# Supplementary Information: Systematic Optimization of Magnesium Force Field Parameters for Biomolecular Simulations with Accurate Solvation, Ion-pairing, and Water Exchange Properties in SPC/E, TIP3P-fb, TIP4P/2005, TIP4P-Ew, and TIP4P-D.

Kara K. Grotz and Nadine Schwierz<sup>\*</sup>

Department of Theoretical Biophysics, Max-Planck-Institute of Biophysics, Frankfurt am Main, Germany.

E-mail:<nadine.schwierz@biophys.mpg.de>

# Contents





S2

# <span id="page-2-0"></span>S1 Supplementary methods

# <span id="page-2-1"></span>S1.1 Literature values for central  $Mg^{2+}$  properties

Previous force fields for  $Mg^{2+}$  failed to provide a quantitative description of central properties of the ion. In Table [S1](#page-2-3) the central properties are listed for  $Mg^{2+}$  as obtained from the literature and compared to experiments.

<span id="page-2-3"></span>Table S1: Values for  $Mg^{2+}$  models from the literature for central properties of the ion in comparison to experiments and one model of the current work.

	$R_1$	$\Delta G_{\rm solv}$	k <sub>i</sub>	$\Delta G_{\rm h}^0$
	$\vert$ mm	[kJ/mol]	$[s^{-1}]$	$[k_B T]$
Mamatkulov-Schwierz (TIP3P) $1$	$0.196$ <sup>1</sup>	$-2\overline{531.1}$ <sup>1</sup>	$24 \pm 9^2$	$-14.53 \pm 0.95$ <sup>3</sup>
Allner-Villa (TIP3P) $4$	$0.204$ <sup>4</sup>	$-2397.35$	« $2.39 \times 10^{5}$ <sup>6</sup>	$-12.17 \pm 0.14$ <sup>6</sup>
Li-Merz (TIP3P) $^7$	$0.194$ <sup>7</sup>	$-2527.0$ <sup>7</sup>	$\approx 3.52 \times 10^{4}$ 6	n.a.
Li-Merz (TIP3P) [12-6-4] $8$	0.208 <sup>8</sup>	$-2519.1$ $8$	$(1.44 \pm 0.03) \times 10^{7}$ 6	n.a.
Mamatkulov-Netz (SPC/E) $^9$	$0.196$ <sup>9</sup>	$-2532.0$ <sup>9</sup>	n.a.	$-14.35 \pm 0.48$
Aquist (SPC/E) $^{10}$	$0.199\ ^{5}$	$-2505.35$	n.a.	n.a.
Li-Merz (SPC/E) $^7$	$0.195$ <sup>7</sup>	$-2521.2$ <sup>7</sup>	n.a.	n.a.
Li-Merz (SPC/E) [12-6-4] $8$	0.208 <sup>8</sup>	$-2524.98$	n.a.	n.a.
Li-Merz (TIP4P-Ew) $^7$	0.199 <sup>7</sup>	$-2520.3$ <sup>7</sup>	n.a.	n.a.
Li-Merz (TIP4P-Ew) [12-6-4] $8$	0.208 <sup>8</sup>	$-2526.2$ $8$	n.a.	$-7.65 \pm 0.35$ <sup>6</sup>
$microMg$ (SPC/E) [current work]	0.209	$-2530.5$	$(9.62 \pm 1) \times 10^5$	$-1.266 \pm 0.62$
Exp.	$0.209 \pm 0.004^{11}$	$-2532^{12}$	$5.3 - \overline{6.7 \times 10^{5}}^{13,14}$	$-1.\overline{036^{15}}$

#### <span id="page-2-2"></span>S1.2 Water models

We used five different rigid water models and listed their partial charges and Lennard-Jones parameters in Table [S2.](#page-2-4) For SPC/E, the bond length is fixed at 0.1 nm and the bond angle at 109.47°. For TIP3P-fb, the bond length is fixed at about 0.10118 nm and the bond angle at about 108.15◦ . Otherwise, the bond length is 0.09572 nm and the bond angle is  $104.52^\circ$ .

<span id="page-2-4"></span>Table S2: Parameters of different water models. The 3-site waters (SPC/E, TIP3P-fb) carry their negative partial charge  $(q_{\text{Ow}/\text{Mw}})$  on the oxygen Ow, whereas for the 4-site waters (TIP4P/2005, TIP4P-Ew, and TIP4P-D), it is located on the slightly shifted additional site Mw.  $q<sub>H</sub>$  is the partial charge of the hydrogen atoms and  $\sigma_{\text{Ow}}$  and  $\varepsilon_{\text{Ow}}$  the Lennard-Jones parameters of the water oxygen.



#### <span id="page-3-0"></span>S1.3 Simulation setups

In all simulations long-range electrostatic forces were handled applying the particle-mesh Ewald method [16](#page-20-3) with periodic boundary conditions, a Fourier spacing of 0.12 nm and a grid interpolation up to order 4. Shortrange Coulomb and Lennard-Jones interactions were cut off at 1.2 nm. In all simulations, a 2 fs timestep was used. In all simulations, ions and rigid water models are used, in those with AMBER we employed  $SETILE$ , <sup>[17](#page-20-4)</sup> and in those with GROMACS we employed LINCS. <sup>[18](#page-20-5)</sup> Hydrogen bonds are constrained in all simulations containing DMP using LINCS. All the simulations started with an energy minimization by steepest decent, which was followed by two equilibration steps of 0.5 ns in NVT and 1 ns in NPT. The temperature was kept at 300 K and the pressure at 1 atm using thermostat and barostat of Berendsen. [19](#page-20-6) Position restraints were applied to all ions to ensure that no ions are pairing before hydration shells can properly form. The restraints were released in the production run. Lorentz-Berthelot combination rules were chosen for all simulations. The simulations to compute solvation free energy  $\Delta G_{solv}$ , radius of the first hydration shell  $R_1$ , and coordination number  $n_1$ , were performed employing the thermostat and barostat of Berendsen<sup>[19](#page-20-6)</sup> with  $\tau = 0.1$  and  $\tau_p = 1.0$  to ensure a temperature of 300 K and roughly an atmospheric pressure of 1 bar. All other simulations were performed with the velocity rescaling thermostat by Bussi et al.<sup>[20](#page-20-7)</sup> with  $\tau = 0.1$ . In case of the simulations in isothermal-isobaric ensemble (NPT in Table [S3](#page-3-1) and [S4\)](#page-4-2) the Parrinello-Rahman barostat<sup>[21](#page-20-8)</sup> with  $\tau_p = 5.0$  was applied. Frames were written out every 0.2 ps if not stated otherwise. For both engines GROMACS and AMBER, the same setups were employed.

<span id="page-3-1"></span>Table S3: This table shows all simulation setups of the 3-site water models SPC/E and TIP3P-fb. In 'Phys. property', the physical properties are listed that we obtained from the respective simulation using the respective 'Method' (FEP: free energy perturbation, unbiased: straight forward simulations without additional biases, US: umbrella sampling). 'System' lists all particles of the respective simulation and 'Duration' their duration (products show for FEP and US simulations the individual simulation windows). 'L' indicates the simulation box size (edge length) for (cu) cubic and (do) rhombic dodecahedron box shapes. 'Ensemble' expresses if an canonical ensemble (with constant number of particles  $N$ , constant volume V, and constant temperature  $T$ ), or an isobaric-isothermal ensemble (with constant pressure P and flexible simulation box size) was used. [GMX and AMBER specify the setup used for the 12-6 based (GMX) or the 12-6-4 based parameters (AMBER)].

