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Supplementary Materials for

Metallic *Mimosa pudica*: A 3D biomimetic buckling structure made of metallic glasses

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Movies S1 to S4

Supplementary Text

1. Laser parallel-line patterning

The *ns*-pulsed laser parallel-line scanning and the processing geometry are schematically illustrated in **fig. S1**, where *x*, *y*, *z* directions stand for length, width and thickness (*d*) of the ribbon, respectively, *s* refers to the laser line spacing (distance between centers of two laser lines), *w* is the laser linewidth. Patterning started from one corner of the ribbon, moved in a zigzag path, and ended in the opposite corner. The laser source was only ON while scanning the ribbon along the *x*-direction but OFF when moving from one line to another line. The procedure is termed parallel-line patterning.



Fig. S1. Schematic of the parallel-line patterning.

The maximum temperature reached during laser patterning can be estimated by

using Dowden's model presented in *ref. 50*. A laser source is simplified as a point source at the surface of a semi-infinite workpiece. The maximum temperature (T_{max}) can be estimated according to:

$$T_{\max} = T_0 + \frac{P}{2\pi\lambda r} \exp\left\{\frac{v}{2\kappa}(x-r)\right\}$$
(S1)

where the parameters are: T_0 – ambient temperature; P – incident power; λ – thermal conductivity; v – scanning speed; κ – thermal diffusivity; the origin is set as the central position of the laser spot; r the distance from any position in the sample to the origin; x= $r\cos\theta$; θ the angle of r to the x-direction. By taking P = 5 W (which counts for the influence of the reflectivity/absorption (50%) of MG and the percentage of power we applied), $T_0 = 300$ K, v = 0.1 m s⁻¹, r = 50 µm (half of the spot size), x = 50 µm, $\lambda =$ 10.6 W m⁻¹ K⁻¹ and $\kappa = 2.58$ mm² s⁻¹ (λ and κ of a similar MG composition, Fe₈₀B₂₀, are taken from *ref.* 51), then this gives $T_{max} = 1802$ K, which is above the melting temperature 1410 K, obtained independently by high-temperature calorimetry of Fe₇₈Si₉B₁₃ MG suggesting that the material was melted under the laser-scanning conditions for successful buckling.

2. Demonstration of the buckling

R of the naturally buckled roll is 10.75 mm, and of the reversed roll is 11.25 mm (**fig. S2**). This difference in *R* can be explained by a small difference in *w* at the top and bottom surfaces: *w* at the bottom was about 5% smaller than that at the top surface, which is consistent with the previous observations (*52*). All the *R* given in the figures were measured for naturally buckled ribbons, i.e., immediately after laser-line

patterning.



Fig. S2. Reversible 3D buckled structures (at *s* of 1.5 mm, *v* of 100 mm s⁻¹). (A) Concave and (B) convex geometries of the buckled ribbon; (C) Naturally buckled roll and (D) reversed roll. *R*: the curvature radius of MGC composite.

3. Analytical solution of R using the 2D model

The *R* of the buckled ribbon is determined by the energy minimum of the sum of bending and stretching energy (Equations 1, 2 & 3 in the manuscript). It includes 7 inputs, f_g , f_x , E_g , E_x , w_g , w_x , and d, and 3 variables, R, r_g and r_x . Here, f is the relative length change of the striped material, E the elastic modulus, w the linewidth, R the radius of curvature for the bulked structure along the laser lines direction, and r the radius of curvature for the modulated curves transverse to laser lines direction. In particular, f_x/f_g can be estimated by the density ratio between the two stripes, i.e., $f_x/f_g = (\rho_x/\rho_g)^{1/3}$; and $w_x = w$, $w_g = s - w$. For Fe-Si-B MG, w is 0.18 mm, s is 1.5 mm, d is 0.024 mm, E_g is 168.9 GPa, E_x is 217.8 GPa, ρ_g is 7.154 g cm⁻³, and ρ_x is 7.375 g cm⁻³. To obtain R corresponding to the minimum sum of energy, a MATLAB program

"fmincon" is utilized to solve the multivariable (R, r_g and r_x) function (Eq. (1)-(3)),

together with the "multistart" solver for multiple local minima calculation. To see the trend of *R* vs. ρ_x/ρ_g , we fixed the E_x/E_g and w/s at default values and allowed ρ_x/ρ_g to change. For *R* vs. E_x/E_g and *R* vs. w/s, similar approaches were applied, too. The results are shown in **fig. S3**.



