- Supplementary Information -

Selective ¹³C labelling reveals the electronic structure of flavocoenzyme radicals

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General remarks

Errors of hyperfine principal values and Euler angles have been estimated based on single-point calculations, in which the respective optimized parameter value was varied systematically and the resulting simulated ENDOR spectrum compared to the experimental one. The error margins of Euler angles were chosen rather large for nuclei with hyperfine tensors of small overall and axial anisotropy, such as ¹³C5a, ¹³C10a and ¹³C1'.

Weighting factors were necessary to account for the specific isotope enrichment and relaxation properties of the individual ¹³C nucleus under consideration.

Analysis of ENDOR data from [6,8α-¹³C₂]FAD bound to *Escherichia coli* DNA photolyase

The frozen-solution ENDOR spectra of $[6,8\alpha$ -¹³C₂]FAD-reconstituted *Ec*PL, recorded at magnetic-field values corresponding to the canonical **g**-principal axes orientations, exhibit one pair of intense and narrow signals (resonances "1", see Fig. S1) from which the following components of one hyperfine tensor of nearly axial symmetry could be extracted: $|A_{X,1}| = 3.6$ MHz, $|A_{Y,1}| = 4.2$ MHz and $|A_{Z,1}| = 3.7$ MHz. A much less intense and broader pair of resonances ("2") is also observed that is described with $|A_{X,2}| = 0.5$ MHz, $|A_{Y,2}| = 5.9$ MHz and $|A_{Z,2}| = 8.9$ MHz.



Figure S1. ¹³C W-band (93.88 GHz) Davies-type ENDOR spectra (black curves) from $[6,8\alpha$ -¹³C₂]FADreconstituted *EcPL*, recorded at 80 K. For improved global least-squares fitting, the data have been symmetrized with respect to the Larmor precession frequency at around 35.85 MHz. For this figure, the spectra have also been normalized to equal maximum intensity with respect to each other. The spectral simulations (red curves) have been obtained with the parameters compiled in Table S1.

From $A_{Y,1}$ being the largest component of hyperfine tensor "1", and from the near axial symmetry $|A_{X,1}| \approx |A_{Z,1}|$, which is typical for the hyperfine coupling tensor of a methyl group that is freely rotating about its C–C bond, even at cryogenic temperatures, we conclude that this set of resonances must arise from ¹³C8 α . This is consistent with predictions from DFT, that yield $A_X = -3.0$ MHz, $A_Y = -3.1$ MHz and $A_Z = -2.9$ MHz. For ¹³C6, DFT predicts much larger hyperfine anisotropy: $A_X = +0.3$ MHz, $A_Y = +0.8$ MHz and $A_Z = +11.6$ MHz, which is in fair agreement with the simulations of ENDOR resonances "2". All parameters used to obtain the overall simulations shown as red curves in Fig. S1 are compiled in Table S1.

Table S1. Hyperfine data used to simulate the ENDOR spectra of $[6,8\alpha-{}^{13}C_2]FAD$ -reconstituted *EcPL*. The principal values are accurate within ± 0.2 MHz, and the Euler angles within ± 0.1 rad. The signs of the hyperfine components were taken from DFT calculations. Γ_G is the full width at half maximum (fwhm) of Gaussian resonance line broadening. The w_N are weighting factors that account, *e. g.*, for the specific relaxation properties of the individual ${}^{13}C$ nuclei under consideration.

Parameter	¹³ C6	¹³ C8a	
A_X / MHz	(+)0.5	(-)3.6	
A_Y / MHz	(+)5.9	(-)4.2	
A_Z / MHz	(+)8.9	(-)3.7	
a _{iso} / MHz	(+)5.1	(-)3.8	
$\Gamma_{\rm G}/{ m mT}$	0.07	0.15	
α / rad	+1.5	-0.4	
β / rad	-0.7	+0.8	
γ / rad	-1.8	-0.4	
WN	0.29	0.71	

Analysis of ENDOR data from [5a,8-¹³C₂]FAD bound to *Escherichia coli* DNA photolyase

