

## S1 Appendix. Energy distribution function

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382

In the case of a particle beam with continuous particle density over time, as shown in Fig. 2a, the fraction of particles  $dN$  in a time interval  $dt$  is constant:  $\frac{dN}{dt} = c$ . The number of particles  $dN$  within an energy bin  $dE$  as a function of the change  $E$  in beam energy after the beam has passed through a buncher cavity is given by

$$\frac{dN}{dE} = \frac{dN}{dt \cdot \frac{\partial E}{\partial t}} = \frac{c}{\frac{\partial E}{\partial t}}. \quad (\text{A.1})$$

For a particle beam modulated by a sinusoidal RF signal, the energy modulation  $E(t)$  of a particle of charge  $q$  that passes the buncher at a time point  $t$  is described by

$$E(t) = qU_b \sin(\omega t). \quad (\text{A.2})$$

$\omega = 2\pi f$  is the angular frequency of the RF signal feeding the buncher.  $U_b$  is the total voltage a particle can see when passing the buncher. The derivative of equation A.2 is

$$\frac{\partial E}{\partial t} = qU_b \omega \cos(\omega t) = qU_b \omega \sqrt{1 - \sin^2(\omega t)}. \quad (\text{A.3})$$

Solving equation A.2 for  $\sin(\omega t)$  and substituting it into equation A.3 results in

$$\frac{\partial E}{\partial t} = qU_b \omega \sqrt{1 - \left(\frac{E(t)}{qU_b}\right)^2} = \omega \sqrt{(qU_b)^2 - E(t)^2} = \omega \sqrt{E_b^2 - E^2}. \quad (\text{A.4})$$

Where  $E_b$  is the maximum energy imposed by the cavity to a particle traversing it. Inserting equation A.4 into equation A.1 results in the energy distribution function

$$\frac{dN}{dE} = \frac{c}{\omega \sqrt{E_b^2 - E^2}} \quad (\text{A.5})$$

for any beam energy  $E < E_b$ . Using  $E = qU$  transforms equation A.5 in the particle distribution function

$$g(U|U_b) = \frac{dN}{dU} = \frac{c}{\omega \sqrt{U_b^2 - U^2}} = \frac{C}{\sqrt{U_b^2 - U^2}} \quad (\text{A.6})$$

that gives the number of detected particles  $N$  in dependence of the acceleration voltage  $U$  for any  $U < U_b$ . Where  $U_b$  is the maximum acceleration voltage (buncher amplitude).

392  
393