1	Supplementary information for		
2	Microfluidic Manipulation by Spiral Hollow-Fiber Actuators		
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16			
17			
18	This PDF file includes:		
19			
20	Supplementary Note 1. Methods (Pages S2–S3)		
21	Supplementary Note 2. Thermomechanical modeling (Pages S3–S7)		
22	Supplementary Note 3. Optimum hollow fiber geometry analysis (Pages S7–S8)		
23	Supplementary Figures 1 to 12 (Pages S9–S15)		
24	Supplementary Tables 1 to 3 (Pages S16-17)		
25	References (Page S18)		
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27 Supplementary Note 1. Methods

28 1.1 Characterizations

Linear coefficients of thermal expansion (CTE) of PEHF were measured from 0 to 100 °C at a heating rate of 3 °C min⁻¹ using a thermomechanical analyzer (TMA, Q400 system, TA instruments). Because the thermal expansion of employed polyethylene hollow fibre is anisotropic¹, the axial and the radial coefficients of thermal expansion were tested individually. For each experiment, three heating scans and two cooling scans were conducted. The linear CTE along each direction was calculated based on the thermomechanical analysis results.

Two-dimensional wide-angle X-ray scattering (2D WAXS) experiments were 36 performed at a sample-to-detector distance of 100 mm using CuK α radiation ($\lambda = 1.54$ 37 38 Å) by deploying a D8 Discover system (Bruker, Germany) equipped with a 2D HI-Star detector and a 1.6 kW generator. Parallel PEHF were assembled, using a 39 video microscope to ensure that the X-ray beam irradiated the centers of the fibres. 40 The image acquisition time was 600 s. One-dimensional wide-angle X-ray scattering 41 42 (1D WAXS) experiments were performed using CuK α radiation and a 2 θ scanning rate of 5° min⁻¹ on a Rigaku Smart Lab system (Japan) equipped with a 3.0 kW 43 generator. 44

45 1.2 Calculation of the degree of orientation of crystals using WAXS analysis

Despite inserting twists into a polymer fibre, the highly oriented semicrystalline form is also the key to provide a fibre the actuation capability. The Herman's orientation parameter (*f*) is employed to represent the polymer chain alignment of the fibre, where f = 1 means that the polymer chains orient along the fibre axis direction completely; f = -0.5 means that the polymer chains orient vertically to the fibre axis direction completely. For a non-oriented sample, f equals to 0.

52 The Herman's orientation parameters (f_a , f_b , and f_c) for the a-axis, b-axis, and 53 c-axis, for a unit cell of polyethylene crystallite with orthorhombic symmetry, were 54 calculated as 0.12, -0.30, and 0.18, respectively, according to the literatures.²⁻³ These 55 values provide the evidences that the fibres are in highly oriented semicrystalline 56 form.

59

57 The calculation details are as follows.

58 The Herman's orientation parameter (f) is defined as:

 $f = \frac{3\langle \cos^2 \Phi \rangle - 1}{2} \tag{1.1},$

(1.2),

60 where $\langle \cos^2 \Phi \rangle$ is an orientation factor defined as:

61
$$\langle \cos^2 \Phi \rangle = \frac{\int_0^{\pi} I_{\Phi} \cos^2 \Phi \sin \Phi \mathrm{d}\Phi}{\int_0^{\pi} I_{\Phi} \sin \Phi \mathrm{d}\Phi}$$

62 where I_{Φ} represents the diffraction intensity. The azimuthal curve obtained by 63 integration of diffraction intensity in the 2D WAXS pattern is used in Eq. (1.2) to 64 calculate $\cos^2 \Phi$ by the numerical integration process. Then the values of $\cos^2 \Phi$ 65 are used in Eq. (1.1) to calculate the orientation function value *f* of the semicrystalline 66 polymers.