The frozen-solution ENDOR spectra of $[5a,8-^{13}C_2]$ FAD-reconstituted *Ec*PL, recorded at magnetic-field values corresponding to the canonical **g**-principal axes orientations, exhibit one pair of intense and relatively narrow signals (resonances "3", see Fig. S2) from which the following components of one hyperfine tensor can be extracted: $|A_{X,3}| = 15.1$ MHz, $|A_{Y,3}| = 13.0$ MHz and $|A_{Z,3}| = 11.6$ MHz. A much weaker and broader pair of resonances ("4") with almost axial symmetry is also observed that is described with $|A_{X,4}| = 0.8$ MHz, $|A_{Y,4}| = 1.8$ MHz and $|A_{Z,4}| = 23.4$ MHz.



Figure S2. ¹³C W-band (93.88 GHz) Davies-type ENDOR spectra (black curves) from $[5a,8-^{13}C_2]$ FAD-reconstituted *Ec*PL, recorded at 80 K. For improved global least-squares fitting, the data have been symmetrized with respect to the Larmor precession frequency at around 35.85 MHz. For this figure, the spectra have also been normalized to equal maximum intensity with respect to each other. The spectral simulations (red curves) have been obtained with the parameters compiled in Table S2.

For ¹³C8, DFT predicts a hyperfine tensor of axial symmetry with very large anisotropy, $A_X =$ +1.1 MHz, $A_Y =$ +1.2 MHz and $A_Z =$ +20.1 MHz, which is in good agreement with the ENDOR resonances "4". The predicted hyperfine anisotropy for ¹³C5a is much smaller: $A_X =$ -10.6 MHz, $A_Y =$ -10.0 MHz and $A_Z =$ -16.2 MHz. Therefore, we assign the resonances "3" to ¹³C5a.

Each spectrum shown in Fig. S2 contains one pair of resonances (marked with asterisks) that is neither assigned to the hyperfine couplings of ¹³C8 nor ¹³C5a. The signals are quite narrow and their position hardly changes on altering the magnetic field value, at which the ENDOR data were recorded. A nearly isotropic hyperfine coupling is extracted if one additional set of

hyperfine principal values is allowed to enter the simulations: $|A_{X,*}| \approx |A_{Y,*}| \approx |A_{Z,*}| \approx |A_{iso,*}|$. This value agrees reasonably well with DFT predictions for ¹³C7: $A_X = -3.8$ MHz, $A_Y = -3.8$ MHz and $A_Z = -6.0$ MHz. Tentatively we assign the additional hyperfine resonances "*" to ¹³C7. This notion is corroborated by the determination of the hyperfine coupling of this nucleus using difference ENDOR spectroscopy; see also below. The resonance of ¹³C7 is observed here because in the synthesis procedure of [5a,8-¹³C₂]FAD the [7,9a-¹³C₂]FAD isotopomer is also be formed to some extent.¹ The signal of ¹³C9a escapes clear detection because its hyperfine tensor is quite anisotropic (see Table S6) and the intensities of its resonances in ENDOR consequently rather low, and therefore, do not exceed the noise level.

All parameters used to obtain the simulations shown as red curves in Fig. S2 are compiled in Table S2.

Table S2. Hyperfine data used to simulate the ENDOR spectra of $[5a,8^{-13}C_2]FAD$ -reconstituted *EcPL*. The principal values are accurate within ±0.2 MHz, and the Euler angles within ±0.1 rad. The signs of the hyperfine components were taken from DFT calculations. Γ_G is the fwhm of Gaussian broadening. The *w*_N are weighting factors that account, *e. g.*, for the specific relaxation properties of the individual ¹³C nuclei.

Parameter	¹³ C5a	¹³ C8	¹³ C7
A_X / MHz	(-)15.1	(+)0.8	(-)4.1
$A_{\rm Y}/{ m MHz}$	(-)13.0	(+)1.8	(-)4.1
$A_{\rm Z}/{ m MHz}$	(-)11.6	(+)23.4	(-)3.7
$a_{\rm iso}/{ m MHz}$	(-)13.2	(+)8.7	(-)4.0
$\Gamma_{ m G}/{ m mT}$	0.35	0.46	0.15
lpha / rad	+0.2	-0.1	0
eta / rad	+0.5	-0.2	0
γ/rad	+0.4	+0. <u>9</u>	0
W _N	0.62	0.35	0.03