Phys. property	Method	System	Duration	L	Ensemble
$\Delta G_{\rm solv}, R_1, n_1$	<b>FEP</b>	$\overline{1 \text{ Mg}^{2+}$ , 506 water	$40 \times 1$ ns	$2.5 \text{ nm}$ (cu)	<b>NPT</b>
$D_0$	unbiased	$1 \text{ Mg}^{2+}$ , 506 water	$100$ ns	$2.5 \text{ nm}$ (cu)	<b>NVT</b>
$a_{\rm cc}, k$	unbiased	$39 \text{ Mg}^{2+}$ , 78 Cl , 2048 water (1 M)	$1 \mu s$	$4 \text{ nm}$ (cu)	<b>NPT</b>
$a_{\rm cc}$	unbiased	$73 \text{ Mg}^{2+}$ , 146 Cl , 1961 water $(2 \text{ M})$	$150$ ns	$4 \text{ nm}$ (cu)	<b>NPT</b>
$a_{\rm cc}$	unbiased	40 $Mg^{2+}$ , 20 Cl , 2105 water (0.5 M)	$150$ ns	$4 \text{ nm}$ (cu)	<b>NPT</b>
$a_{\rm cc}$	unbiased	$20 \text{ Mg}^{2+}$ , 10 Cl , 2135 water (0.25 M)	$150$ ns	$4 \text{ nm}$ (cu)	<b>NPT</b>
$F(r_{\rm MgOw})$	US(GMX)	$1 \text{ Mg}^{2+}$ , 505 water	$68 \times 3$ ns	$2.5 \text{ nm}$ (cu)	<b>NPT</b>
$F(r_{\rm MgOw})$	US (AMBER)	$1 \text{ Mg}^{2+}$ , 500 water	$68 \times 3$ ns	$2.5 \text{ nm}$ (cu)	<b>NPT</b>
$F(r_{\rm MgOw_1}, r_{\rm MgOw_2})$	US.	$1 \text{ Mg}^{2+}$ , 506 water	$946 \times 5$ ns	$2.5 \text{ nm}$ (cu)	<b>NVT</b>
$\Delta G_{\rm b}^0$	<b>FEP</b>	1 DMP, $1 \text{ Mg}^{2+}$ , $1 \text{ Cl}$ , $1492 \text{ water}$	$20 \times 15$ ns	$4 \text{ nm}$ (do)	<b>NPT</b>
$\Delta G_{\rm b}^0$ , $F(r_{\rm MgOP})$	US	1 DMP, $1 \text{ Mg}^{2+}$ , $1 \text{ Cl}$ , $1492 \text{ water}$ ,	$67 \times 20$ ns	$4 \text{ nm}$ (do)	<b>NVT</b>
$\Delta G_{\text{ref}}^0$ , $F(r_{\text{ref}-\text{OP}})$	US	1 DMP, 1 ref atom, 1 Cl <sup>-</sup> , 1492 water	$67 \times 20$ ns	$4 \text{ nm}$ (do)	<b>NVT</b>

<span id="page-4-2"></span>Table S4: This table lists all simulation setups using the 4-site water models TIP4P/2005, TIP4P-Ew, and TIP4P-D that differ from Table [S4.](#page-4-2) If the required 'Phys. property' is not listed here, the same setup is applied as in Table [S3.](#page-3-1)

Phys. property	Method	System	Duration		Ensemble
$a_{\rm cc}, k$	unbiased	39 $Mg^{2+}$ , 78 Cl <sup>-</sup> , 2055 water (1 M)		$150 \text{ ns}$ 4 nm (cu)	<b>NPT</b>
$a_{\rm cc}$	unbiased	73 Mg <sup>2+</sup> , 146 Cl <sup>-</sup> , 1909 water $(2 M)$		$150 \text{ ns}$ 4 nm (cu)	<b>NPT</b>
$a_{cc}$	unbiased	$20 \text{ Mg}^{2+}$ , 40 Cl <sup>-</sup> , 2135 water (0.5 M)		$150 \text{ ns}$ 4 nm (cu)	<b>NPT</b>
$a_{cc}$	unbiased	$10 \text{ Mg}^{2+}$ , $20 \text{ Cl}$ <sup>-</sup> , $2171 \text{ water}$ (0.25 M)		$150 \text{ ns}$ 4 nm (cu)	<b>NPT</b>
$\Delta G_{\rm h}^0$	FEP	1 DMP, 1 $Mg^{2+}$ , 1 Cl <sup>-</sup> , 1536 water	$20 \times 15$ ns 4 nm (do)		<b>NPT</b>
$\Delta G_{\rm b}^0$ , $F(r_{\rm MgOP})$	US	1 DMP, 1 $Mg^{2+}$ , 1 Cl <sup>-</sup> , 1536 water	$67 \times 20$ ns 4 nm (do)		<b>NVT</b>
$\Delta G_{\text{ref}}^0$ , $F(r_{\text{ref}-\text{OP}})$	US	$1 \text{ DMP}, 1 \text{ ref atom}, 1 \text{ Cl}^-, 1536 \text{ water}$	$67 \times 20$ ns 4 nm (do)		<b>NVT</b>

### <span id="page-4-0"></span>S1.4 Diffusion coefficient

Diffusion coefficients  $D_0$  were calculated from 10 ns long NVT simulations. The first 1 ns was excluded in each trajectory from the analysis. We calculated the coefficients  $D_{\text{pbc}}(L)$  from a straight line fit of the slope of the mean-squared displacement of the single ion and took only the linear part into account. The obtained diffusion coefficient was size-corrected  $^{22}$  $^{22}$  $^{22}$  by

$$
D_0 = \frac{\eta_{\rm W}}{\eta_{\rm water}} \left[ D_{\rm pbc}(L) + \frac{k_B T \zeta_{\rm ew} \alpha}{6\pi \eta L} \right]. \tag{S1}
$$

We explicitly took the different viscosities  $\eta_W$  of the different water models W into account (Table [S5\)](#page-4-3).

<span id="page-4-3"></span>Table S5: Viscosities  $\eta_W$  for each water model W. Values for SPC/E and TIP4P/2005 are taken from ref.,  $^{23}$  $^{23}$  $^{23}$  the value for TIP4P-D is taken from ref.,  $^{24}$  $^{24}$  $^{24}$  the other values are from their original publications  $(TIP3P-fb, <sup>25</sup> TIP4P-Ew<sup>26</sup>).$  $(TIP3P-fb, <sup>25</sup> TIP4P-Ew<sup>26</sup>).$  $(TIP3P-fb, <sup>25</sup> TIP4P-Ew<sup>26</sup>).$  $(TIP3P-fb, <sup>25</sup> TIP4P-Ew<sup>26</sup>).$  $(TIP3P-fb, <sup>25</sup> TIP4P-Ew<sup>26</sup>).$ 



#### <span id="page-4-1"></span>S1.5 Rate constant of water exchange

The most popular theory to calculate reaction rates is transition state theory (TST). [27](#page-21-2)[,28](#page-21-3) In simple systems for which the reaction coordinate is exactly known, TST gives an accurate estimate of the rate. However, in complex many body systems as the ones presented here, TST can fail due to the violation of the nonrecrossing hypothesis, one of the fundamentals of the theory. Therefore, in the following an alternative approach is presented (for more details see ref.  $6$ ).

For the exchange of water from the first hydration shell of  $Mg^{2+}$ , the rate constant k is defined by <sup>[13](#page-20-0)</sup>

$$
rate = 6 \cdot k \cdot [Mg(H_2O)62+], \qquad (S2)
$$

where 6 is the coordination number of the first hydration shell and  $[Mg(H_2O)_6^{2+}]$  is the concentration of hexa-coordinated  $\rm Mg^{2+}$  ions.