For polyethylene with orthorhombic symmetry of the unit cell, the a-axis (f_a) and b-axis (f_b) orientation parameters were calculated from the intense reflections of (200) and (020) planes based on the relation:

70 $\langle \cos^2 \Phi_{110} \rangle = 0.692 \langle \cos^2 \Phi_{020} \rangle + 0.308 \langle \cos^2 \Phi_{200} \rangle$ (1.3),

where $\langle \cos^2 \Phi_{110} \rangle$ and $\langle \cos^2 \Phi_{200} \rangle$ are obtained from the azimuthal intensity measurements on the (110) and (200) reflections, respectively. The parameter f_c is estimated based on the orthogonal symmetry relationship: $f_a + f_b + f_c = 0.4$

The detailed calculations for f_a , f_b , and f_c are as follows. First, an azimuthal 74 75 curve was obtained by integration of diffraction intensity in the 2D WAXS pattern 76 over the azimuthal angle using the software of the GADDS (general area detector diffraction system) for diffraction angles $(2\theta_d)$ of 21.6° (110 plane) and 24° (200 77 78 plane), as shown in Supplementary Fig. 1c. Two identical peaks were observed in the 79 azimuthal curve (Supplementary Fig.1d) because of the symmetric diffraction pattern in 2D WAXS. The f_a , f_b , and f_c of the PEHF are respectively 0.12, -0.30, and 0.18, 80 which indicate that the fibres are in highly oriented semicrystalline form. 81

83 Supplementary Note 2. Thermomechanical modeling of hollow fibre actuator

84 1.1 Transient heat conduction analysis

The temperature change of the hollow fibre actuator is contributed by the forced 85 86 convection of the internal flowing water as well as the natural convection of the 87 ambient (external) air. To study this transient heat conduction problem, we establish a 88 two-step model. First, we model the radial heat conduction at an arbitrary cross-section of the hollow fibre actuator. Second, we build a simple axial heat 89 90 conduction model based on the assumption that heat conduction rate along the hollow 91 fibre axis is much slower compared to the water flow rate and the rate of heat 92 conduction in the radial direction. With this two-step model, we determine the hollow 93 fibre actuator temperature as a function of radial and axial coordinates and time.

The DSC results showed that there was no observable phase transition in the PEHF for the investigated temperature range (Supplementary Fig. 1b), we assume that there is no contribution of the phase transition to the volume change and the heat transfer contributed to the volume expansion.

98 First, the governing equation of the transient heat conduction in radial direction99 at an arbitrary cross-section of the hollow fibre is given by

100
$$\frac{\partial^2 \theta}{\partial r^2} + \frac{1}{r} \frac{\partial \theta}{\partial r} = \frac{1}{\alpha} \frac{\partial \theta}{\partial t} \quad \text{in } R_i \le r \le R_o, \qquad (2.1)$$

101 where θ (*r*, *t*) is the temperature, R_i and R_o are the inner and outer radius of the hollow 102 fibre, respectively. Thermal diffusivity α is defined by

103
$$\alpha = \frac{k}{\rho c_p}, \qquad (2.2)$$

104 in which k is the thermal conductivity, ρ is the mass density, c_p is the specific heat 105 capacity⁵⁻⁶. For PE, we have k = 0.33 W (m·K)⁻¹,⁷ $\rho = 920$ kg m⁻³,⁸ $c_p = 2300$ J 106 (kg·K)⁻¹⁸.

107 The convective boundary conditions at the interior and exterior surfaces of the108 hollow fibre are given by

109
$$k \frac{\partial \theta}{\partial r}\Big|_{r=R_i} = h_w \Big(\theta\Big|_{r=R_i} - \theta_w\Big), \qquad (2.3)$$

110
$$-k \frac{\partial \theta}{\partial r}\Big|_{r=R_o} = h_a \Big(\theta\Big|_{r=R_o} - \theta_a\Big), \qquad (2.4)$$

111 where h_w and h_a are the heat transfer coefficients of water and air, respectively. 112 Similarly, θ_w and θ_a are respectively the temperatures of water and air. The heat 113 transfer coefficient of water depends on many factors such as hollow fibre internal 114 geometry, surface roughness, water physical properties as well as the flow state 115 (laminar or turbulent), *etc.* In this work, we employ $h_w = 2000 \text{ W} (\text{m}^2 \cdot \text{K})^{-1}$, which 116 provides the best fit of theoretical predictions to the experimental measurements.