Analysis of ENDOR data from [1',7α,9-¹³C₃]FAD bound to *Escherichia coli* DNA photolyase

The frozen-solution ENDOR spectra of $[1',7\alpha,9^{-13}C_3]$ FAD-reconstituted *Ec*PL, recorded at magnetic-field values corresponding to the canonical **g**-principal axes orientations, exhibit two pairs of quite intense signals (resonances "5" and "6", see Fig. S3) from which the following components of two hyperfine tensors can be extracted: $|A_{X,5}| = 6.1$ MHz, $|A_{Y,5}| = 5.6$ MHz and $|A_{Z,5}| = 4.6$ MHz, and $|A_{X,6}| = 5.5$ MHz, $|A_{Y,6}| = 5.9$ MHz and $|A_{Z,6}| = 11.6$ MHz. Due to the strong overlap of the resonances, a higher error margin is assumed for these hyperfine components. A much weaker pair of resonances ("7") with virtually no variation of the splitting with respect to the magnetic-field position is also observed that is described with $|A_{X,7}| \approx |A_{Y,7}| \approx |A_{Z,7}| \approx |A_{iso,7}| = 0.5$ MHz.



Figure S3. ¹³C W-band (93.88 GHz) Davies-type ENDOR spectra (black curves) from $[1',7\alpha,9-^{13}C_3]FAD$ -reconstituted *EcPL*, recorded at 80 K. For improved global least-squares fitting, the data have been symmetrized with respect to the Larmor precession frequency at around 35.85 MHz. For this figure, the spectra have also been normalized to equal maximum intensity with respect to each other. The spectral simulations (red curves) have been obtained with the parameters compiled in Table S3.

For ¹³C1', DFT predicts a hyperfine tensor with small anisotropy, $A_X = -2.5$ MHz, $A_Y = -3.5$ MHz and $A_Z = -3.2$ MHz, and for ¹³C9 one with larger anisotropy and of axial symmetry: $A_X = -4.4$ MHz, $A_Y = -3.9$ MHz and $A_Z = -10.1$ MHz. Therefore, we assign resonances "5" to ¹³C1', and resonances "6" to ¹³C9. Consequently, resonances "7" must arise from ¹³C7 α , for which only small spin-density is predicted from quantum-chemical calculations, as is evident from its nearly isotropic hyperfine tensor: $A_X = -0.3$ MHz, $A_Y = -0.5$ MHz and $A_Z = -0.1$ MHz. All

parameters used to obtain the simulations shown as red curves in Fig. S3 are compiled in Table S3.

Table S3. Hyperfine data used to simulate the ENDOR spectra of $[1',7\alpha,9^{-13}C_3]$ FAD-reconstituted *Ec*PL. The principal values of ¹³C1' are accurate within ±0.2 MHz, the ones from ¹³C9 to ±0.4 MHz, and the ones from ¹³C7 α to ±0.1 MHz. The Euler angles are accurate within ±0.2 rad. The signs of the hyperfine components were taken from DFT calculations. Γ_G is the fwhm of Gaussian broadening. The *w*_N are weighting factors that account, *e. g.*, for the specific relaxation properties and isotope enrichment of the individual ¹³C nuclei.

Parameter	¹³ C1'	¹³ C9	¹³ C7a
A_X / MHz	(-)6.1	(–)5.5	(-)0.5
A_Y/MHz	(–)5.6	(-)5.9	(-)0.5
A_Z /MHz	(-)4.6	(-)11.6	(-)0.5
$a_{\rm iso}/{ m MHz}$	(–)5.4	(–)7.7	(-)0.5
$\Gamma_{\rm G}/{ m mT}$	0.18	0.25	0.10
α / rad	-0.6	+0.8	0
β / rad	-1.2	+0.4	0
γ / rad	-1.2	0.6	0
WN	0.57	0.42	0.01

Analysis of ENDOR data from [2,4a-¹³C₂]FAD bound to *Escherichia coli* DNA photolyase