In this work, the water exchange rate constant  $k$  is calculated by counting the total number of transitions that are observed within a 1  $\mu$ s long trajectories of a 1 M MgCl<sub>2</sub> solution. As a transition N, we regard the exchange of waters between the first and second hydration shell of Magnesium (the exchange from first to second hydration shell and the reverse transition are counted as individual events). The water exchange rate constant  $k$  is hence given by

$$
k = \frac{1}{N_{\rm H_2O}} \cdot \frac{N}{2 \cdot t_{\rm B}} \tag{S3}
$$

where  $N_{\text{H}_2\text{O}}$  is the number of water molecules in the simulation box.  $t_B = N_{\text{Mg}} \cdot p_B \cdot t_{\text{sim}}$  is the cumulative time the water molecule spends in the first hydration shell of any  $Mg^{2+}$  ion.  $N_{Mg}$  is the number of  $Mg^{2+}$ ions in the simulation box,  $p_B = 6/(N_{\text{H}_2\text{O}}-6)$  is the probability of water to be in the first hydration shell and  $t_{\text{sim}}$  is the total simulation time. Errors are calculated from block averaging <sup>[29](#page-21-4)</sup> by dividing the trajectory into 2 blocks.

### <span id="page-5-0"></span>S1.6 Calculation of activity derivatives

The derivative  $a_{\rm cc}$  of the activity  $a_{\rm c} = \rho_{\rm c} y_{\rm c}$  with activity coefficient  $y_{\rm c}$  is defined via

$$
a_{\rm cc} = \left(\frac{\partial \ln a_{\rm c}}{\partial \ln \rho_{\rm c}}\right)_{\rm P,T} = 1 - \left(\frac{\partial \ln y_{\rm c}}{\partial \ln \rho_{\rm c}}\right)_{\rm P,T} = \frac{1}{1 + \rho_{\rm c}(G_{\rm cc} - G_{\rm co})},
$$
(S4)

with respect to the natural logarithm of the number density  $\rho_c$ . The expressions  $G_{cc}$  and  $G_{co}$  are for divalent cations obtained [30,](#page-21-5)[31](#page-21-6) from

$$
G_{\rm cc} = \frac{1}{9} \left[ G_{++} + 4(G_{--} + G_{+-}) \right]
$$
 (S5)

and

$$
G_{\rm co} = G_{\rm oc} = \frac{1}{3}G_{\rm +o} + \frac{2}{3}G_{\rm -o} \,,\tag{S6}
$$

where  $G_{ij}$  are the Kirkwood-Buff (KB) integrals  $31-33$  $31-33$  (eq [S7\)](#page-6-1) with +, −, and o denoting the cation, anion, and water oxygen, respectively. To obtain these KB integrals we followed that same strategy as in ref.,  $^6$  $^6$ 

<span id="page-6-1"></span>
$$
G_{ij} = 4\pi \int_0^\infty \left[ g_{ij}^{\mu \text{VT}}(r_{ij}) - 1 \right] r_{ij}^2 dr_{ij} , \qquad (S7)
$$

where  $g_{ij}^{\mu VT}(r_{ij})$  is the radial distribution function of the two species in the grand canonical ensemble, with  $r_{ij}$  being the center of mass distance between the two. Note that the simulation were done in the NPT ensemble. Therefore, we introduce a correction factor such that the radial distribution function used in the calculations of the KB integrals shows the correct asymptotic behavior at large distances (for more details see refs.  $^{1,6,31,33}$  $^{1,6,31,33}$  $^{1,6,31,33}$  $^{1,6,31,33}$  $^{1,6,31,33}$  $^{1,6,31,33}$ .

#### <span id="page-6-0"></span>S1.7 Calculation of free energy profiles

We computed all free energy profiles  $F(R)$  with umbrella sampling  $34,35$  $34,35$  and the weighted histogram analysis method (WHAM) [36](#page-21-10) to combine the individual umbrella windows.

1D  $Mg^{2+}$  - water: One dimensional free energy profiles as a function of the distance between  $Mg^{2+}$  and the oxygen atom of the leaving water molecule for the 12-6 type force fields were computed with GROMACS<sup>[37](#page-22-0)</sup> (version 2020). The profile of the 12-6-4 type force field in  $SPC/E$  water was computed using AMBER<sup>[38](#page-22-1)</sup> (version 2018) and PLUMED<sup>[39](#page-22-2)</sup> (version 2.5). Force constants and window spacing were in both setups  $k = 400,000 \text{ kJ/(mol nm}^2) \text{ and } 0.005 \text{ nm}$  [0.17  $\le R_{\text{MgOx}} < 0.4 \text{ nm}$ ] and  $k = 100,000 \text{ kJ/(mol nm}^2)$  and 0.01 nm  $[0.4 \leq R_{\rm MgOx} < 0.6$  nm]. Frames for the analysis were considered every 0.5 ps. We took a bin width of 5.4 x  $10^{-4}$  nm into account for WHAM.

To ensure convergence, windows of size 50 ns were obtained and subdivided into blocks of  $(A)$  10 x 5 ns, (B) 5 x 10 ns, and (C) 2 x 25 ns for the Mamatkulov-Netz parameters in SPC/E water (Figure [S1\)](#page-7-0). Of each individual block a free energy profile is calculated. Afterwards, we transformed the free energy profiles into potentials of mean force  $V^{PMF}(R)$  by applying a Jacobian correction,

<span id="page-6-2"></span>
$$
V^{\text{PMF}}(R) = F(R) + 2k_{\text{B}}T \ln R,\tag{S8}
$$

to take radial distances rather the Cartesian coordinates into account. We find that with windows of duration (A) 5 ns and (B) 10 ns the individual profiles show deviations whereas with a duration of (C) 25 ns the results are converged.

All other water exchange free energy profiles were thus simulated for 25 ns (see Table [S3](#page-3-1) and [S4](#page-4-2) for more

details on the simulation setups).

As an additional validation step, we obtained the radial distribution function  $g(R)$  from the 1  $\mu$ s 1 M MgCl<sup>2</sup> trajectory and the free energy profile via Boltzmann inversion. Note that the free energy profile from the radial distribution function is used only to validate the depth of the two minima. The barrier height cannot be derived from this method as the transitions of water exchange are too rare. Further insight into the quality of the umbrella sampling method is provided by the evenly spaced and overlapping histograms (Figure [S1D](#page-7-0)).

<span id="page-7-0"></span>

Figure S1: Convergence check. Potential of mean force  $V^{\text{PMF}}(R)$  along the distance between a  $Mg^{2+}$  ion and the oxygen atom of the leaving water molecule, obtained from windows of 50 ns that are subdivided into blocks of (A) 10 x 5 ns, (B) 5 x 10 ns, and (C) 2 x 25 ns and  $V^{PMF}(R)$  obtained from  $-\ln g(R)$ . In (D) additional to  $V^{PMF}(R)$  obtained from 25 ns long windows, the histograms are shown. Parameters used are Mamatkulov-Netz and SPC/E water.

1D Mg<sup>2+</sup> - phosphate oxygen: The free energy profile along the distance between Mg<sup>2+</sup> and one of the two non-bridging phosphate oxygens of the dimethylphosphate (DMP) was obtained with force constants and window spacing of  $k = 300,000 \text{ kJ/(mol nm}^2)$  and  $0.0075 \text{ nm}$   $[0.15 \leq R_{\rm MgOP} < 0.525 \text{ nm}]$  and  $k = 5,000$  $\mathrm{kJ/(mol\;nm^2)}$  and 0.02 nm  $[0.525\leq R_{\rm MgOP}< 0.885$  nm], respectively, using GROMACS<sup>[40](#page-22-3)</sup> (version 2020). PLUMED<sup>[39](#page-22-2)</sup> was employed to include an additional bias avoiding any direct interaction with any of the DMP atoms but the selected phosphate oxygen (see ref.<sup>[6](#page-19-6)</sup> for more details). Frames for the analysis were considered every 0.5 ps. For WHAM a bin width of  $9.2 \times 10^{-4}$  nm was considered.