117 The initial condition is

$$\theta(t=0) = \theta_0, \qquad (2.5)$$

119 where θ_0 is the initial temperature of the hollow fibre actuator.

The above 1D radial transient heat conduction equations can be solved either analytically using methods of separation of variables and superposition, which are well documented in ref. 10, or using numerical methods such as finite difference or finite element. In this work, we solved this 1D partial differential equation (PDE) using the PDE solver provided by MATLAB^{*}.

Next, we assume that the axial heat conduction rate is much lower compared to the water flow rate. As such, the heat transfer in the axial direction is mainly contributed by the flowing water and the radial conduction. Therefore, by substituting water temperature as a function of axial coordinate and time, i.e. $\theta_w = \theta_w(x,t)$, in which x is the axial coordinate, into the above radial heat conduction equations, we can solve the hollow fibre temperature distribution in the axial direction.

131 Water temperature $\theta_w(x, t)$ depends on a few factors, such as initial water 132 temperature, water flow rate, heat loss due to transfer to hollow fibre actuator and

^{*}https://www.mathworks.com/help/matlab/math/partial-differential-equations.html

eventually to the ambient air, *etc*. For simplicity, we assume that the heat loss is
relatively low when the flow rate is high. Therefore, we can write the water
temperature as

136
$$\theta_{w}(x,t) = \theta_{w0}H(vt-x), \qquad (2.6)$$

137 where θ_{w0} is the initial water temperature, *v* is the water flowrate and $H(\cdot)$ is the 138 Heaviside step function.

139

144

140 *2.2 Torsional actuation modeling*

141 Torsional actuation mechanism of the twisted polymer fibres was studied in ref. 142 10 and 11. The hollow fibre length l, the hollow fibre outer diameter d, and the 143 helically oriented polymer chain length λ are related through

$$\lambda^2 = \left[\left(\pi dT \right)^2 + 1 \right] l^2 , \qquad (2.7)$$

145 where *T* is the twist density.

146 Upon heating, polymer chains will expand in the radial direction and contract in 147 the axial direction, which can be accommodated by the changes in the hollow fibre 148 length, the diameter and the twist density. In ref. 10, the change of hollow fibre length 149 is considered small and the following equation for untwisting can be obtained by

150
$$\Delta T = \left(\frac{\Delta \lambda}{\lambda} \frac{1}{\cos^2 \alpha_f} - \frac{\Delta d}{d}\right) T, \qquad (2.8)$$

151 where bias angle $\alpha_f = \tan^{-1}(\pi dT)^{12}$.

152 Ref. 11 further assumed that λ is also constant. Therefore, the untwist ΔT is 153 linearly proportional to the change in the hollow fibre diameter, i.e.

154
$$\Delta T = -\frac{\Delta d}{d}T, \qquad (2.9)$$

155 In this work, we employ Eq. (2.8) for torsional actuation modeling. It can be 156 rewritten as

157
$$\Delta T = \left(\alpha_{\lambda} \frac{1}{\cos^2 \alpha_f} - \alpha_d\right) \Delta \theta \cdot T, \qquad (2.10)$$

where α_{λ} and α_{d} are the coefficients of thermal expansion of polymer chain in the 158 axial direction and that of the hollow fibre in the radial direction, respectively. The $\Delta\theta$ 159 is the temperature change obtained from the previous transient heat conduction 160 analysis. In this work, the average exterior hollow fibre temperature change is 161 employed as the actuation temperature $\Delta \theta$. Axial and radial coefficients of thermal 162 were employed as $\alpha_{\lambda} = -(5.3 \pm 0.4) \times 10^{-4} \text{ K}^{-1}$ and expansion of PEHFs 163 $\alpha_d = (5.2 \pm 0.7) \times 10^{-4} \text{ K}^{-1}$, which were obtained from thermomechanical analysis 164 (TMA). 165

- 166
- 167 *2.3 Tensile actuation modeling*

The tensile actuation mechanism of the twisted, mandrel-coiled polymer fibre is
well documented in ref. 10, which is adopted in this work. The equation for tensile
actuation is

171
$$\frac{\Delta L}{L} = \frac{l^2}{NL} \Delta T \quad , \tag{2.11}$$

where *L* is the coil length, *l* is the twisted hollow fibre length, *N* is the number of coil turns and ΔT is the amount of untwisting that can be obtained from Eq. (2.10).