The frozen-solution ENDOR spectra of $[2,4a^{-13}C_2]$ FAD-reconstituted *Ec*PL, recorded at magnetic-field values corresponding to the canonical **g**-principal axes orientations, reveal two pairs of quite intense signals (resonances "8" and "9", see Fig. S4) from which the following components of two hyperfine tensors can be extracted: $|A_{X,8}| = 11.1$ MHz and $|A_{Y,8}| = 12.2$ MHz (a third component could not be determined; the fitting routine revealed a value of 9.6 MHz, which did yield the best possible agreement with the experimental ENDOR data, but this value is certainly an artefact given the limited signal-to-noise ratios of all three spectra), and $|A_{X,9}| = 2.0$ MHz, $|A_{Y,9}| = 1.9$ MHz and $|A_{Z,9}| = 0.3$ MHz.



Figure S4. ¹³C W-band (93.88 GHz) Davies-type ENDOR spectra (black curves) from $[2,4a-^{13}C_2]FAD$ -reconstituted *EcPL*, recorded at 80 K. For improved global least-squares fitting, the data have been symmetrized with respect to the Larmor precession frequency at around 35.85 MHz. For this figure, the spectra have also been normalized to equal maximum intensity with respect to each other. The spectral simulations (red curves) have been obtained with the parameters compiled in Table S4.

DFT predicts a hyperfine tensor with very large anisotropy and axial symmetry for ¹³C4a, $A_X = -10.1$ MHz, $A_Y = -9.5$ MHz and $A_Z = +59.2$ MHz, and for ¹³C2 one with much smaller (absolute) hyperfine components: $A_X = -4.2$ MHz, $A_Y = -2.8$ MHz and $A_Z = -5.7$ MHz. Therefore, we assign resonances "9" to ¹³C2, and resonances "8" to ¹³C4a. However, the contribution of A_Z to the powder pattern of ¹³C4a could not be detected in the experimental ENDOR spectrum. DFT predicts values >50 MHz for this principal value. Such strong anisotropy leads to a distribution of signal intensity over a wide radio-frequency range, and hence, very low absolute signal amplitude. Therefore, we exhibit only two out of the three hyperfine components from ENDOR

spectroscopy. All parameters used to obtain the overall simulations shown as red curves in Fig. S4 are compiled in Table S4.

Table S4. Hyperfine data used to simulate the ENDOR spectra of $[2,4a^{-13}C_2]$ FAD-reconstituted *Ec*PL. The principal values are accurate within ±0.2 MHz, and the Euler angles within ±0.1 rad. The signs of the hyperfine components were taken from DFT calculations. Γ_G is the fwhm of Gaussian broadening. The *w*_N are weighting factors that account, *e. g.*, for the specific relaxation properties of the individual ¹³C nuclei.

Parameter	¹³ C2	¹³ C4a
A_X / MHz	(-)2.0	(-)12.2
A_Y/MHz	(-)1.9	(-)11.1
A_Z / MHz	(-)0.3	n.d.
a _{iso} / MHz	(-)1.4	_a
$\Gamma_{\rm G}/{ m mT}$	0.16	0.22
α / rad	+1.2	-2.1
β / rad	+0.5	+0.6
γ / rad	-0.9	-0.1
WN	0.309	0.691

^aQuantum-chemical calculations predict an isotropic hyperfine coupling constant of about +12 MHz for ¹³C4a. Recent determinations of $A_{iso}(^{13}C4a)$ using photo-CIDNP on the neutral FMN radical in bulk aqueous solution question such a large isotropic hyperfine coupling and suggest values near 0 [N. Pompe *et al.*, unpublished data].

Analysis of ENDOR data from [4,10a-¹³C₂]FAD bound to *Escherichia coli* DNA photolyase

The frozen-solution ENDOR spectra of $[4,10a^{-13}C_2]$ FAD-reconstituted *Ec*PL, recorded at magnetic-field values corresponding to the canonical **g**-principal axes orientations, reveal two pairs of overlapping powder patterns, one of large anisotropy (resonances "10", see Fig. S5) and the other of small anisotropy (resonances "11"), from which the following hyperfine components can be extracted: $|A_{X,10}| = 11.2$ MHz, $|A_{Y,10}| = 13.8$ MHz and $|A_{Z,10}| = 4.0$ MHz, and $|A_{X,11}| = 15.4$ MHz, $|A_{Y,11}| = 13.6$ MHz and $|A_{Z,11}| = 13.1$ MHz.



