

Supplementary material S1: the two-pool MT model

The two-pool magnetisation transfer (MT) model is used in this paper in both simulations and in vivo analyses. The model describes the evolution of the magnetisation of two exchanging ^1H pools, i.e. free protons in bulk water and protons bound to macromolecules (Henkelman et al., 1993), in presence of off-resonance irradiation $b_{1,\text{off}}(t)$ with carrier frequency $f_0 + \Delta f_c$ (with Δf_c being the offset with respect to water resonance frequency f_0), assumed to be played out at times $0 \leq t \leq t_{\text{off}}$. On-resonance excitation is then assumed to be played out at $t > t_{\text{off}}$, after a generic delay t_{delay} .

Let us indicate with $\mathbf{m}^{\text{F}}(t) \stackrel{\text{def}}{=} [m_x^{\text{F}}(t) \ m_y^{\text{F}}(t) \ m_z^{\text{F}}(t)]^{\text{T}}$ and $\mathbf{m}^{\text{B}}(t) \stackrel{\text{def}}{=} [m_x^{\text{B}}(t) \ m_y^{\text{B}}(t) \ m_z^{\text{B}}(t)]^{\text{T}}$ the free and bound proton magnetisations during the off-resonance irradiation, and with T_1^{F} and T_1^{B} and with T_2^{F} and T_2^{B} their respective longitudinal and transverse relaxation times (Portnoy and Stanisiz, 2007). Let us also indicate with k the exchange rate between free to bound protons, with BPF the bound pool fraction (fraction of protons belonging to the bound pool) and with γ the proton gyromagnetic ratio ($\frac{\gamma}{2\pi} \cong 42.577 \frac{\text{MHz}}{\text{T}}$).

The MT-weighting factor w in Eq. 5 (main manuscript) is defined as the value of the free pool longitudinal magnetisation at the end of off-resonance irradiation, i.e.:

$$w = m_z^{\text{F}}(t = t_{\text{off}}). \quad (\text{S1.1})$$

The value of w in Eq. S1.1 is obtained directly by numerical integration of the two-pool Bloch equations, which are written as:

$$\frac{d}{dt} \begin{bmatrix} m_x^{\text{F}} \\ m_y^{\text{F}} \\ m_z^{\text{F}} \\ m_z^{\text{B}} \end{bmatrix} (t) = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & -\frac{1}{T_2^{\text{F}}} & 2\pi \Delta f_c & 0 & 0 \\ 0 & -2\pi \Delta f_c & -\frac{1}{T_2^{\text{F}}} & \gamma b_{1,\text{off}}(t) & 0 \\ \frac{1}{T_1^{\text{F}}} & 0 & -\gamma b_{1,\text{off}}(t) & -\left(\frac{1}{T_1^{\text{F}}} + k\right) & k \frac{(1 - \text{BPF})}{\text{BPF}} \\ \frac{1}{T_1^{\text{B}}} \frac{\text{BPF}}{(1 - \text{BPF})} & 0 & 0 & k & -\left(\frac{1}{T_1^{\text{B}}} + R_{\text{RF}}^{\text{B}} + k \frac{(1 - \text{BPF})}{\text{BPF}}\right) \end{bmatrix} \begin{bmatrix} \frac{1}{2} \\ m_x^{\text{F}} \\ m_y^{\text{F}} \\ m_z^{\text{F}} \\ m_z^{\text{B}} \end{bmatrix} (t). \quad (\text{S1.2})$$

For the integration of equation S1.2, it is assumed that $\mathbf{m}^F(t=0) \stackrel{\text{def}}{=} [0 \ 0 \ (1 - \text{BPF})]^T$, $\mathbf{m}^B(t=0) \stackrel{\text{def}}{=} [0 \ 0 \ \text{BPF}]^T$, $m_x^B(t = t_{\text{off}}) = m_y^B(t = t_{\text{off}}) \approx 0$ and that the bound pool off-resonance absorption can be modelled based on a super-Lorentzian line shape. Under this assumption, the absorption term R_{RF}^B is a function of $(t, \Delta f_c, T_2^B)$ and is modelled as (Portnoy and Stanisz, 2007):

$$R_{\text{RF}}^B(t, \Delta f_c, T_2^B) = \pi \gamma^2 b_{1,\text{off}}^2(t) \int_0^{\frac{\pi}{2}} \sqrt{\frac{2}{\pi}} \frac{\sin(\varphi)}{|3 \cos^2(\varphi) - 1|} T_2^B \exp\left(-2 \left(\frac{2\pi \Delta f_c T_2^B}{3 \cos^2(\varphi) - 1}\right)^2\right) d\varphi \quad (\text{S1.3}).$$

For the practical integration of Eq. S1.2, it is assumed that $T_1^B \approx 1$ s (Battiston et al., 2018), and that the free pool longitudinal relaxation time (i.e. T_1^F) is linked to the observable T_1 relaxation time as (Portnoy and Stanisz, 2007):

$$\frac{1}{T_1^F} \approx \frac{1}{T_1} - \frac{k \left(\frac{1}{T_1^B} - \frac{1}{T_1}\right)}{\frac{1}{T_1^B} - \frac{1}{T_1} + k \left(\frac{1 - \text{BPF}}{\text{BPF}}\right)}. \quad (\text{S1.4})$$

References

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