra at the excitation wavelength of 318 nm for the $(PMA)_2PbI_4$ film with pump fluence of 4.0 $\mu J/cm^2$. A new photobleaching signal appears at 408 nm due to a phase change.

Supplementary Table 4. Fitting parameters of the transient absorption kinetics probed at 554 nm for the (3AMP)PbI4 film (excitation wavelength of 330 nm) and at 518 nm for the (4AMP)PbI4 film (excitation wavelength of 316 nm) with different pump fluences.

Supplementary Table 5. Fitting parameters of the transient absorption kinetics probed at 554 nm for the (3AMP)PbI4 film (excitation wavelength of 500 nm) and at 518 nm for the (4AMP)PbI4 film (excitation wavelength of 470 nm) with different pump fluences.

Supplementary Figure 11. Average nonadiabatic couplings (NACs, in meV) between electronic states for (3AMP)PbI4, (4AMP)PbI4, (PMA)2PbI4, and (PEA)2PbI4, as obtained from nonadiabatic molecular dynamics (NAMD) simulations.

Supplementary Figure 12. Low-frequency vibrational modes. (**a, b**) Experimental Raman spectra (grey lines) and calculated Raman peaks and intensities for (3AMP)PbI4 and (4AMP)PbI4 obtained at LDA level of theory. (**c, d**) Spectral densities obtained by calculating the Fourier transforms of the autocorrelation functions for the fluctuations of the energy differences between two conduction bands in (3AMP)PbI4 and (4AMP)PbI4.

Supplementary Figure 13.Intraband relaxations in 2D RP perovskites. (**a, b**) Electronic bands and energy levels of (PMA)2PbI4 and (PEA)2PbI4 at the high-symmetry Γ-point (the VBM is set to zero energy). (**c, d**) Time evolution of hot electron relaxation starting from the CBM+1 for $(PMA)_2PbI_4$ and from the CBM+2 for $(PEA)_2PbI_4$.

Supplementary Table 6. Average electron/hole effective masses of bare electron bands (*m**), optical and static dielectric constants (ε∞ and εS), angular frequencies of a characteristic LO phonon mode (ω, in THz), and electron-phonon coupling constant (*α*) for (3AMP)PbI4, (4AMP)PbI4, (PMA)2PbI4, and (PEA)2PbI4.

Compound	m^* electron	m^*_{hole}	ϵ_{∞}	ϵ s	ω (THz)	α _{electron-} phonon
(3AMP)PbI ₄	0.182	0.343	4.81	29.86	2.83	2.53
(4AMP)PbI ₄	0.319	0.808	5.89	20.32	2.92	2.29
(PMA) ₂ $PbI4$	0.243	0.697	4.23	22.89	2.64	3.36
(PEA) ₂ $PbI4$	0.323	0.558	3.38	20.85	3.38	4.29

Supplementary Figure 14. Experimental setup used for circular polarized TA measurements.

The circular polarized pump pulse is set at an angle of 45° and another circular polarizer is set either parallel or perpendicular to the pump pulse.

Supplementary Figure 15. Circular polarized TA spectroscopy. Normalized TA spectra of (3AMP)PbI₄ by (a, b) co-circular polarized pump and probe ($\sigma^+\sigma^+$ and $\sigma\sigma^-$) and (c, d) countercircular polarized pump and probe $(\sigma^+\sigma^-$ and $\sigma^-\sigma^+$).

Supplementary Figure 16. Circular polarized TA spectroscopy. Normalized TA spectra of (4AMP)PbI₄ by (a, b) co-circular polarized pump and probe ($\sigma^+\sigma^+$ and $\sigma\sigma^-$) and (c, d) countercircular polarized pump and probe ($\sigma^+ \sigma^-$ and $\sigma^- \sigma^+$).

Supplementary Table 7. Fitting parameters of circular polarized transient absorption kinetics probed at 554 nm for the (3AMP)PbI4 film and at 518 nm for the (4AMP)PbI4 film following excitation at 308 nm with a pump fluence of 4 μ J/cm².

Film	Pump Type	A ₁	τ_1 (p _s)	A ₂	τ_2 (p _s)	A ₃	τ_3 (p _s)
(3AMP)PbI ₄	$\sigma^*\sigma^+$	53%	2.10	32%	34	15%	496
	$\sigma^*\sigma^-$	60%	2.50	26%	35	14%	485
	$\sigma^-\sigma^+$	52%	1.85	31%	30	17%	420
	$\sigma^-\sigma^-$	52%	1.63	31%	29	17%	423
(4AMP)PbI ₄	$\sigma^*\sigma^*$	38%	3.27	38%	62	24%	672
	$\sigma^+\sigma^-$	46%	2.00	33%	30	21%	421
	$\sigma^-\sigma^+$	42%	3.50	38%	72	20%	643
	$\sigma^{-}\sigma^{-}$	50%	3.00	30%	46	20%	550

Supplementary References

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