Figure S5. ¹³C W-band (93.97 GHz) Davies-type ENDOR spectra (black curves) from $[4,10a-{}^{13}C_2]FAD$ -reconstituted *EcPL*, recorded at 80 K. For improved global least-squares fitting, the data have been symmetrized with respect to the Larmor precession frequency at around 35.88 MHz. For this figure, the spectra have also been normalized to equal maximum intensity with respect to each other. The spectral simulations (red curves) have been obtained with the parameters compiled in Table S5.

For ¹³C4, DFT predicts a hyperfine tensor with moderate anisotropy, $A_X = -11.6$ MHz, $A_Y = -12.5$ MHz and $A_Z = -9.4$ MHz, and for ¹³C10a one with smaller anisotropy and of axial symmetry: $A_X = -12.5$ MHz, $A_Y = -14.3$ MHz and $A_Z = -14.1$ MHz. Therefore, we assign resonances "10" to ¹³C4, and resonances "11" to ¹³C10a. All parameters used to obtain the overall simulations shown as red curves in Fig. S5 are compiled in Table S5.

Table S5. Hyperfine data used to simulate the ENDOR spectra of $[4,10a^{-13}C_2]FAD$ -reconstituted *EcPL*. The principal values are accurate within ±0.2 MHz, and the Euler angles within ±0.2 rad. The signs of the hyperfine components were taken from DFT calculations. Γ_G is the fwhm of Gaussian broadening. The *w*_N are weighting factors that account, *e. g.*, for the specific relaxation properties of the individual ¹³C nuclei.

Parameter	¹³ C4	¹³ C10a
A_X / MHz	(-)11.2	(-)15.4
A_Y / MHz	(-)13.8	(-)13.6
A_Z / MHz	(-)4.0	(-)13.1
a _{iso} / MHz	(-)9.7	(-)14.0
$\Gamma_{\rm G}/{ m mT}$	0.08	0.37
α / rad	+0.3	+0.3
β / rad	+0.5	+0.6
γ / rad	+1.0	+0.6
WN	0.49	0.51

Determination of the hyperfine tensors of ¹³C7 and ¹³C9a using difference-ENDOR spectroscopy

To determine hyperfine data for the two remaining ¹³C nuclei in the 7,8-dimethyl isoalloxazine moiety of the neutral flavin radical, ¹³C7 and ¹³C9a, we applied difference-ENDOR spectroscopy. Successive subtraction of the ENDOR spectra from EcPL reconstituted with $[5a, 8^{-13}C_2]FAD$, $[6, 8\alpha^{-13}C_2]FAD$ and $[1^{\circ}, 7\alpha, 9^{-13}C_3]FAD$ from the ENDOR spectra of *EcPL* reconstituted with [xylene- ${}^{13}C_8$]FAD (= [5a,6,7,7\alpha,8,8\alpha,9,9a- ${}^{13}C_8$]FAD) results in a set of difference-ENDOR spectra with at the end only positive signals from the remaining nuclei ¹³C7 and ¹³C9a, see Fig. S6. The negative signals in Fig. S6 arise from ¹³C1' absent in [xylene-¹³C₈]FAD and will not be discussed further, since the hyperfine principal values of this nucleus have been determined directly using the $[1',7\alpha,9-^{13}C_3]FAD$ isotopolog (see above). Two sets of positive difference-ENDOR amplitudes are observed, resonances "12" and "13". The sets of very narrow lines with a splitting of around 4 MHz are assigned to ¹³C7. The hyperfine components extracted from the difference ENDOR spectra, $|A_X| = |A_Y| = 4.12$ MHz and $|A_Z| =$ 3.74 MHz, agree reasonably well with predictions from DFT for this nucleus: $A_X = -3.8$ MHz, $A_Y = -3.8$ MHz and $A_Z = -6.0$ MHz. (These resonances have also been observed as an "impurity" (resonances marked with "*" in Fig. S2) in the *Ec*PL reconstituted [5a,8-¹³C₂]FAD, see above). The more anisotropic set of hyperfine resonances, "13", is consequently assigned to ¹³C9a. Due to distortions arising from the overlap with the difference-ENDOR signals from ¹³C1', its hyperfine components can be read out from the spectra with only moderate accuracy: $|A_X| = 4.8$ MHz, $|A_Y| = 4.6$ MHz, and $|A_Z| = 1.7$ MHz. DFT predicts $A_X = -1.9$ MHz, $A_Y = -1.6$ MHz and $A_Z = +12.2$ MHz for this nucleus.