To ensure convergence, windows of size 100 ns were obtained and subdivided into blocks of  $(A)$  10 x 10

ns, (B) 5 x 20 ns, and (C) 2 x 50 ns for the Mamatkulov-Netz SPC/E parameters (Figure [S2\)](#page-8-1). The results show that 20 ns per window are sufficient to provide converged results. For all other profiles for DMP and  $Mg^{2+}$ , windows were simulated for 20 ns (see Table [S3](#page-3-1) and [S4](#page-4-2) for more details on the simulation setups). Additionally, Figure [S2D](#page-8-1) shows the overlapping histograms of neighboring simulation windows.

<span id="page-8-1"></span>

Figure S2: Convergence check. Free energy profile  $F(R)$  along the distance between one of the non-bridging oxygens of the DMP and a  $Mg^{2+}$  ion, obtained from windows of 100 ns that are subdivided into blocks of (A) 10 x 10 ns, (B)  $5 \times 20$  ns, and (C)  $2 \times 50$  ns. In (D) additional to  $F(R)$  obtained from 20 ns long windows, overlapping distance histograms are shown. Parameters used are Mamatkulov-Netz and SPC/E water.

### <span id="page-8-0"></span>S1.8 Binding affinity towards DMP

Calculating binding affinities from molecular dynamics simulations can be challenging since the results depend sensitively on accurate sampling. In principle, one can obtain the binding affinity of a metal cation toward the phosphate oxygen in three different ways: (i) Via association/dissociation rates, (ii) via free energy perturbation calculations, or (iii) via integrating potentials of mean force (PMFs). All three methods yield consistent results if accurately sampled (see Table S6 for results for  $Ca^{2+}$  from previous works). To avoid possible pitfalls, we strongly recommend to use different methods to verify binding affinities obtained by computer simulations.

#### <span id="page-9-0"></span>S1.8.1 General strategies for computing binding affinities

**Association/dissociation rates** If association  $k_{on}$  and dissociation  $k_{off}$  rates are available, the free energy of binding can be calculated from

<span id="page-9-4"></span>
$$
\Delta G_{\rm b}^0 = -k_{\rm B}T \cdot \ln\left(\frac{c^0}{[M]} \cdot \frac{k_{\rm on}}{k_{\rm off}}\right),\tag{S9}
$$

with respect to the standard concentration  $c^0 = 1$  M. For  $Ca^{2+}$ , the rates were calculated in our previous work<sup>[41](#page-22-4)</sup> and the results for the resulting free energy are shown in Table [S6.](#page-10-1)

Free energy perturbation The binding free energy can be obtained from free energy perturbation using the double decoupling method  $(DDM)^{42}$  $(DDM)^{42}$  $(DDM)^{42}$  or alchemical transformation calculations.  $43,44$  $43,44$  In the former (eq [S10\)](#page-9-1), both van der Waals as well as electrostatic interactions are decoupled for the ion in solution  $(\Delta G_{\text{elec}+\text{vdW}}^{\text{solv}})$  and turned back on for an ion at the binding site  $(\Delta G_{\text{elec}+\text{vdW}}^{\text{BS}})$ . The effect of the restraints to keep the ion in the binding site during the decoupling in bulk  $(\Delta G_{\text{rest}}^{\text{solv}})$  and in the binding site  $(\Delta G_{\text{rest}}^{\text{BS}})$ has to be included [42](#page-22-5)

<span id="page-9-1"></span>
$$
\Delta G_{\rm b}^0 = \Delta G_{\rm elec+vdW}^{\rm solv} + \Delta G_{\rm elec+vdW}^{\rm BS} + \Delta G_{\rm rest}^{\rm solv} + \Delta G_{\rm rest}^{\rm BS}
$$
\n(S10)

In the alchemical transformation (eq [S11\)](#page-9-2), a reference ion of the same valency is used. The ion of interest is alchemically transformed into the reference ion, both in bulk  $(\Delta G_{\text{ion}\rightarrow\text{ref}}^{\text{solv}})$  as well as at the binding site  $(\Delta G_{\text{ref}\rightarrow \text{ion}}^{\text{B}})$ 

<span id="page-9-2"></span>
$$
\Delta G_{\rm b}^0 = \Delta G_{\rm ion \to ref}^{\rm solv} + \Delta G_{\rm ref}^0 + \Delta G_{\rm ref \to ion}^{\rm BS} \tag{S11}
$$

where  $\Delta G_{\rm ref}^0$  is the binding affinity of the reference ion. To check the convergence, the transformation is done in opposite directions (i.e. forward and backward). For  $Ca^{2+}$ , the values for both methods were obtained from our previous works [3](#page-19-3)[,41](#page-22-4) and are listed in Table [S6.](#page-10-1)

Integration of a potential of mean force Here, the binding affinity is obtained by integrating the potential of mean force  $V^{PMF}(r)$  (eq [S8\)](#page-6-2) along the distance r between the binding site and the ion of interest,

<span id="page-9-3"></span>
$$
\Delta G_{\rm b}^{0} = -k_{\rm B}T \cdot \ln \left( \frac{c^0}{[M]} \cdot \frac{\int_0^{\rm T^{\dagger}} r^2 e^{-V^{\rm PMF}(r)/k_{\rm B}T} dr}{\int_{\rm T^{\dagger}}^{\rm T_{\rm L}} r^2 e^{-V^{\rm PMF}(r)/k_{\rm B}T} dr} \right) , \qquad (S12)
$$

where  $[M]$  is the ion concentration of the simulation box,  $r<sub>L</sub>$  is the radius of a sphere that contains the same number of water molecules as the simulation box and  $r^{\dagger}$  is the position of the maximum of the PMF. Note that both  $[M]$  and  $r<sub>L</sub>$  are dependent on the number of waters in the simulation such that eq [S12](#page-9-3) becomes

independent of the box size used. The value for  $Ca^{2+}$  listed in Table [S6](#page-10-1) is taken from our previous work. <sup>[6](#page-19-6)</sup>

As shown in Table S6 the results for  $Ca^{2+}$  for all four methods match within error. This demonstrates that the methods are converged and that all of these setups to obtain the binding free energy for  $Ca^{2+}$ (which is used as reference ion in this current work) are sufficiently accurate.

<span id="page-10-1"></span>Table S6: Comparison of different methods to compute the free binding energy of  $Ca^{2+}$  to a non-bridging

phosphate oxygen using the parameters by Mamatkulov-Schwierz in TIP3P. [1](#page-19-1) Method  $\overline{C}^0$  $Ca<sup>2+</sup>$ 

11100110U	----------	$ C_{\rm a}$ <sup>2+</sup> $[k_B T]$
Rates (TREMD)	(S9)	$-4.51 \pm 1.27$ <sup>41</sup>
FEP DDM	(S10)	$-4.57 \pm 0.90$ <sup>41</sup>
FEP alchem. trafo.	(S11)	$-4.86 \pm 0.85$ <sup>3</sup>
PMF	(S12)	$-4.84 \pm 0.18$ <sup>6</sup>

#### <span id="page-10-0"></span>S1.8.2 Methods in this work for computing binding affinities

In this present work, to fine-tune the interaction and find scaling factors between the different  $Mg^{2+}$  models and the phosphate oxygen, we primarily computed the binding affinities via integration of PMFs (eq [S12\)](#page-9-3), as this strategy also includes a more mechanistic picture due to the obtained energetic profiles. The free energy profiles (most of which are shown in Figure 6B,C in the main text) were obtained using umbrella sampling. Details on all the simulation parameters are given in Section S1.2 and all parameters for the umbrella sampling in Section S1.6. (All simulation parameters are exactly the same as in our previous work before.<sup>[6](#page-19-6)</sup>)

The concentrations of  $Mg^{2+}$  in the simulation boxes used in eq [S12](#page-9-3) are  $[M] = 0.037$  M and 0.036 M for the 3-site and 4-site waters, respectively. Errors are calculated by dividing the 20 ns long windows into 4 blocks, and subsequent calculation of PMFs and thus  $\Delta G_{\rm b}^0$  from each block and block averaging.