175 Supplementary Note 3. The analysis for the optimum hollow fibre geometry

To provide an estimation of the optimum configuration of the hollow fibre, we assume steady heat solution and constant temperature due to the small thickness of the hollow fibre. Based on the linear elastic solution for a thin-walled hollow fibre under pure torsion¹³, the generated torque *M* is related to the twist through

180

$$M = 2\pi G J T, \tag{3.1}$$

181 in which *G* is shear modulus, *J* is the polar second moment of area. For the hollow 182 fibre section $J = \frac{\pi}{2} (r_o^4 - r_i^4)$ where r_i and r_o are the inner and outer radius of the 183 hollow fibre, respectively.

184 Applying Eq. (3.1) to the hollow fibre configurations before and after the

185 torsional actuation provides the torque release ΔM :

186
$$\Delta M = \pi G A_0 \left(r_i^2 + r_o^2 \right) \left[(1 + \alpha_d \Delta \theta)^4 (T + \Delta T) - T \right], \qquad (3.2)$$

187 in which A_0 is the cross-section area of the hollow fibre.

188 We now introduce a simplified version of torsional actuation based on Ref. 10 :

$$\Delta T = \left(\frac{d_0}{d} - 1\right) T. \tag{3.3}$$

190 Substituting Eq. (3.3) into (3.2) gives

191
$$\Delta M = \pi G A_0 (r_i^2 + r_o^2) [(I + \alpha_d \Delta \theta)^3 - 1] T.$$
(3.4)

Based on Eq. (3.4), we conclude that the optimum geometry for maximizing the 192 torque output is to maximize $(r_i^2 + r_o^2)$, assuming that the cross-section area A_0 (or 193 equivalently the material consumption) is constant. This means that the material 194 195 should be distributed as far away as possible from the center of the cross-section, which justifies the use of hollow fibre over solid cross-sections. It should be noted 196 197 that the above theoretically derived optimum geometry is subjected to the constraint 198 that the hollow fibre will not buckle under the applied torque. Thus, in practice a 199 hollow fibre with finite thickness should be used for the proposed actuation 200 application.

201





Supplementary Figure 1. (a) The stress-strain curve of the PEHF₅₈₀₋₉₉₀. The gauge 204 length was 34 mm, and the extension rate was 20 mm min⁻¹. (b) Differential scanning 205 calorimetry (DSC) curve of the PEHF₅₈₀₋₉₉₀ with a heating rate of 5°C/min in N_2 206 atmosphere. The melting enthalpy ($H_{\rm m}$ =126.4 J g⁻¹) was obtained by integration of the 207 endothermic melting peak. By considering that the melting enthalpy of the 100% 208 crystalline polyethylene was $H_c=287.3$ J g^{-1 14}, we can obtain the crystallinity of 209 low-density polyethylene (χ_c) from the equation $\chi_c = H_m/H_c = 44\%$. (c) The intensity of 210 211 diffraction peaks versus Bragg angle (2 θ) in WAXS for PEHF₅₈₀₋₉₉₀. (d) 2D WAXS 212 diffraction pattern for PEHF₅₈₀₋₉₉₀ showing anisotropic structure. (e, f) The azimuthal curves of (110) and (200) planes in PEHF₅₈₀₋₉₉₀ were obtained from (d). 213