Table	S6 .	Hyperfine	data	extracted	from th	e differer	ce-ENDOR	spectra	of EcPL	with	"[1',7,9a- ¹	³ C ₃]FAD"
depicte	d in	Fig. S6. Th	he pri	ncipal val	ues are a	ccurate w	thin ± 0.1 M	Hz (¹³ C7) and ± 1	MHz ((¹³ C9a).	

Parameter	¹³ C7	¹³ C9a	
A_X / MHz	(-)4.1	(-)4.8	
$A_Y/$ MHz	(-)4.1	(-)4.6	
A_Z / MHz	(-)3.7	(+)1.7	
a _{iso} / MHz	(-)4.0	(-)2.6	



Figure S6. ¹³C W-band difference-ENDOR data (green curves) from *Ec*PL with the isotopolog species marked on the left sides of the sets of spectra.

Table S7. Compilation of isotropic hyperfine couplings of 13 C nuclei in the neutral flavin radical in *Ec*PL from the various ENDOR experiments, DFT calculations based on the PDB structure of *Ec*PL, and DFT calculations following geometry relaxation using QM/MM.

Nucleus	A _{iso} / MHz (experiment) ^a	A _{iso} / MHz (DFT)	A _{iso} / MHz (QM/MM/DFT) ^b
¹³ C2	(-)1.4	-4.2	-3.7
¹³ C4	(-)9.7	-11.2	-13.1
¹³ C4a	_	+13.2	+15.1
¹³ C5a	(-)13.2	-12.3	-14.0
¹³ C6	(+)5.1	+4.2	+5.0
¹³ C7	(-)4.0	-4.5	-4.5
¹³ C7α	(+)0.5	+0.3	0.0
¹³ C8	(+)8.7	+7.5	+7.7
¹³ C8α	(-)3.8	-3.0	-3.3
¹³ C9	(–)7.7	-6.1	-5.6
¹³ C9a	(+)2.6	+2.9	+0.6
¹³ C10a	(-)14.0	-13.6	-17.1
¹³ C1'	(-)5.4	-3.1	-2.3

^aSign of hyperfine couplings taken from DFT calculations.

^bTo obtain hyperfine data from quantum mechanics/molecular mechanics (QM/MM) computations, we performed in a first step molecular-dynamics (MD) simulations using GROMACS to obtain "equilibrated" structure data for photolyase. The force fields Gromos53a6 and Gromos G96 53a6 united-atoms² were used for the protein and the FAD cofactor, respectively. An MD trajectory of 4 ns was computed and representative "snapshots" of the FAD cofactor and its immediate surroundings (comprising the five amino acids Ser-238, Arg-344, Asp-372, Asn-378 and Trp-382) selected that were then used, in the second step, as input coordinates for a QM treatment at the DFT level. Principal values of ¹³C hyperfine tensors were computed using ORCA with the B3LYP hybrid functional in combination with the EPR-II basis set, and subsequently averaged to yield A_{iso} values.



Figure S7. Compilation of (absolute) isotropic ¹³C hyperfine couplings from ENDOR (white squares), QM/MM/DFT (blue squares) and a "simple" DFT treatment (red squares). Values have been taken from Table S7.



Figure S8. ENDOR/QM/MM/DFT (blue circles) and ENDOR/DFT (red circles) correlation plots for the isotropic ¹³C hyperfine coupling constants of the neutral FAD radical in *E. coli* photolyase. The dashed blue line ($R^2 = 0.9321$; slope 1.0619) and the dashed red line ($R^2 = 0.9576$; slope 0.9440) represent linear regression fits constrained to go through the origin.

References

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