As an independent validation step, we selected a second method to obtain the binding affinities. Association/dissociation rates are very difficult to obtain for  $Mg^{2+}$  due to the long timescales involved. Therefore, we chose alchemical transformation calculations (eq [S11\)](#page-9-2), as their accuracy is comparable to other methods (Table [S6\)](#page-10-1), while they are computationally more efficient than the double decoupling method (eq [S10\)](#page-9-1) as fewer steps are involved. Employing the final scaling factors (Table 1, main text), we calculate  $Mg^{2+}$ binding affinities toward the DMP by transforming the  $Mg^{2+}$  ion for each water models into a reference ion (eq [S11\)](#page-9-2) and using forward and backward transformations. For the reference ion, we chose the parameters by Mamatkulov-Schwierz for  $Ca^{2+}$  in TIP3P water<sup>[1](#page-19-1)</sup> since we already had well converged results based on our previous work (Table [S6\)](#page-10-1).

The transformation was performed in 20 iterative steps gradually switching the Lennard-Jones parameters from the respective  $Mg^{2+}$  model to the reference ion and vice versa both in solution as well as at its binding site (yielding  $\Delta G_{\text{ion}\rightarrow \text{ref}}^{\text{solv}}$  and  $\Delta G_{\text{ref}\rightarrow \text{ion}}^{\text{BS}}$  from eq [S11\)](#page-9-2).

 $\Delta G_{\text{ref}}^0$  for each water model was obtained from integration of the potentials of mean force (Figure [S8\)](#page-17-1) using eq [S12](#page-9-3) (simulation parameters are again given in Section S1.2 and Section S1.6) and the values are listed in Table [S8.](#page-17-2)

Employing the respective reference binding affinities, we check for convergence by performing the transformation for each  $Mg^{2+}$  model corresponding to one of the five waters both in forward as well as in backward direction (Tables [S9](#page-18-0) and [S10\)](#page-18-1).

Both methods yielded identical results within error (Tables [S8,](#page-17-2) [S9,](#page-18-0) and [S10\)](#page-18-1). Errors were calculated from block averaging by dividing the trajectory of the alchemical transformation into 5 blocks.

# <span id="page-11-0"></span>S2 Supplementary results

#### <span id="page-11-1"></span>S2.1 Transferability

<span id="page-11-2"></span>Table S7: Parameters and single-ion properties after converting  $microMg(TIP3P)$  and  $nanoMg(TIP3P)$ into  $microMq(TIP3P|W)$  and  $nanoMq(TIP3P|W)$  using Lorentz-Berthelot combination rules (eq 2, main text, unscaled).  $\sigma_{ii}$ ,  $\varepsilon_{ii}$ ,  $\sigma_{io}$ ,  $\varepsilon_{io}$  are the ion-ion and ion-water LJ parameters single-ion properties.  $\Delta G_{\text{solv}}$ ,  $R_1$ , and  $n_1$  are the solvation free energy of the neutral MgCl<sub>2</sub> ion-pair, the radius of the first hydration shell, and the coordination number of the first hydration shell, respectively.

	$\sigma_{\rm ii}$	$\varepsilon$ <sub>ii</sub>	$\sigma_{\rm io}$	$\varepsilon_{\rm io}$	$\Delta G_{\rm solv}$	$R_1$	n <sub>1</sub>
	nm	[kJ/mol]	$\vert$ nm $\vert$	[kJ/mol]	kJ/mol	$\left  \text{nm} \right $	
microMg(TIP3P SPC/E)	0.1019	235.80	0.2101	13.75	$-2534.9 \pm 1$	$0.206 \pm 0.004$	6
nanoMg(TIP3P SPC/E)	0.1025	389.80	0.2106	17.50	$-2535.0 \pm 1$	$0.210 \pm 0.004$	6
$microMg(TIP3P TIP3P-fb)$	0.1019	235.80	0.2107	13.77	$-2536.1 \pm 1$	$0.208 \pm 0.004$	6
$nanoMg(TIP3P TIP3P-fb)$	0.1025	389.80	0.2112	17.52	$-2535.5 \pm 1$	$0.211 \pm 0.004$	6
$microMg$ (TIP3P TIP4P/2005)	0.1019	235.80	0.2097	15.01	$-2451.2 \pm 1$	$0.210 \pm 0.004$	6
$nanoMq$ (TIP3P TIP4P/2005)	0.1025	389.80	0.2102	19.10	$-2455.8 \pm 1$	$0.213 \pm 0.004$	6
$microMg(TIP3P TIP4P-D)$	0.1019	235.80	0.2101	14.49	$-2474.7 \pm 1$	$0.211 \pm 0.004$	6
$nanoMg(TIP3P TIP4P-D)$	0.1025	389.80	0.2106	20.99	$-2483.1 \pm 1$	$0.214 \pm 0.004$	6
$microMg$ (TIP3P TIP4P-Ew)	0.1019	235.80	0.2100	14.07	$-2462.9 \pm 1$	$0.208 \pm 0.004$	6
$nanoMg(TIP3P TIP4P-Ew)$	0.1025	389.80	0.2105	17.90	$-2465.4 \pm 1$	$0.212 \pm 0.004$	6
exp.					$-2532^{12}$	$0.209 \pm 0.004^{11}$	$6^{11}$

<span id="page-12-0"></span>

Figure S3: Solvation free energy  $\Delta G_{\text{solv}}$ , radius of the first hydration shell  $R_1$ , and coordination number  $n_1$ isosurfaces in  $\sigma_{io}$  -  $\varepsilon_{io}$  space for the interaction of the *microMg* and *nanoMg* parameters of various water models, including the force fields for (A)  $TIP3P^{45}$  $TIP3P^{45}$  $TIP3P^{45}$  from the literature,  $6$  (B)  $SPC/E$ ,  $^{46}$  $^{46}$  $^{46}$  (C)  $TIP3P$ -fb,  $^{25}$  $^{25}$  $^{25}$  (D) TIP4P/2005,<sup>[47](#page-22-10)</sup> (E) TIP4P-Ew,<sup>[26](#page-21-1)</sup> and (F) TIP4P-D.<sup>[48](#page-23-0)</sup> Transparent circles indicate converted  $microMq(TIP3P|W)$  and  $nanoMq(TIP3P|W)$  parameters, using the ion-ion parameter sets,  $\sigma_{ii}$  and  $\varepsilon_{ii}$ , of TIP3P water [6](#page-19-6) and Lorentz-Berthelot combination rules (eq 2, main text, unscaled) to gain the effective water-ion parameter sets,  $\sigma_{io}$  and  $\varepsilon_{io}$ , of the respective water model W (Table [S7\)](#page-11-2). In parameter regions far off the  $\Delta G_{\text{solv}}$  isolines, grid points are sparser. Therefore, the  $R_1$  and  $n_1$  surfaces are less accurate in those regions.