Fig. S3 The *R* is calculated by energy-minimum calculation. (A) *R* vs. ρ_x/ρ_g , (B) *R* vs. E_x/E_g , and (C) *R* vs. *w/s*.

From **fig. S3**, it can be seen that *R* decreases with increasing ρ_x/ρ_g and/or E_x/E_g ; an optimum $w/s \sim 0.4$ exists corresponding to a minimum *R* (**fig. S3C**). While w/s is too large or too small the *R* goes to infinity, in these cases, buckling does not occur.

4. Kerr analysis

Fig. S4 demonstrates the evolution of magnetic domains upon the application of external tensile stress along the direction of the ribbon. In the stress-free state (fig. S4A), the ribbon had highly disordered domains governed by strongly varied local internal stresses. On applying external stress, the domains got more ordered and aligned along with the increasing stress principal (fig. S4B to D). The increasing ordering was accompanied by an increasing domains width. This kind of domains re-ordering were

expected due to the positive magnetostriction constant of the as-spun Fe₇₈Si₉B₁₃ ribbon.



Fig. S4. Kerr microscopy overview of the magnetic domains on the surface of the as-spun Fe₇₈Si₉B₁₃ MG ribbon, part A, upon the application of increasing tensile stress along the ribbon axis, parts B-D, which agrees to *ref. 49*. The Kerr sensitivity is horizontal with the respect to the image plane (black arrow). The applied force, the direction is indicated by the red arrow and the tensile stress values are given in the labels, is oriented along the same axis.

The magnetic domains, shown in **fig. S5** in the as-spun state (A) and after successive laser patterning (B), demonstrates how far the stress propagates in the ribbon. The well-oriented in-plane wide domains induced by the tensile stress can be found as far as several millimeters away from the laser line.



Fig. S5. The magnetic domains on the surface of Fe₇₈Si₉B₁₃ ribbon ($\nu = 100 \text{ mm s}^{-1}$) contain a single laser line (enveloped by yellow dashed line). (A) As-spun ribbon, and (B) the ribbon after laser-patterning. The Kerr sensitivity is along the vertical direction (transverse to laser scanning) concerning the image plane (black arrow).

5. Experimental *R* of buckling structure

To further compare the extent of buckling, the curvature radius of composite structure (*R*) of the buckled ribbons was evaluated. Laser scanning speed *v* influences crystal stripe width *w* (**fig. S6**). *R* is a function of *w* and *v* (see **fig. S7**). At a constant *s* of 2 mm, *R* increases with faster *v* from 40 to 120 mm s⁻¹. Similarly, at a constant *v* of 100 mm s⁻¹, the minimum spacing to obtain apparent buckling was 1.5 mm for *R* = 10 mm. *R* firstly increased with larger *s* and then leveled off at a constant value ~34 mm at *s* of about 10 mm. The *R* can be controlled by the ratio of *w*/*s*. Fig. 4F plots the dependence of *R* on the *w*/*s* ratio, and a universal correlation was observed. When *w*/*s* was higher than 25%, the ribbon buckled slightly due to the excessive penetration. Also, for *w*/*s* < 1%, buckling was not detected because the contribution of the crystalline part

is too small comparing with the line spacing. Therefore, for the ribbons of 24 μ m thick used in this work, the achievable range of *R* is between 10–34 mm.



Fig. S6 w increases with decreasing v.



Fig. S7 s and v influences R of buckling structures.

Fig. S8 shows the 3D shape-change of $Fe_{78}Si_9B_{13}$ scanned ribbons after thermal annealing (5 min at 743-753 K). The ribbon was flattened after annealing, attributed to the removal of density difference between the relaxed glass and the crystallized regions.



Fig. S8 Flattening of the buckled Fe₇₈**Si**₉**B**₁₃ **ribbon after annealing. (A)** As-spun ribbon via the laser patterning. (B) The buckled ribbon in (A) after thermal annealing (5 min at 743-753 K).

6. Structural and compositional characterizations of buckling ribbons



Fig. S9. Buckling of different composite ribbons. (A) Buckling shapes of $Fe_{78}Si_9B_{13}$, $Cu_{46}Zr_{46}Al_8$, $La_{55}Ni_{20}Al_{25}$ MGs and $Fe_{80}Cr_{20}$ polycrystal alloy ribbon after one-line laser patterning. The laser scanning speed v was set at 100 mm s⁻¹. (B) XRD of the laser scanned and unscanned ribbons. (C) XRD of as-spun and the laser scanned Fe₇₈Si₉B₁₃ ribbons at different v.