215 Supplementary Figure 2. The torsional actuation of the PEHF₅₈₀₋₉₉₀ actuator by flowing hot water at 1.72 g s⁻¹ at different actuation conditions: (a) at different 216 environment temperatures, (b) at different environmental relative humidity, and (c) at 217 different wind speed. (d) The torsional actuation of the PEHF₅₈₀₋₉₉₀ actuator by 218 flowing different types of 40 °C liquid at a flow rate of 1.72 g s⁻¹. The δ and $\Delta\delta$ are the 219 220 rotation angle during actuation and its change. If not specified, the room temperature is 25 °C, the environmental relative humidity is 40%, and the water flow rate is 1.72 g 221 s^{-1} . If not specified, error bars in this figure and in the following figures indicate 222 standard deviations. EtOH: ethanol; EA: ethyl acetate; DMSO: dimethyl sulfoxide; 223 tetrahydrofuran; DMF: N, N-dimethylformamide; DCM: dichloromethane; 224 THF: 225 IPA: iso-propanol; PE: petroleum ether.



Supplementary Figure 3. (a) Contraction and hollow fibre surface temperature as a function of time for a homochiral PEHF₅₈₀₋₉₉₀ actuator by flowing 95 °C water with different flow rates. (b) Response time and hollow fibre surface temperature as a function of water flow rate for a homochiral PEHF₅₈₀₋₉₉₀ actuator by flowing 95 °C water. (c) Coil length and actuation stroke of the homochiral PEHF₅₈₀₋₉₉₀ actuator by

flowing 25 °C water at different flow rates. The spring index was 4.0, and the inserted
twist density was 200 turns m⁻¹.



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Supplementary Figure 4. (a) The initial coil length and actuation stroke of the 236 homochiral PEHF₅₈₀₋₉₉₀ actuator by flowing 25 °C glycerol/water solution with 237 238 different viscosity. (b) Actuation stroke of the PEHF₅₈₀₋₉₉₀ actuator by flowing 80 °C glycerol/water solution with different viscosity. (c) The initial coil length and 239 240 actuation stroke of the PEHF₅₈₀₋₉₉₀ actuator by flowing 25 °C aqueous NaCl solution with different densities. (d) Actuation stroke of the PEHF₅₈₀₋₉₉₀ actuator by flowing 241 80 °C aqueous NaCl solution with different densities. The spring index was 4.0, and 242 the twist density was 200 turns m⁻¹. 243



245 **Supplementary Figure 5**. The contraction as a function of temperature for 246 PEHF₅₈₀₋₉₉₀ actuator driven by consecutively flowing hot water and cold water with a 247 temperature range. The spring index was 4.0, the inserted twist density was 200 turns

248 m^{-1} , and the flow rate was 1.72 g s⁻¹. The solid symbols indicate the heating process, 249 and the open symbols indicate the cooling process.



250

Supplementary Figure 6. (a, b) Serial infrared images of the homochiral
PEHF₃₈₀₋₁₀₉₀ (a) and PEHF₂₈₀₋₆₄₀ (b) actuators by flowing 95 °C water at different time
(0 s, 0.2 s, 0.4 s, and 0.6 s). (c) Experimental results and theoretical analysis of the
actuation strain and surface temperature of the PEHF actuators as a function of time.



Supplementary Figure 7. (a) Maximum actuation stroke for ten consecutive
heating/cooling cycles for heterochiral and homochiral load-free PEHF₅₈₀₋₉₉₀ actuators
driven by alternatively flowing hot water and 25 °C water. (b) Maximum actuation

stroke as a function of twist density for load-free $PEHF_{580-990}$ homochiral actuators with different spring indexes driven by 80 °C flowing water.



Supplementary Figure 8. (a) Maximum contraction at different mass of the load for ten consecutive heating/cooling cycles of homochiral PEHF₅₈₀₋₉₉₀ actuators with a spring index of 4.0 and an inserted twist of 100, 150, 200, 250 and 300 turns m⁻¹ by alternatively flowing 90 °C and 25 °C water. (b) The actuation stroke and the work

266 capacity as a function of the mass of the applied load for the homochiral PEHF₅₈₀₋₉₉₀ 267 actuators with different twist densities, which were derived from graphs in (a).