<span id="page-13-0"></span>S2.3 Lennard-Jones interaction potentials

Figure S4: Lennard-Jones interaction potential  $V^{\rm LJ}$  as function of the Mg<sup>2+</sup>-oxygen distance  $r_{\rm MgOw}$  for different Mg<sup>2+</sup> force fields and (A) TIP3P-fb,<sup>[25](#page-21-0)</sup> (B) TIP4P/2005,<sup>[47](#page-22-10)</sup> (C) TIP4P-Ew,<sup>[26](#page-21-1)</sup> and (D) TIP4P-D<sup>[48](#page-23-0)</sup> water models.



<span id="page-14-0"></span>S2.4 One-dimensional free energy profiles for  $Mg^{2+}$ -water interactions

Figure S5: One-dimensional free energy profiles as a function of the distance between  $Mg^{2+}$  and the leaving water molecule  $R_{\text{Mg-Ox}}$  for different force fields in combination with (A) TIP3P-fb,<sup>[25](#page-21-0)</sup> (B) TIP4P/2005,<sup>[47](#page-22-10)</sup> (C) TIP4P-Ew,  $^{26}$  $^{26}$  $^{26}$  and (D) TIP4P-D<sup>[48](#page-23-0)</sup> water models.



<span id="page-15-0"></span>S2.5 Two-dimensional free energy profiles for  $Mg^{2+}$ -water interactions

Figure S6: Two-dimensional free energy profiles as a function of the distance between  $Mg^{2+}$  and the leaving water molecule  $r_{\text{MgOw}_1}$  and  $\text{Mg}^{2+}$  and the entering water molecule  $r_{\text{MgOw}_2}$  for different force fields in combination with  $(A,B)$  SPC/E and  $(C,D)$  TIP3P-fb water models. The free energy profiles were calculated via umbrella sampling using additional restraints (see ref. [6](#page-19-6) for more details).



Figure S7: Two-dimensional free energy profiles as a function of the distances between  $Mg^{2+}$  and the leaving water molecule  $r_{\text{MgOw}_1}$  and  $\text{Mg}^{2+}$  and the entering water molecule  $r_{\text{MgOw}_2}$  for different force fields in combination with (A,B) TIP4P-Ew, (C,D) TIP4P-D and (E,F) TIP4P/2005 water models. The free energy profiles were calculated via umbrella sampling using additional restraints (see ref. <sup>[6](#page-19-6)</sup> for more details).

#### <span id="page-17-0"></span>S2.6 Binding affinities

Two strategies were proceeded to validate binding affinities values, as the calculation of binding affinities is prone to errors even for small systems, like the one we are using here. Both methods yielded nearly identical results within error (Tables 5, main text, [S9,](#page-18-0) and [S10\)](#page-18-1).

First strategy to calculate binding affinities: Integration of the free energy profiles along the distance between the  $Mg^{2+}$  ion and the dimethylphosphate (DMP) molecule. Resulting properties are listed in Table 5 in the main text.

Second strategy to calculate binding affinities: Alchemical transformation between  $Mg^{2+}$  and an divalent reference atom of known binding affinity. As reference ion (Table [S8\)](#page-17-2), we choose the parameters of  $Ca^{2+}$  in TIP3P water obtained earlier.<sup>[1](#page-19-1)</sup>

<span id="page-17-1"></span>

Figure S8: One-dimensional potential of mean force  $V^{\text{PMF}}$  as a function of the distance  $R_{\text{OP-ref}}$  between the divalent reference ion and one of the non-bridging phosphate oxygens of the DMP in different water models. The free energy profile was calculated via umbrella sampling using additional restraints (see ref. [6](#page-19-6) for more details). Afterwards a Jacobian correction was applied (eq [S8\)](#page-6-2).

<span id="page-17-2"></span>Table S8: Binding affinity  $\Delta G_{\text{ref}}^0$  obtained for the divalent reference ion (parameters for Ca<sup>2+</sup> from Mamatkulov-Schwierz in TIP3P<sup>[1](#page-19-1)</sup>) in the different water models via integration of their potentials of mean force (Figure [S8\)](#page-17-1).



<span id="page-18-0"></span>Table S9: Binding affinity  $\Delta G_{\rm b}^0$  obtained from forward alchemical transformation. The values given for  $\Delta G_{\text{Mg}^2+\rightarrow \text{ref}}^{\text{ol}}$  and  $\Delta G_{\text{ref}\rightarrow \text{Mg}^2+\text{}}^{\text{bind}}$  are obtained from block averaging for 5 blocks of 3 ns long windows each.  $\Delta G_{\text{ref}\to\text{Mg}^{2+}}^{\text{bind}}$ ,  $\Delta G_{\text{Mg}^{2+}\to \text{ref}}^{\text{bind}}$  and hence  $\Delta G_{\text{b}}^0$  are shown for the final scaling factor combination  $\lambda_{\sigma,\varepsilon}^{\text{RNA}}$ (Table 2, main text).  $R<sub>b</sub>$  is obtained from the last 15 ns window of the alchemical transformation calculation, the error here indicates the standard deviation of the distribution.

	$\Delta G_{\rm b}^0$	$\Delta G_{\rm Mg^{2+}\rightarrow ref}^{\rm solv}$	$\Delta G^{\rm bind}_{\rm ref \rightarrow Mg^{2+}}$	$R_{\rm b}$
	$[k_B T]$	$[k_B T]$	$[k_B T]$	${\rm [nm]}$
microMq(SPC/E)	$-0.873 \pm 0.4$	$131.776 \pm 0.003$	$-128.088 \pm 0.003$	$0.209 \pm 0.004$
nanoMq(SPC/E)	$-1.033 \pm 0.4$	$131.766 \pm 0.003$	$-128.238 + 0.003$	$0.208 \pm 0.004$
$microMg(TIP3P-fb)$	$-0.789 \pm 0.4$	$133.557 \pm 0.006$	$-130.465 \pm 0.003$	$0.208 \pm 0.004$
$nanoMg(TIP3P-fb)$	$-0.976 \pm 0.4$	$133.674 \pm 0.006$	$-130.770 \pm 0.006$	$0.208 \pm 0.004$
microMg(TIP4P/2005)	$0.138 \pm 0.4$	$161.893 \pm 0.008$	$-154.165 \pm 0.007$	$0.207 \pm 0.004$
nanoMg(TIP4P/2005)	$-0.632 \pm 0.4$	$161.396 \pm 0.007$	$-154.437 + 0.007$	$0.207 \pm 0.004$
$microMg(TIP4P-Ew)$	$-0.897 \pm 0.4$	$154.691 \pm 0.007$	$-148.532 \pm 0.007$	$0.208 \pm 0.004$
$nanoMg(TIP4P-Ew)$	$-0.945 \pm 0.4$	$154.497 \pm 0.003$	$-148.384 + 0.006$	$0.208 \pm 0.004$
$microMq(TIP4P-D)$	$-0.677 \pm 0.4$	$158.560 \pm 0.007$	$-152.535 \pm 0.006$	$0.207 \pm 0.004$
$nanoMg(TIP4P-D)$	$-0.775 \pm 0.4$	$158.550 \pm 0.006$	$-152.622 \pm 0.006$	$0.207 \pm 0.004$
exp.	$-1.036^{15}$	n.a.	n.a.	$0.206 - 0.208^{49}$

<span id="page-18-1"></span>Table S10: Binding affinity  $\Delta G_b^0$  obtained from backward alchemical transformation. The values given for  $\Delta G_{\text{ref}\to\text{Mg}^{2+}}^{\text{olub}}$  and  $\Delta G_{\text{Mg}^{2+}\to\text{ref}}^{\text{bind}}$  are obtained from block averaging for 5 blocks of 3 ns long windows each.  $\Delta G_{\text{Mg}^{2+}\to \text{ref}}^{\text{bind}}$  and hence  $\Delta G_{\text{b}}^{0}$  are shown for the final scaling factor combination  $\lambda_{\sigma,\varepsilon}^{\text{RNA}}$  (Table 2, main text).  $R<sub>b</sub>$  is obtained from the first 15 ns window of the alchemical transformation calculation, the error here indicates the standard deviation of the distribution.