Fig. S9A shows the buckling structure of different ribbons after laser patterning of one line. Fe₇₈Si₉B₁₃ and Cu₄₆Zr₄₆Al₈ MGs show small *R* denoting evident buckling. But La₅₅Ni₂₀Al₂₅ MG and Fe₈₀Cr₂₀ polycrystal alloy show slightly buckling. **Table S1**

shows the density and elastic modulus change of Fe₇₈Si₉B₁₃, and Zr₄₆Cu₄₆Al₈ MGs before and after laser parallel-line processing. The successful buckling structure occurs in these two MGs, respectively. **Fig. S9B** presents the XRD patterns of the laser scanned and unscanned MG ribbons. Peaks of the crystallinities can be observed for all the samples after laser scanning. **Fig. S9C** shows the detected crystallinity of the buckled Fe₇₈Si₉B₁₃ ribbons. Taking the crystalline peak at $2\theta = 44.96$ ° for comparison, the size (*D*) of the crystals can be estimated by using the Scherrer equation, where $D = \frac{0.89\lambda_{\text{KRD}}}{2}$.

Bcosθ

Because the wavelength of X-ray is 1.54 Å, θ is 22.48 °, and the full-width-halfmaximum (*B*) is 0.345±0.023 ° for v = 100 mm s⁻¹ and 0.377±0.008 ° for v = 40 mm s⁻¹, which gives *D* of 24.6±1.8 nm and 22.5±0.6 nm for v = 100 mm s⁻¹ and 40 mm s⁻¹, respectively.

Fig. S10A shows slight buckling in Fe₇₈Si₉B₁₃ and Cu₄₆Zr₄₆Al₈ crystallized ribbon. The crystallized ribbons of the Fe₇₈Si₉B₁₃ and Cu₄₆Zr₄₆Al₈ before laser scanned are checked by XRD (**fig. S10B**) presenting polycrystal alloy, and then with one laser line post-scanned.



Figure S10 The slight buckling of the crystallized ribbon of Fe₇₈**Si**₉**B**₁₃ **and Zr**₄₆**Cu**₄₆**Al**₈ **laser post-scanned.** (A) Slight buckling of the crystallized ribbon. (B) The crystallized ribbon before laser scanned is examined by XRD.

Fig. S11 shows the SEM/EDS compositional maps of the laser-processed Fe₇₈Si₉B₁₃ MG ribbons. There is no detectable compositional difference between the MG matrix and the written line on the macroscopic length scale.



Fig. S11. Energy-dispersive X-ray analysis of the laser-processed Fe₇₈Si₉B₁₃ ribbon at $v = 100 \text{ mm s}^{-1}$. (A) SEM image of the top surface. Concentration mapping of (B) iron, (C) oxygen, (D) silicon, and (E) boron.

Table S1 Density and elastic modulus change of Fe₇₈Si₉B₁₃, and Zr₄₆Cu₄₆Al₈ MGs before and after laser parallel-line processing.

Materials	$ ho_{ m g}$ (g cm ⁻³)	$ ho_{\rm x}$ (g cm ⁻³)	$ ho_{ m x}/ ho_{ m g}$	Eg (GPa)	Ex (GPa)	$E_{\rm x}/E_{\rm g}$	Buckling
Fe ₇₈ Si ₉ B ₁₃	7.154±0.003	7.375±0.003	1.031	168.9±3.7	217.8±10.8	1.29	Evident
Zr ₄₆ Cu ₄₆ Al ₈	7.068±0.005	7.134±0.011	1.009	91.4±6.1	115.0±13.7	1.26	Evident

Movie S1 Metallic mimosa pudica opens and closes upon the applied external magnetic field stimuli. Metallic mimosa pudica opens and closes upon applied external magnetic field stimuli with various magnetic directions supplied by a magnet bar.

Movie S2 **The various postures of metallic mimosa pudica.** Seven convex and concave postures of the metallic petals are reshaped by manual.

Movie S3 Comparison of responses of laser patterned and un-laser patterned petals upon applied magnetic field stimuli. The laser patterned metallic mimosa pudica can open and close their petals upon magnetic stimuli easily, whereas the unlaser patterned petals can't do so.

Movie S4 **The laser patterning process.** The laser patterning process, the formation of 3D buckling structure of the Fe-based MG ribbon spontaneously, and easy rolling of the ribbon by manual.

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