Supplementary Figure 9. (a) Comparison of the work capacity at different isobaric 269 stress for the homochiral PEHF₅₈₀₋₉₉₀ actuator driven by 90 °C water and for the PE 270 solid fibre actuator driven by 90 °C water. (b) Comparison of the time-dependence of 271 work capacity for the homochiral PEHF₅₈₀₋₉₉₀ actuator driven by 90 °C water and for 272 the PE solid fibre driven by 90 °C air at the stress of 107 kPa, and the inset was the 273 274 comparison of power density. The twist density was 300 turns m⁻¹, and the spring index was 4.0. The isobaric stress was calculated as the weight of the load divided by 275 276 the cross-sectional area of sheath of the hollow fibre.



277

Supplementary Figure 10. (a) Schematic diagram of a conventional microfluidic
manipulation system containing liquid sensing, transporting and actuating devices.
Optical images (b) and schematic illustration (c) of the contractile hollow fibre
actuators used for sensing the liquid temperature and sorting the liquid into the

desired vessels. (d) The length of the hollow fibre actuators as a function of the
temperature of the flowing water for the homochiral PEHF₅₈₀₋₉₉₀ actuators.



285 Supplementary Figure. 11 Schematic of the twisted and coiled hollow-fibre286 actuators



Supplementary Figure 12. The twist loss of the hollow fibre after twist insertion,
thermal annealing with both-end tethered at 108 °C for 1h, and cooling down to room
temperature and removing the tethering.

293 Supplementary Tables:

Matariala	Twist density	$\Delta heta$	Flowrate	Response time
Materials	(turns m ⁻¹)	(°C)	$(g s^{-1})$	(s)
PEHF580-990	100	70	0.5	2.5
PEHF580-990	200	70	0.5	2.1
PEHF580-990	300	70	0.5	1.9
PEHF580-990	400	70	0.5	1.7
PEHF580-990	400	55	0.5	1.8
PEHF580-990	400	45	0.366	2.7
PEHF580-990	400	35	0.327	1.1
PEHF580-990	400	5	1.72	1.0

Supplementary Table 1. Response time and actuation parameters of the torsionalPEHF actuators in this work.

296

297 Supplementary Table 2. Response time and actuation parameters of the tensile PEHF

actuators in this work.

Matariala	Twist density	$\Delta heta$	Flowrate	Response time
Materials	(turns m ⁻¹)	(°C)	$(g s^{-1})$	(s)
PEHF580-990	200	35	1.72	2.5
PEHF580-990	200	45	1.72	1.6
PEHF580-990	200	55	1.72	1.4
PEHF580-990	200	70	1.72	0.88
PEHF580-990	200	70	0.5	1.4
PEHF580-990	200	70	0.366	1.6
PEHF580-990	200	70	0.327	1.7
PEHF580-990	200	70	0.245	2.8
PEHF ₂₈₀₋₆₄₀	310	70	0.366	1
PEHF380-1090	182	70	0.366	3.0

299

301 Supplementary Table 3. Resolution and the coil parameters of the homochiral PEHF

Fibre Length	Coil Length	Number of	Twist density	Resolution
(<i>l</i> , mm)	(<i>L</i> , mm)	Turns (N)	$(T, turns mm^{-1})$	$(l^2 \Delta T / NL, 10^{-3})$
200	105	11	0.30	-3.0%
200	105	11	0.25	-2.0%
200	105	11	0.20	-1.3%
200	105	11	0.15	-0.78%
200	105	11	0.10	-0.43%
200	115	12	0.30	-2.5%
200	115	12	0.25	-1.7%
200	115	12	0.20	-1.1%
200	115	12	0.15	-0.66%
200	115	12	0.10	-0.36%
200	130	15	0.25	-1.2%
200	130	15	0.20	-0.78%
200	130	15	0.10	-0.26%
200	145	18	0.20	-0.58%
200	145	18	0.10	-0.19%

302 hollow fibre actuators in this work.

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