	$\Delta G_{\rm b}^0$	$\Delta G_{\rm ref \rightarrow Mg^{2+}}^{\rm solv}$	$\Delta G_{\rm Mg^{2+}\rightarrow ref}^{\rm bind}$	$R_{\rm b}$
	$[k_B T]$	$[k_B T]$	$[k_B T]$	[nm]
microMq(SPC/E)	$-0.987 \pm 0.4$	$-131.606 \pm 0.003$	$128.032 \pm 0.003$	$0.209 \pm 0.004$
nanoMq(SPC/E)	$-1.097 \pm 0.4$	$-131.867 \pm 0.003$	$128.403 \pm 0.003$	$0.208 \pm 0.004$
$microMg(TIP3P-fb)$	$-0.953 \pm 0.4$	$-133.431 \pm 0.006$	$130.503 \pm 0.003$	$0.208 \pm 0.004$
$nanoMg(TIP3P-fb)$	$-0.821 \pm 0.4$	$-133.687 \pm 0.006$	$130.627 \pm 0.003$	$0.209 \pm 0.004$
microMg(TIP4P/2005)	$-0.435 \pm 0.4$	$-161.357 \pm 0.008$	$154.200 \pm 0.007$	$0.207 \pm 0.004$
nanoMg(TIP4P/2005)	$-0.974 \pm 0.4$	$-161.162 + 0.007$	$154.544 \pm 0.007$	$0.207 \pm 0.004$
$microMg(TIP4P-Ew)$	$-0.770 \pm 0.4$	$-154.764 \pm 0.003$	$148.478 \pm 0.003$	$0.208 \pm 0.004$
$nanoMg(TIP4P-Ew)$	$-0.589 \pm 0.4$	$-154.582 \pm 0.003$	$148.114 \pm 0.006$	$0.208 \pm 0.004$
$microMg(TIP4P-D)$	$-0.444 \pm 0.4$	$-158.512 \pm 0.007$	$152.254 \pm 0.006$	$0.207 \pm 0.004$
$nanoMg(TIP4P-D)$	$-0.687 \pm 0.4$	$-158.678 \pm 0.006$	$152.662 \pm 0.006$	$0.207 \pm 0.004$
exp.	$-1.036^{15}$	n.a.	n.a.	$0.206 - 0.208$ <sup>49</sup>

# <span id="page-19-0"></span>S3 Bibliography

- <span id="page-19-1"></span>(1) Mamatkulov, S.; Schwierz, N. Force fields for monovalent and divalent metal cations in TIP3P water based on thermodynamic and kinetic properties. J. Chem. Phys. 2018, 148, 74504.
- <span id="page-19-2"></span>(2) Schwierz, N. Kinetic pathways of water exchange in the first hydration shell of magnesium. J. Chem. Phys. 2020, 152, 224106.
- <span id="page-19-3"></span>(3) Cruz-León, S.; Grotz, K. K.; Schwierz, N. Extended magnesium and calcium force field parameters for accurate ion-nucleic acid interactions in biomolecular simulations. J. Chem. Phys. 2021, 154, 171102.
- <span id="page-19-4"></span>(4) Nilsson, L.; Villa, A. Magnesium Ion Water Coordination and Exchange in Biomolecular Simulations. 2012,
- <span id="page-19-5"></span>(5) Panteva, M. T.; Giambasu, G. M.; York, D. M. Comparison of structural, thermodynamic, kinetic and mass transport properties of  $Mg^{2+}$  ion models commonly used in biomolecular simulations. J. Comput. Chem. 2015, 36, 970–982.
- <span id="page-19-6"></span>(6) Grotz, K. K.; Cruz-León, S.; Schwierz, N. Optimized Magnesium Force Field Parameters for Biomolecular Simulations with Accurate Solvation, Ion-Binding, and Water-Exchange Properties. J. Chem. Theory Comput. 2021, 17, 2530–2540.
- <span id="page-19-7"></span>(7) Li, P.; Roberts, B. P.; Chakravorty, D. K.; Merz, K. M. Rational Design of Particle Mesh Ewald Compatible Lennard-Jones Parameters for  $+2$  Metal Cations in Explicit Solvent. J. Chem. Theory Comput. 2013, 9, 2733–2748.
- <span id="page-19-8"></span>(8) Li, P.; Merz, K. M. Taking into account the ion-induced dipole interaction in the nonbonded model of ions. J. Chem. Theory Comput. 2014, 10, 289–297.
- <span id="page-19-9"></span>(9) Mamatkulov, S.; Fyta, M.; Netz, R. R. Force fields for divalent cations based on single-ion and ion-pair properties. J. Chem. Phys. 2013, 138, 24505.
- <span id="page-19-10"></span>(10) Åqvist, J. Ion-water interaction potentials derived from free energy perturbation simulations. J. Phys. Chem. 1990, 94, 8021–8024.
- <span id="page-19-11"></span>(11) Marcus, Y. Ionic Radii in Aqueous solutions. Chem. Rev. 1988, 88, 1475–1498.
- <span id="page-19-12"></span>(12) Marcus, Y. Ion Properties; Marcel Dekker, Inc.: New York, Basel, 1997.
- <span id="page-20-0"></span>(13) Neely, J.; Connick, R. Rate of Water Exchange from Hydrated Magnesium Ion. J. Am. Chem. Soc. 1970, 92, 3476–3478.
- <span id="page-20-1"></span>(14) Bleuzen, A.; Pittet, P.-A.; Helm, L.; Merbach, A. E. Water exchange on magnesium(II) in aqueous solution: a variable temperature and pressure  $^{17}$ O NMR study. *Magn. Reson. Chem.* 1997, 35, 765– 773.
- <span id="page-20-2"></span>(15) Sigel, R. K.; Sigel, H. A stability concept for metal ion coordination to single-stranded nucleic acids and affinities of individual sites. Acc. Chem. Res. 2010, 43, 974–984.
- <span id="page-20-3"></span>(16) Darden, T.; Pearlman, D.; Pedersen, L. G. Ionic charging free energies: Spherical versus periodic boundary conditions. J. Chem. Phys. 1998, 109, 10921–10935.
- <span id="page-20-4"></span>(17) Ryckaert, J. P.; Ciccotti, G.; Berendsen, H. J. Numerical integration of the cartesian equations of motion of a system with constraints: molecular dynamics of n-alkanes. J. Comp. Phys. 1977, 23, 327–341.
- <span id="page-20-5"></span>(18) Hess, B.; Bekker, H.; Berendsen, H. J. C.; Fraaije, J. G. E. M. LINCS: a linear constraint solver for molecular simulations. J. Comput. Chem. 1997, 18, 1463-1472.
- <span id="page-20-6"></span>(19) Berendsen, H. J.; Postma, J. P.; Van Gunsteren, W. F.; Dinola, A.; Haak, J. R. Molecular dynamics with coupling to an external bath. J. Chem. Phys.  $1984$ ,  $81$ ,  $3684-3690$ .
- <span id="page-20-7"></span>(20) Bussi, G.; Donadio, D.; Parrinello, M. Canonical sampling through velocity rescaling. J. of Chem. Phys. 2007, 126 .
- <span id="page-20-8"></span>(21) Parrinello, M.; Rahman, A. Polymorphic transitions in single crystals: A new molecular dynamics method. J. Appl. Phys. 1981, 52, 7182–7190.
- <span id="page-20-9"></span>(22) Yeh, I.-C.; Hummer, G. System-Size Dependence of Diffusion Coefficients and Viscosities from Molecular Dynamics Simulations with Periodic Boundary Conditions. J. Phys. Chem. B 2004, 108, 15873– 15879.
- <span id="page-20-10"></span>(23) Vega, C.; Abascal, J. L. Simulating water with rigid non-polarizable models: A general perspective. Phys. Chem. Chem. Phys. 2011, 13, 19663-19688.
- <span id="page-20-11"></span>(24) von Bülow, S.; Siggel, M.; Linke, M.; Hummer, G. Dynamic cluster formation determines viscosity and diffusion in dense protein solutions. Proc. Natl. Acad. Sci. 2019, 116, 9843-9852.
- <span id="page-21-0"></span>(25) Wang, L. P.; Martinez, T. J.; Pande, V. S. Building force fields: An automatic, systematic, and reproducible approach. J. Phys. Chem. Lett. 2014, 5, 1885–1891.
- <span id="page-21-1"></span>(26) Horn, H. W.; Swope, W. C.; Pitera, J. W.; Madura, J. D.; Dick, T. J.; Hura, G. L.; Head-Gordon, T. Development of an improved four-site water model for biomolecular simulations: TIP4P-Ew. J. Chem. Phys. 2004, 120, 9665-9678.
- <span id="page-21-2"></span>(27) Wynne-Jones, W. F. K.; Eyring, H. The Absolute Rate of Reactions in Condensed Phases. J. Chem. Phys. 1935, 3, 492-502.
- <span id="page-21-3"></span>(28) Wigner, E. The Transition State Method. Trans. Faraday Soc. 1937, 29–41.
- <span id="page-21-4"></span>(29) Frenkel, D.; Smit, B. Understanding Molecular Simulation: From Algorithms to Applications, 2nd ed.; Elsevier Science, 2001.
- <span id="page-21-5"></span>(30) Weerasinghe, S.; Smith, P. E. A Kirkwood-Buff derived force field for sodium chloride in water. J. Chem. Phys. 2003, 119, 11342–11349.
- <span id="page-21-6"></span>(31) Fyta, M.; Netz, R. R. Ionic force field optimization based on single-ion and ion-pair solvation properties: Going beyond standard mixing rules. J. Chem. Phys. 2012, 136.
- (32) Kirkwood, J. G.; Buff, F. P. The Statistical Mechanical Theory of Solutions. I. J. Chem. Phys. 1951, 19, 774–777.
- <span id="page-21-7"></span>(33) Lee Warren, G.; Patel, S. Hydration free energies of monovalent ions in transferable intermolecular potential four point fluctuating charge water: An assessment of simulation methodology and force field performance and transferability. J. Chem. Phys. 2007, 127 .
- <span id="page-21-8"></span>(34) Torrie, G. M.; Valleau, J. P. Monte Carlo free energy estimates using non-Boltzmann sampling: Application to the sub-critical Lennard-Jones fluid. Chem. Phys. Lett. 1974, 28, 578–581.
- <span id="page-21-9"></span>(35) and Torrie, G.; Valleau, J. Nonphysical sampling distributions in Monte Carlo free-energy estimation: Umbrella sampling. J. Comput. Phys. 1977, 23, 187.
- <span id="page-21-10"></span>(36) Kumar, S.; Rosenberg, J. M.; Bouzida, D.; Swendsen, R. H.; Kollman, P. A. The weighted histogram analysis method for free-energy calculations on biomolecules. I. The method. J. Comput. Chem. 1992, 13, 1011–1021.
- <span id="page-22-0"></span>(37) Hess, B. P-LINCS: A Parallel Linear Constraint Solver for Molecular Simulation. J. Chem. Theory Comput. 2008, 4, 116–122.
- <span id="page-22-1"></span>(38) Case, D. A.; Belfon, K.; Ben-Shalom, I. Y.; Brozell, S. R.; Cerutti, D. S.; Cheatham, T. E. I.; Cruzeiro, V. W. D.; Darden, T.; Duke, R. E.; Giambasu, G.; Gilson, M. K.; Gohlke, H.; Goetz, A. W.; Harris, R.; Izadi, P. A.; Izmailov, S.; Kasavajhala, K.; Kovalenko, A.; Krasny, R.; Kurtzman, T.; Lee, T. S.; LeGrand, S.; Li, P.; Lin, C.; Liu, J.; Luchko, T.; Luo, R.; Man, V.; Merz, K. M.; Miao, Y.; Mikhailovskii, O.; Monard, G.; Nguyen, H.; Onufriev, A.; Pan, F.; Pantano, S.; Qi, R.; Roe, D. R.; Roitberg, A.; Sagui, C.; Schott-Verdugo, S.; Shen, J.; Simmerling, C. L.; Skrynnikov, N. R.; Smith, J.; Swails, J.; Walker, R. C.; Wang, J.; Wilson, L.; Wolf, R. M.; Wu, X.; Xiong, Y.; Xue, Y.; York, D. M.; Kollman, P. A. Amber 2018. 2018; <https://ambermd.org/AmberTools.php>.
- <span id="page-22-2"></span>(39) Tribello, G. A.; Bonomi, M.; Branduardi, D.; Camilloni, C.; Bussi, G. Plumed 2: New feathers for an old bird. Comput. Phys. Commun. 2014, 185, 604–613.
- <span id="page-22-3"></span>(40) Hess, B.; van der Vegt, N. F. A. Cation specific binding with protein surface charges. Proc. Natl. Acad. Sci. 2009, 106, 13296–13300.
- <span id="page-22-4"></span>(41) Cruz-León, S.; Schwierz, N. Hofmeister Series for Metal-Cation-RNA Interactions: The Interplay of Binding Affinity and Exchange Kinetics. Langmuir 2020, 36, 5979–5989.
- <span id="page-22-5"></span>(42) Gilson, M. K.; Given, J. A.; Bush, B. L.; McCammon, J. A. The statistical-thermodynamic basis for computation of binding affinities: a critical review. *Biophys. J.* **1997**, 72, 1047–1069.
- <span id="page-22-6"></span>(43) Chipot, C. Free energy calculations; Springer Science Business Media: Heidelberg, Berlin, 2007.
- <span id="page-22-7"></span>(44) Cruz-León, S.; Vanderlinden, W.; Müller, P.; Forster, T.; Staudt, G.; Lin, Y.; Lipfert, J.; Schwierz, N. Twisting DNA by Salt.  $bioRxiv2021.07.14.452306$  **2021**, 21-26.
- <span id="page-22-8"></span>(45) Jorgensen, W. L.; Chandrasekhar, J.; Madura, J. D.; Impey, R. W.; Klein, M. L. Comparison of simple potential functions for simulating liquid water. J. Chem. Phys. 1983, 79, 926–935.
- <span id="page-22-9"></span>(46) Berendsen, H. J.; Grigera, J. R.; Straatsma, T. P. The missing term in effective pair potentials. J. Phys. Chem. 1987, 91, 6269–6271.
- <span id="page-22-10"></span>(47) Abascal, J. L.; Vega, C. A general purpose model for the condensed phases of water: TIP4P/2005. J. Chem. Phys. 2005, 123 .
- <span id="page-23-0"></span>(48) Piana, S.; Donchev, A. G.; Robustelli, P.; Shaw, D. E. Water dispersion interactions strongly influence simulated structural properties of disordered protein states. J. Phys. Chem. B 2015, 119, 5113–5123.
- <span id="page-23-1"></span>(49) Leonarski, F.; D'Ascenzo, L.; Auffinger, P.  $Mg^{2+}$  ions: Do they bind to nucleobase nitrogens? Nucleic Acids Res. 2017, 45, 987–1004.